

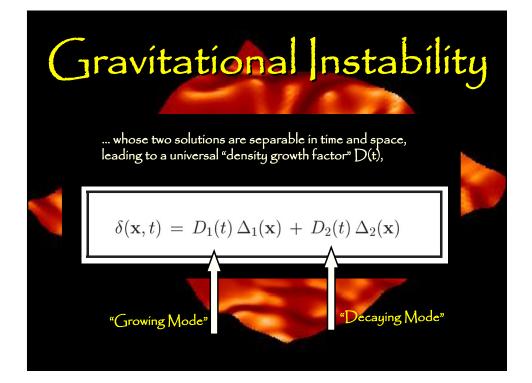
$$\begin{aligned} \frac{\partial \delta}{\partial t} + \frac{1}{a} \nabla \cdot (1 + \delta) \mathbf{v} &= 0\\ \frac{\partial \mathbf{v}}{\partial t} + \frac{\dot{a}}{a} \mathbf{v} + \frac{1}{a} (\mathbf{v} \cdot \nabla) \mathbf{v} &= -\frac{1}{a} \nabla \phi\\ \nabla^2 \phi &= 4\pi G \bar{\rho} a^2 \delta(\mathbf{x}, t) \end{aligned}$$

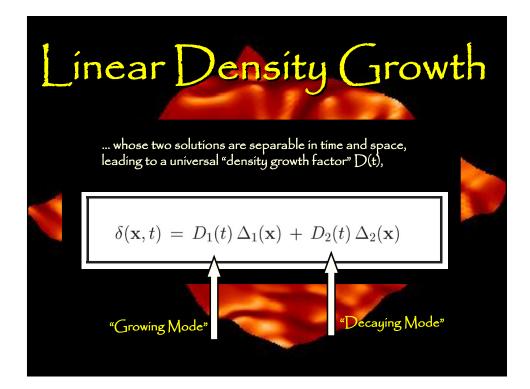
$$\begin{aligned} \frac{\partial \delta}{\partial t} + \frac{1}{a} \nabla \cdot \mathbf{v} &= 0\\ \frac{\partial \mathbf{v}}{\partial t} + \frac{\dot{a}}{a} \nabla \cdot \mathbf{v} &= 0\\ \frac{\partial \mathbf{v}}{\partial t} + \frac{\dot{a}}{a} \mathbf{v} &= -\frac{1}{a} \nabla \phi\\ \nabla^2 \phi &= \frac{3}{2} \Omega H^2 a^2 \, \delta(\mathbf{x}, t) \end{aligned}$$

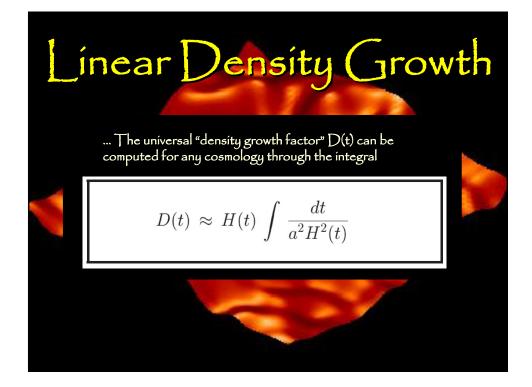
Gravitational Instability

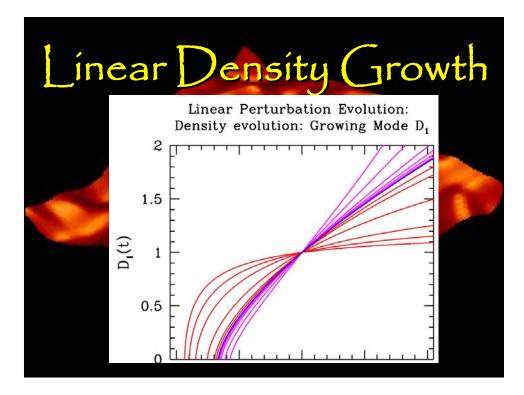
The linear system of structure growth equations can be written in terms of a second order differential equation,

$$\frac{\partial^2 \delta}{\partial t^2} + 2\frac{\dot{a}}{a}\frac{\partial \delta}{\partial t} = \frac{3}{2}\Omega_0 H_0^2 \frac{1}{a^3}\delta$$





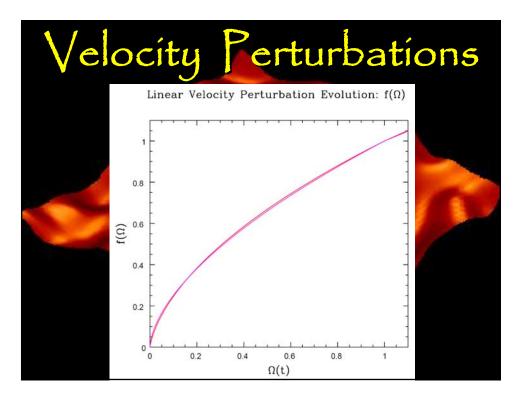


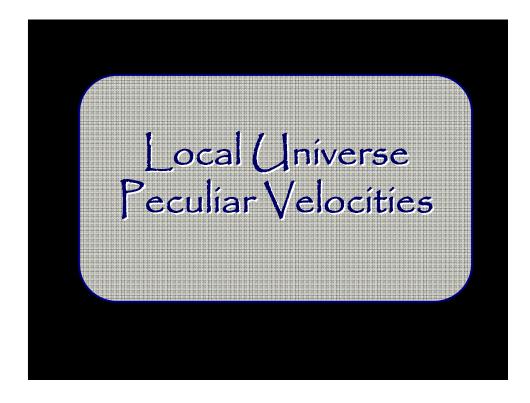


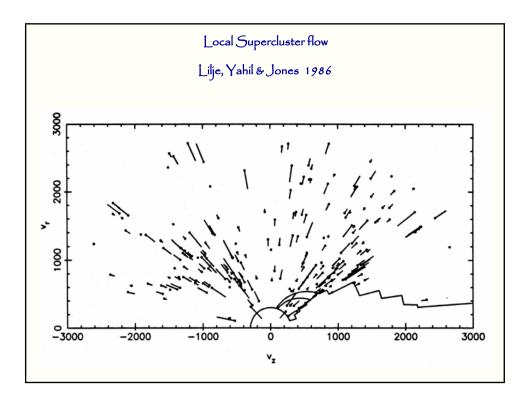
$$Gravitational Instability$$

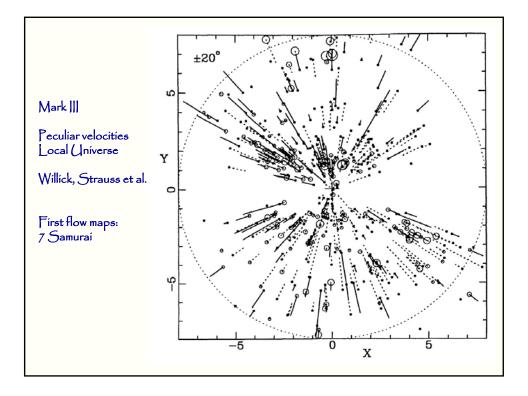
$$GRAVITY PERTURBATIONS$$

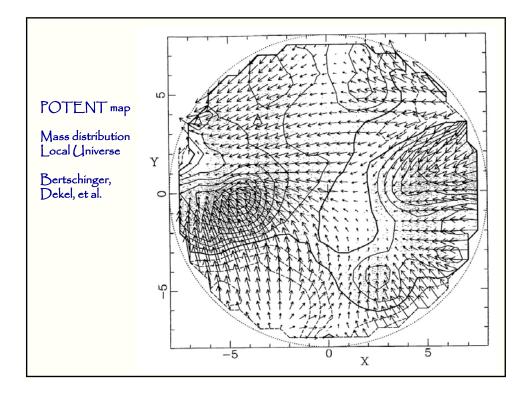
$$g(\mathbf{r}, t) = -\frac{1}{a} \nabla \phi = \frac{3\Omega H^2}{8\pi} \int d\mathbf{x}' \, \delta(\mathbf{x}', t) \frac{(\mathbf{x}' - \mathbf{x})}{|\mathbf{x}' - \mathbf{x}|^2}$$

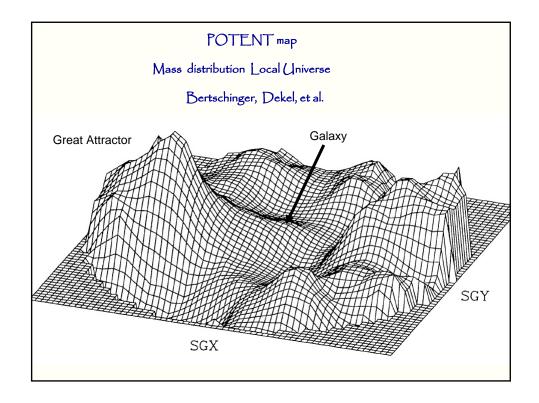


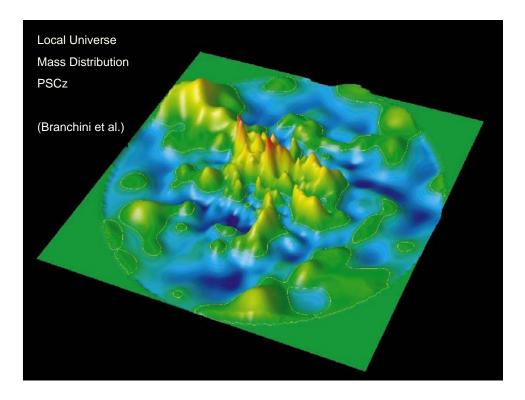




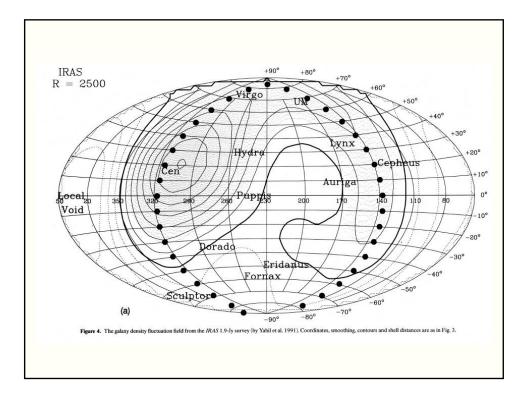


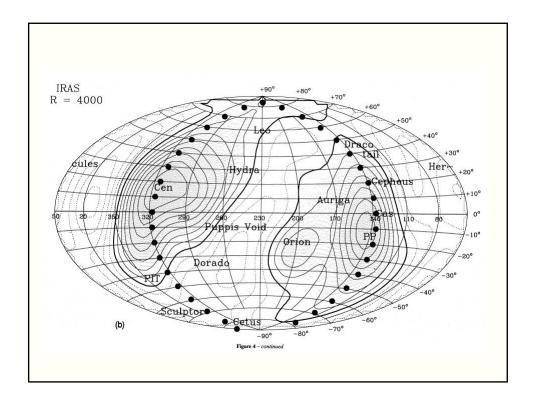


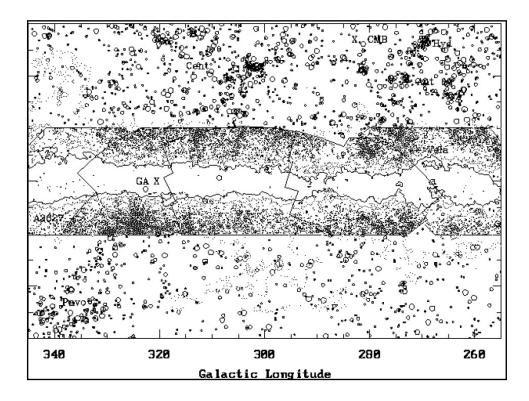


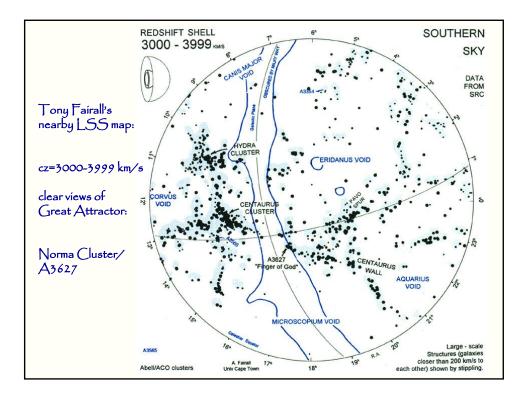


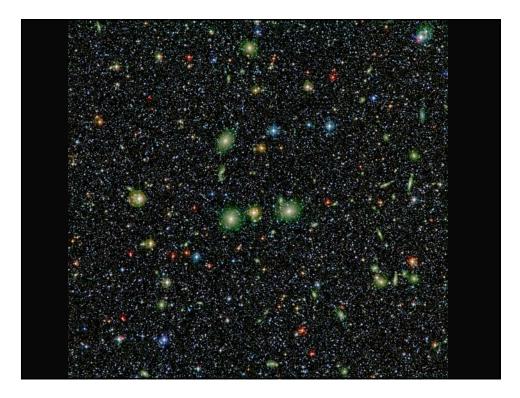












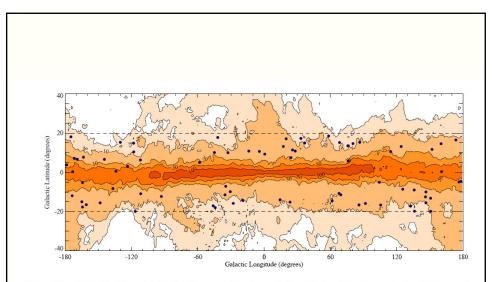


Fig. 16. Distribution in Galactic coordinates of the 76 by Ebeling et al. [129] so far spectroscopically confirmed X-ray clusters (solid dots) of which 80% were previously unknown. Superimposed are Galactic HI column densities in units of 10^{20} cm⁻² (Dickey & Lockman 1990). Note that the region of relatively high absorption ($N_{\rm HI} > 5 \times 10^{21}$ cm⁻²) actually is very narrow and that clusters could be identified to very low latitudes

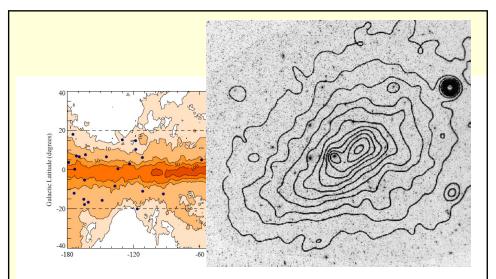


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