

An Analytic Expression for the Evolution of Voronoi Features

Vincent Icke

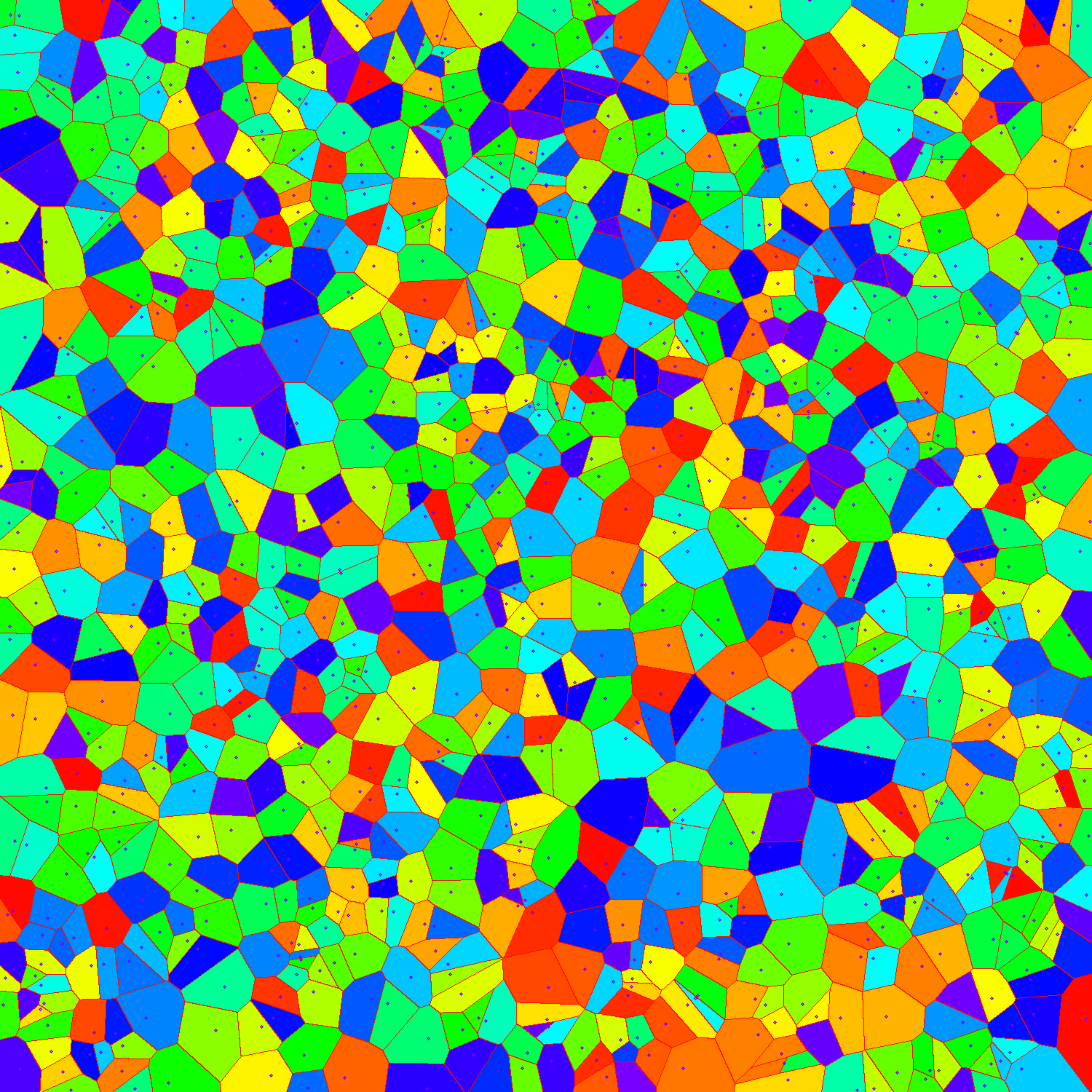
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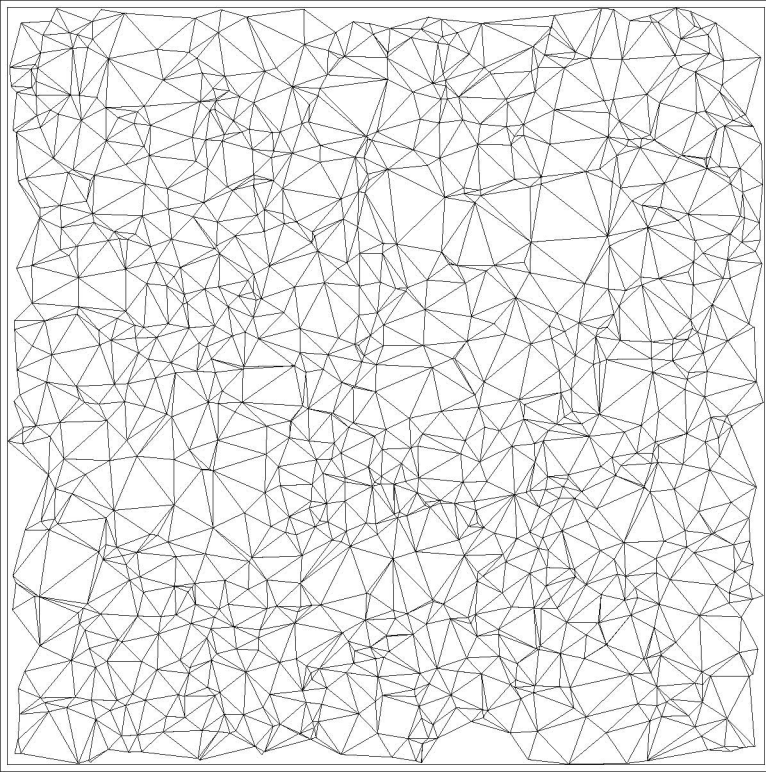


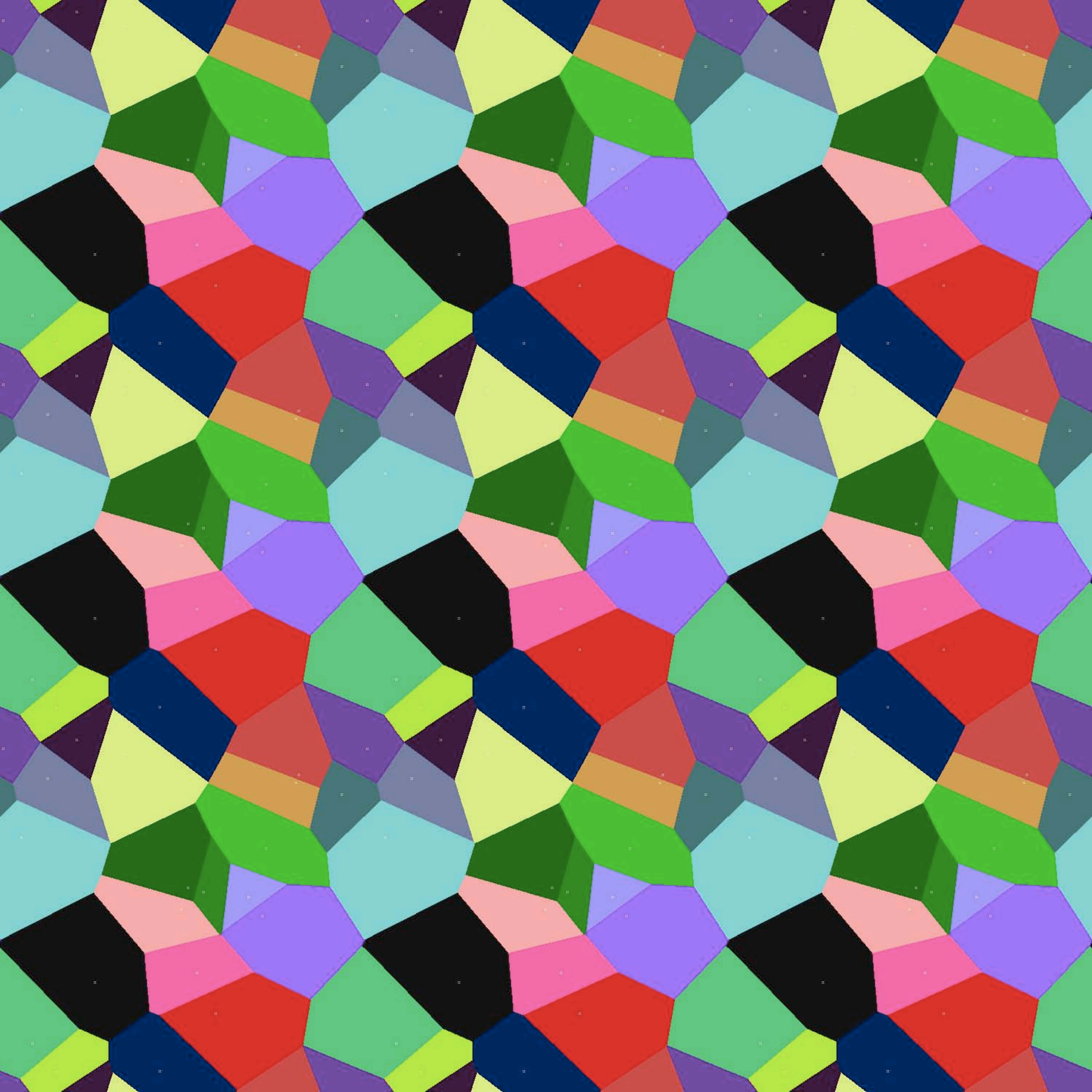
**Voronoi
Tessellation
Primer**

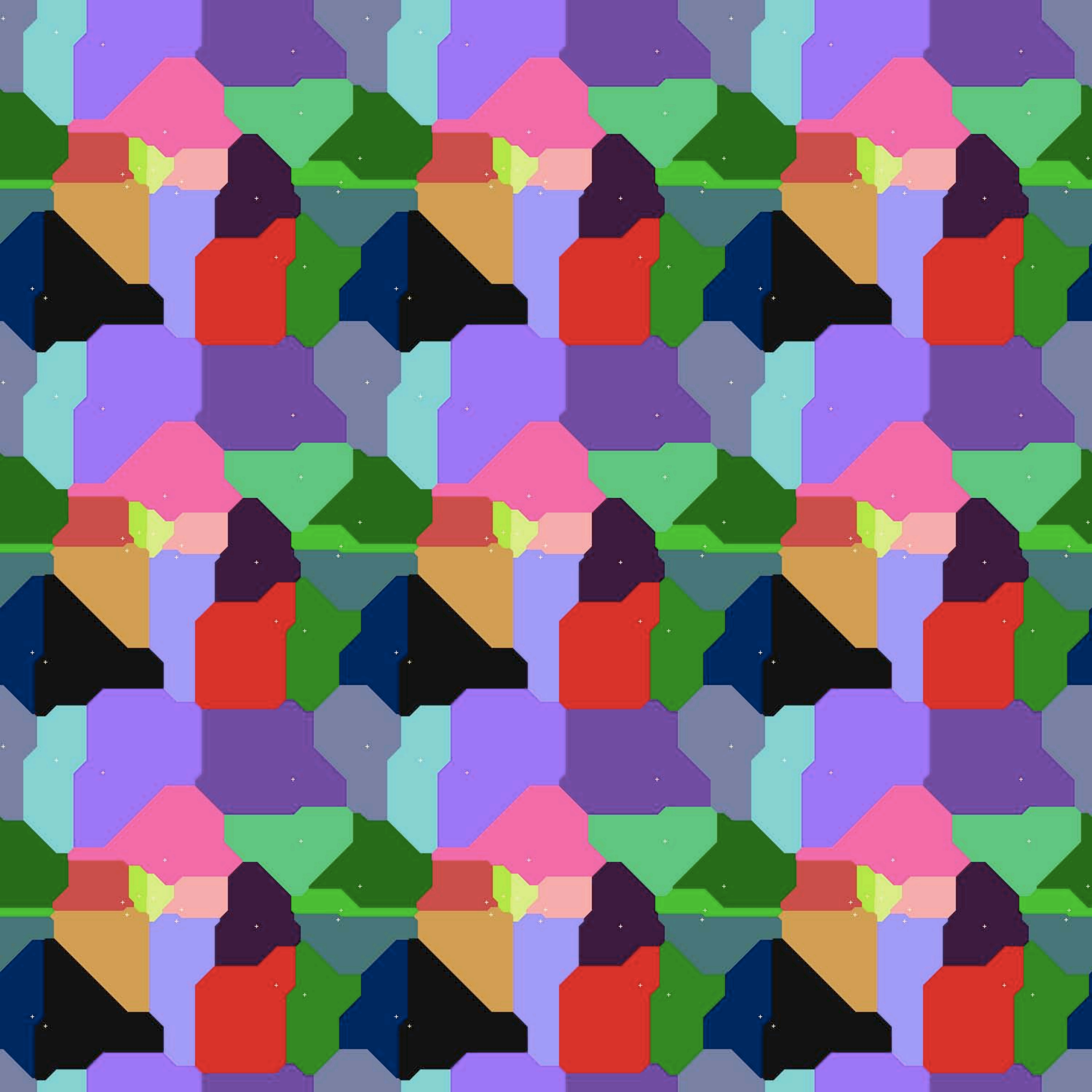
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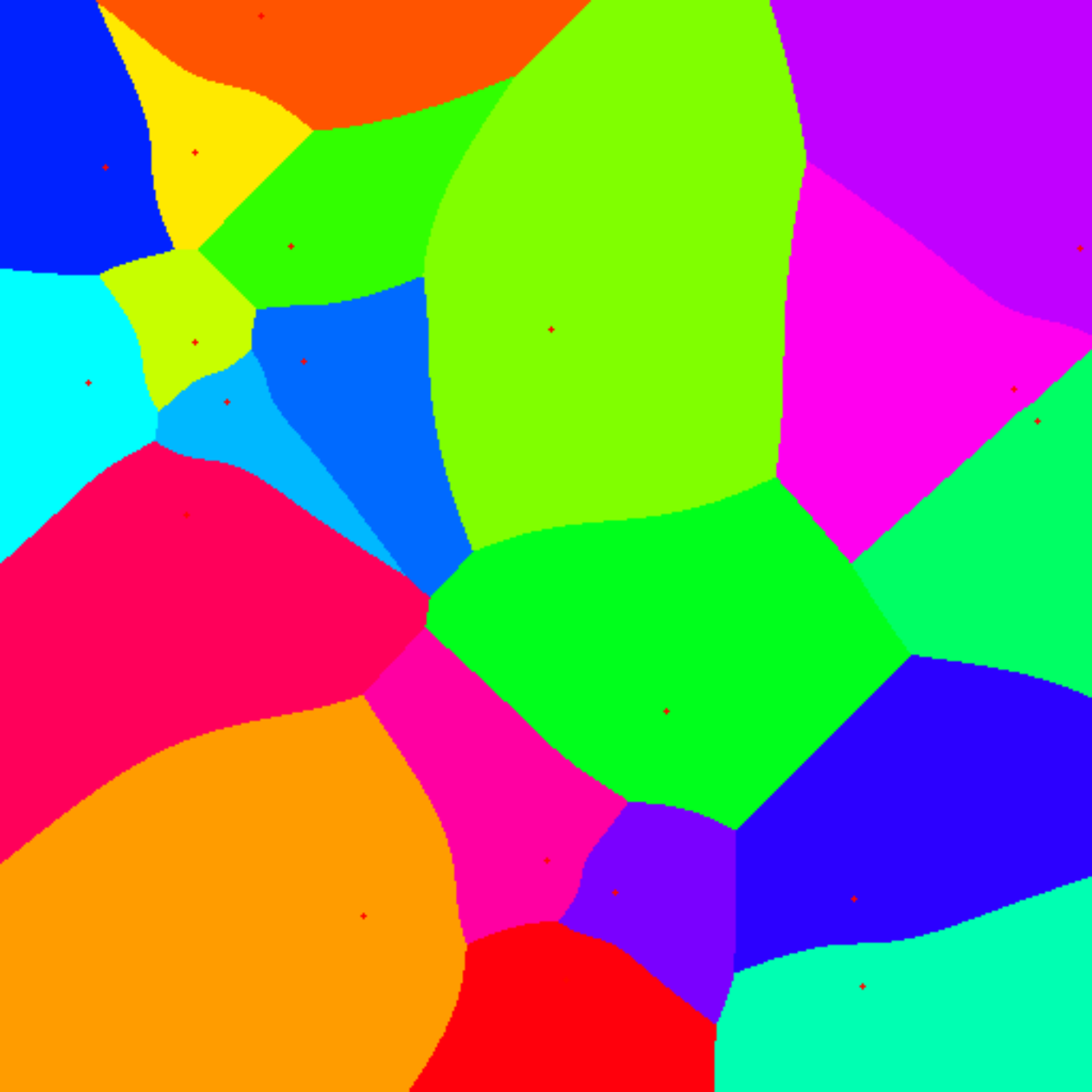


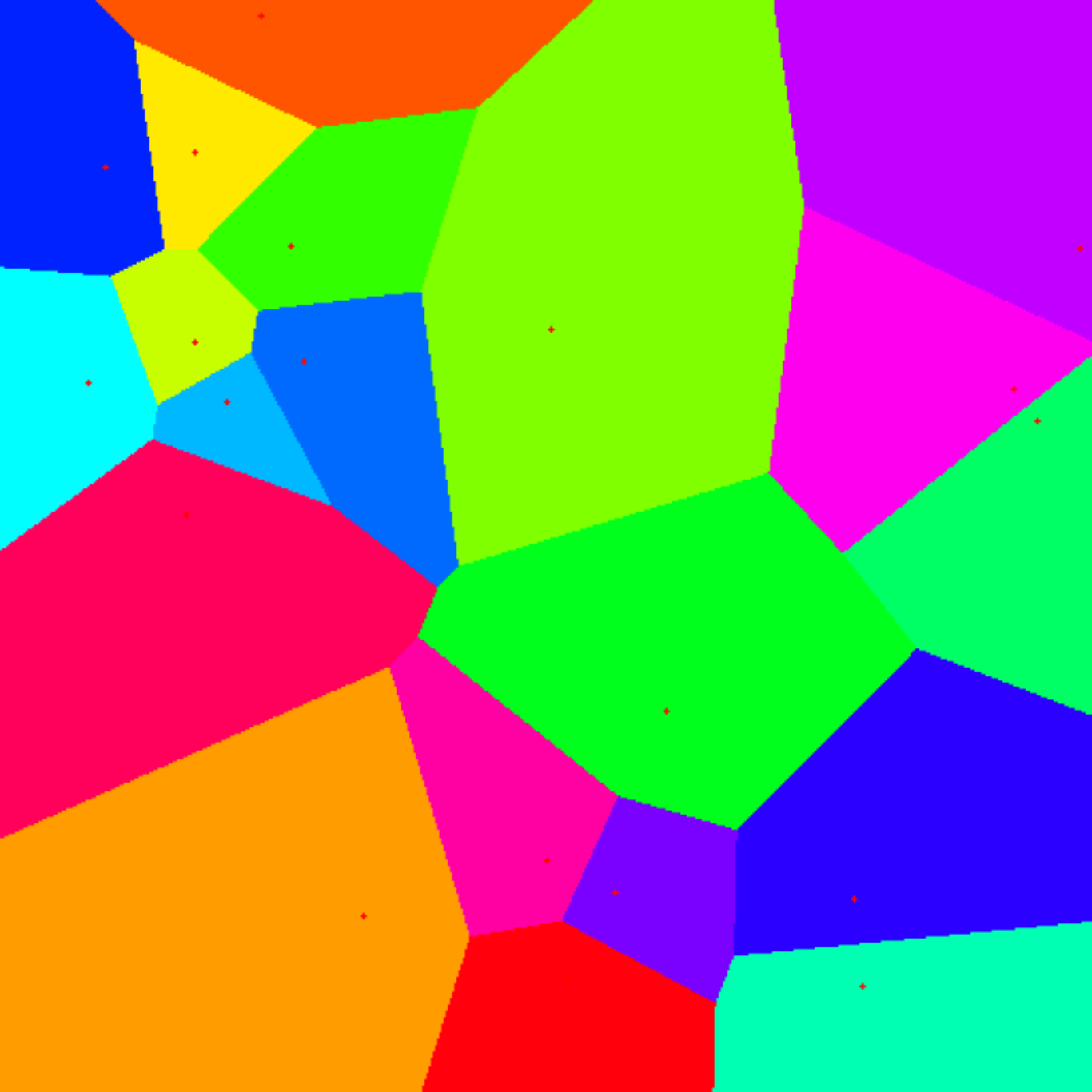


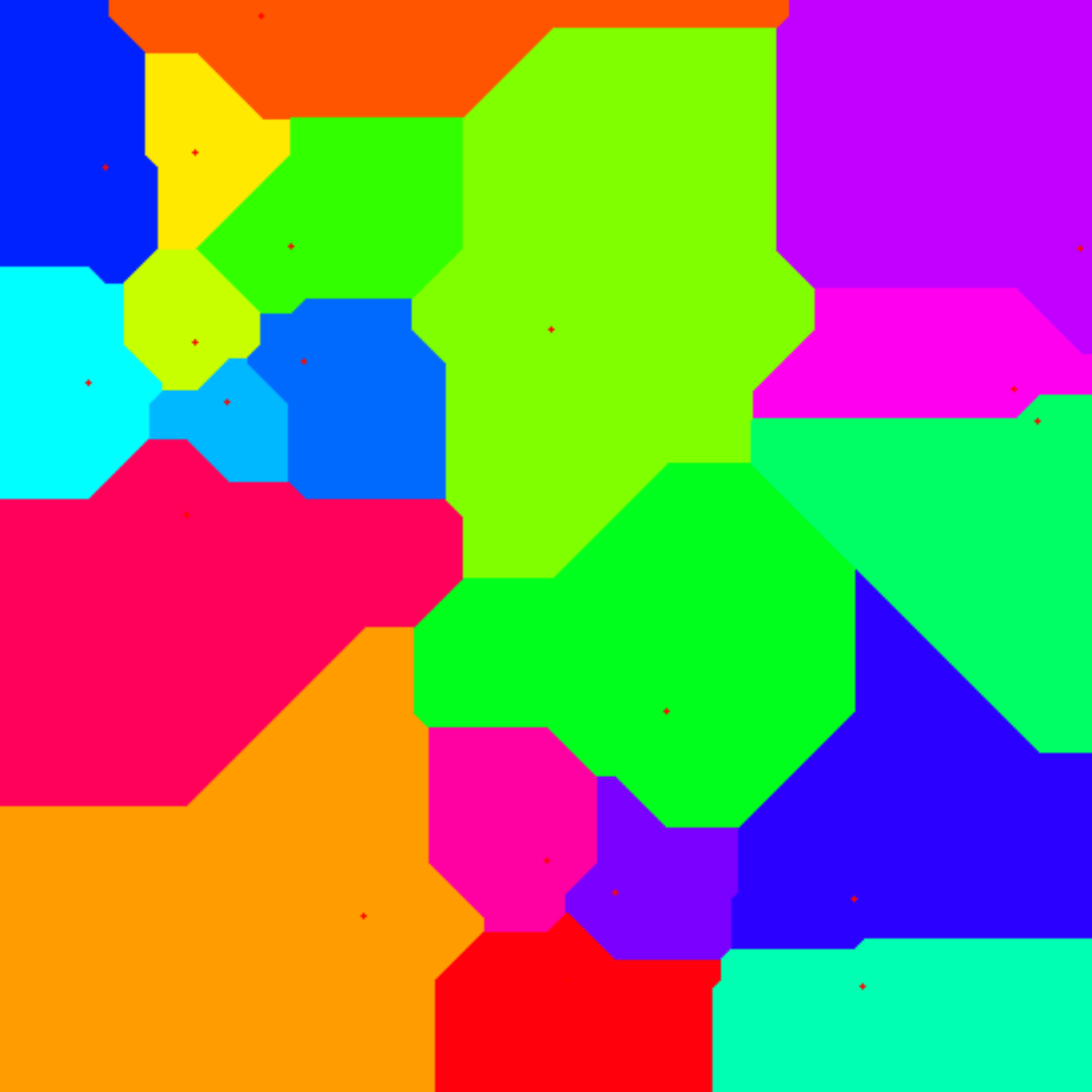


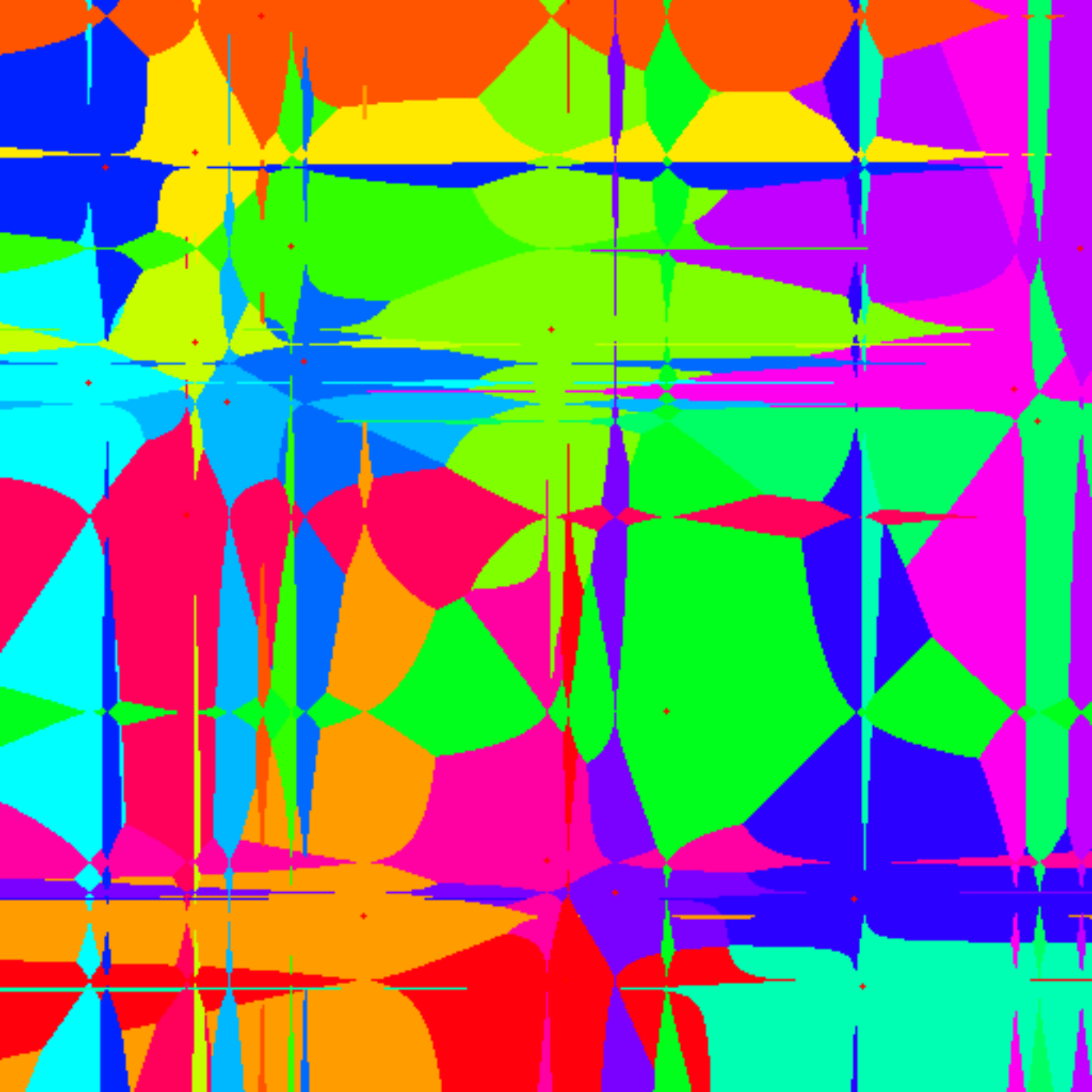


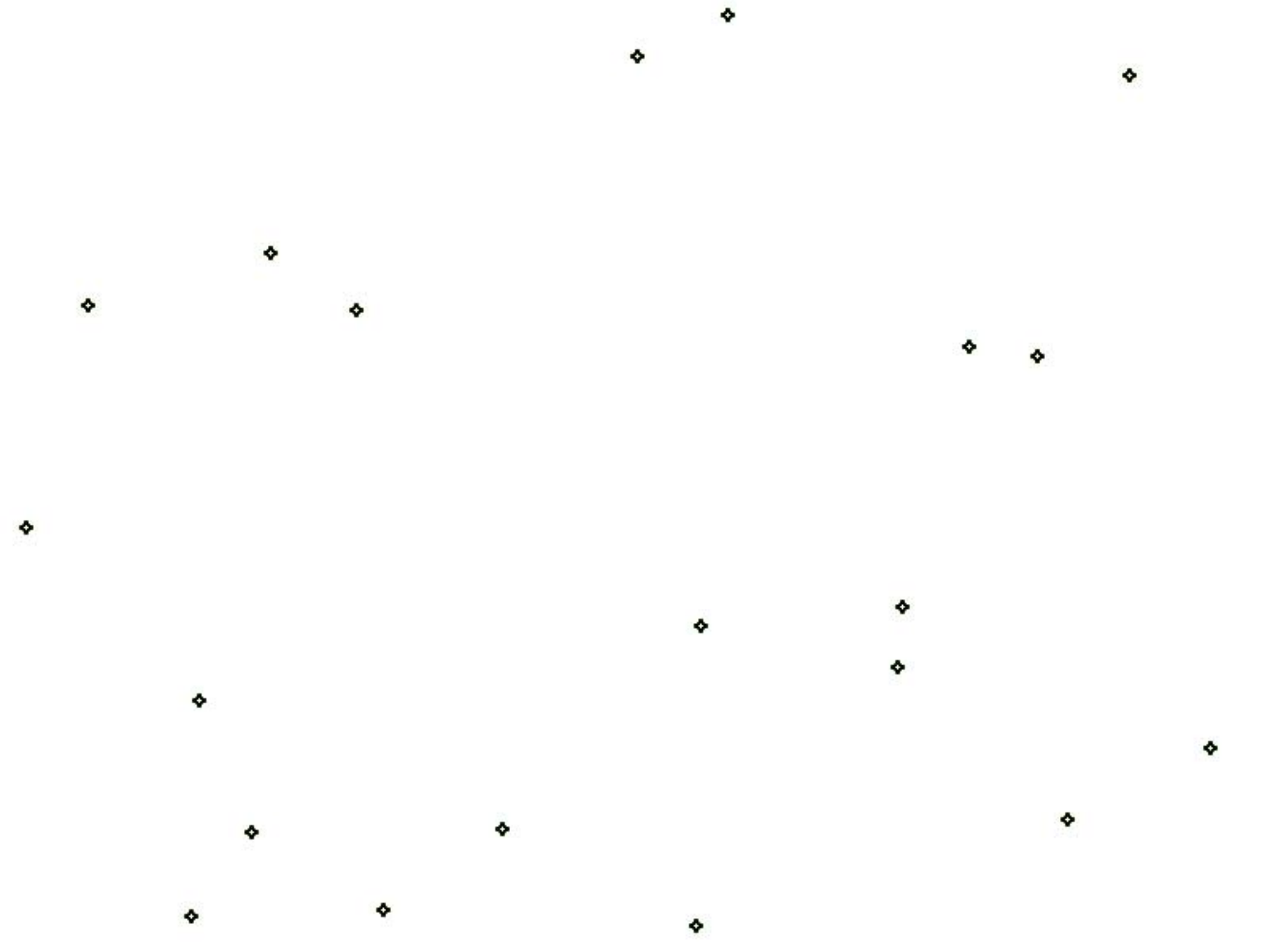


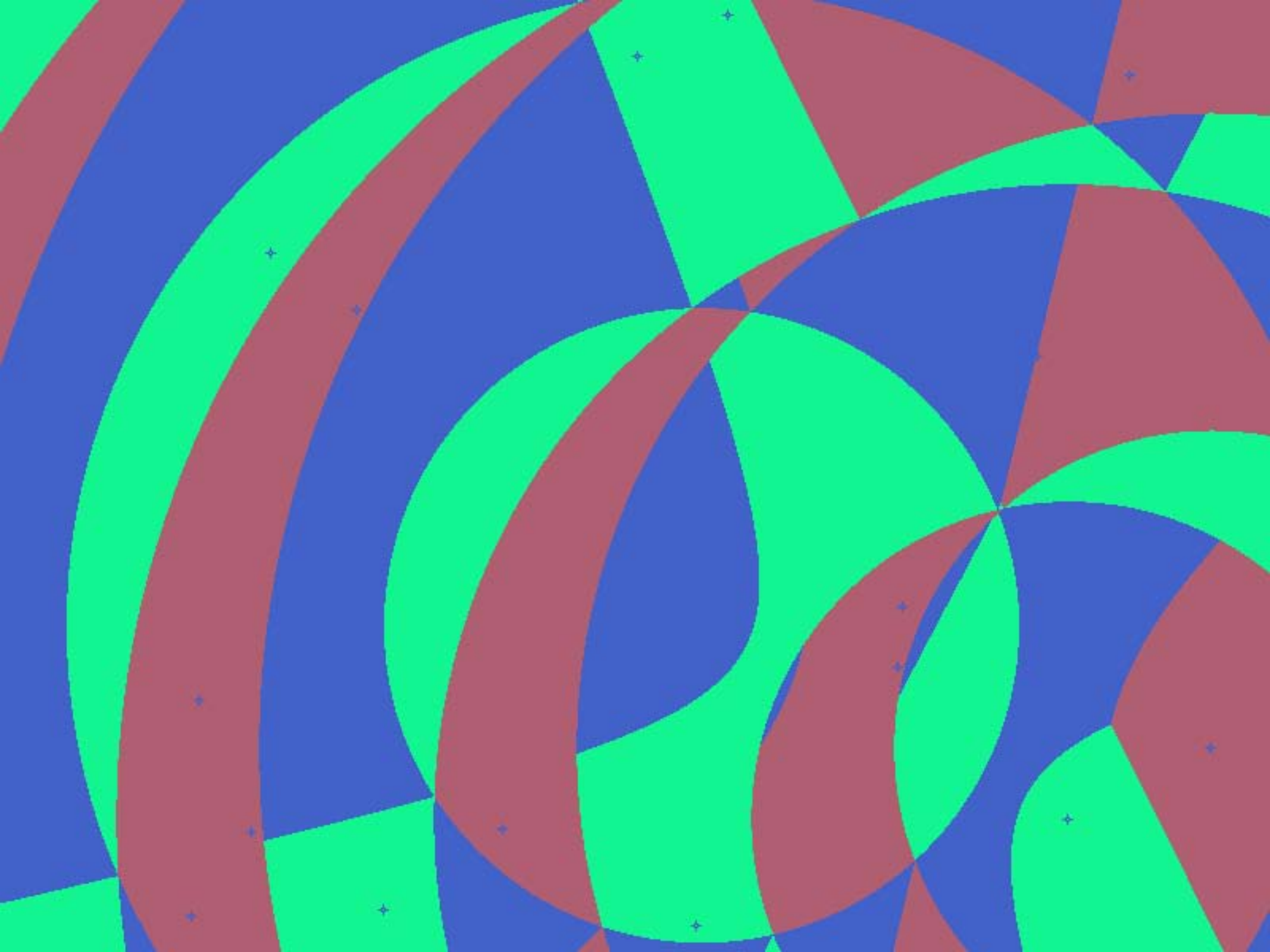


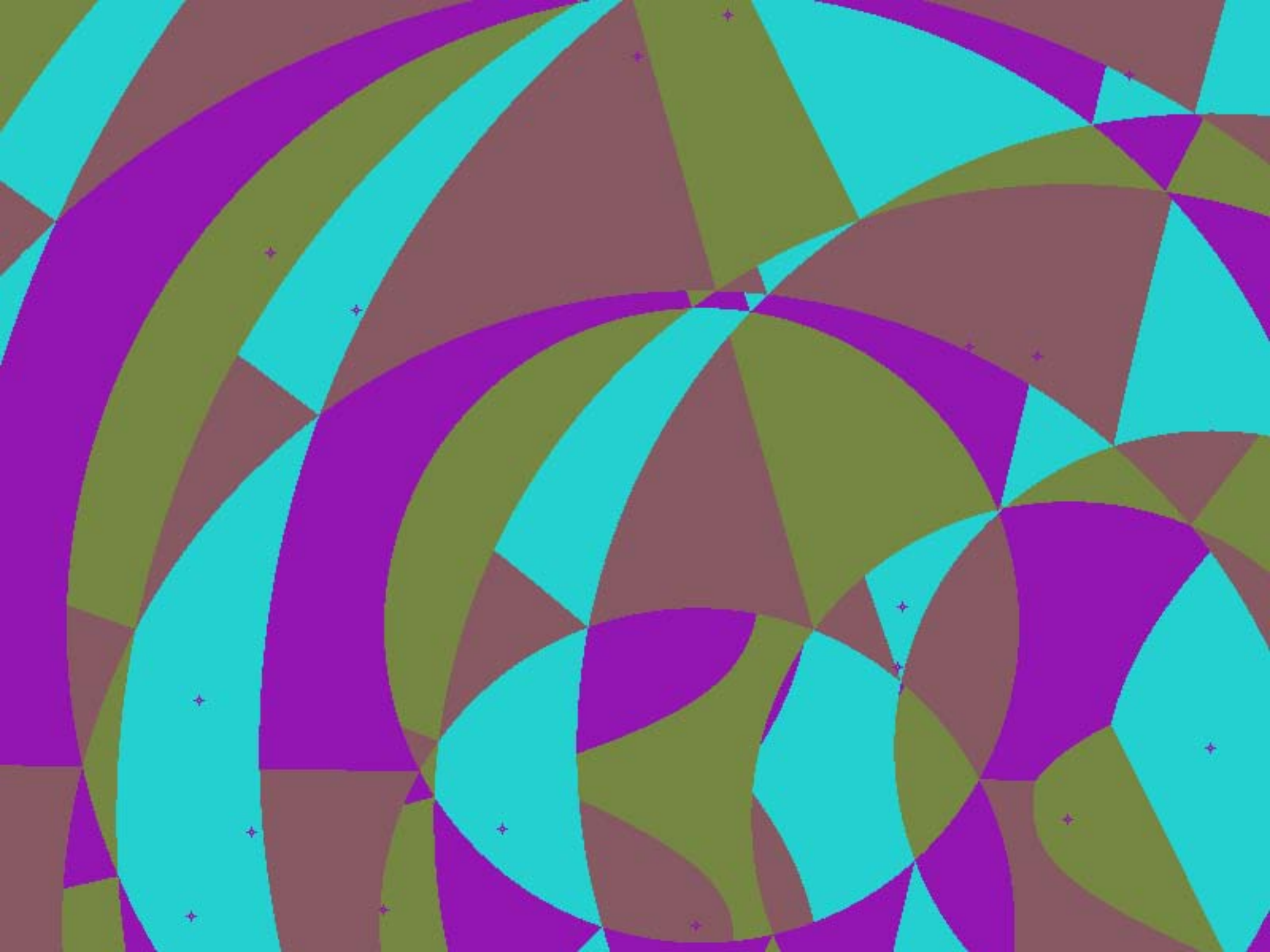


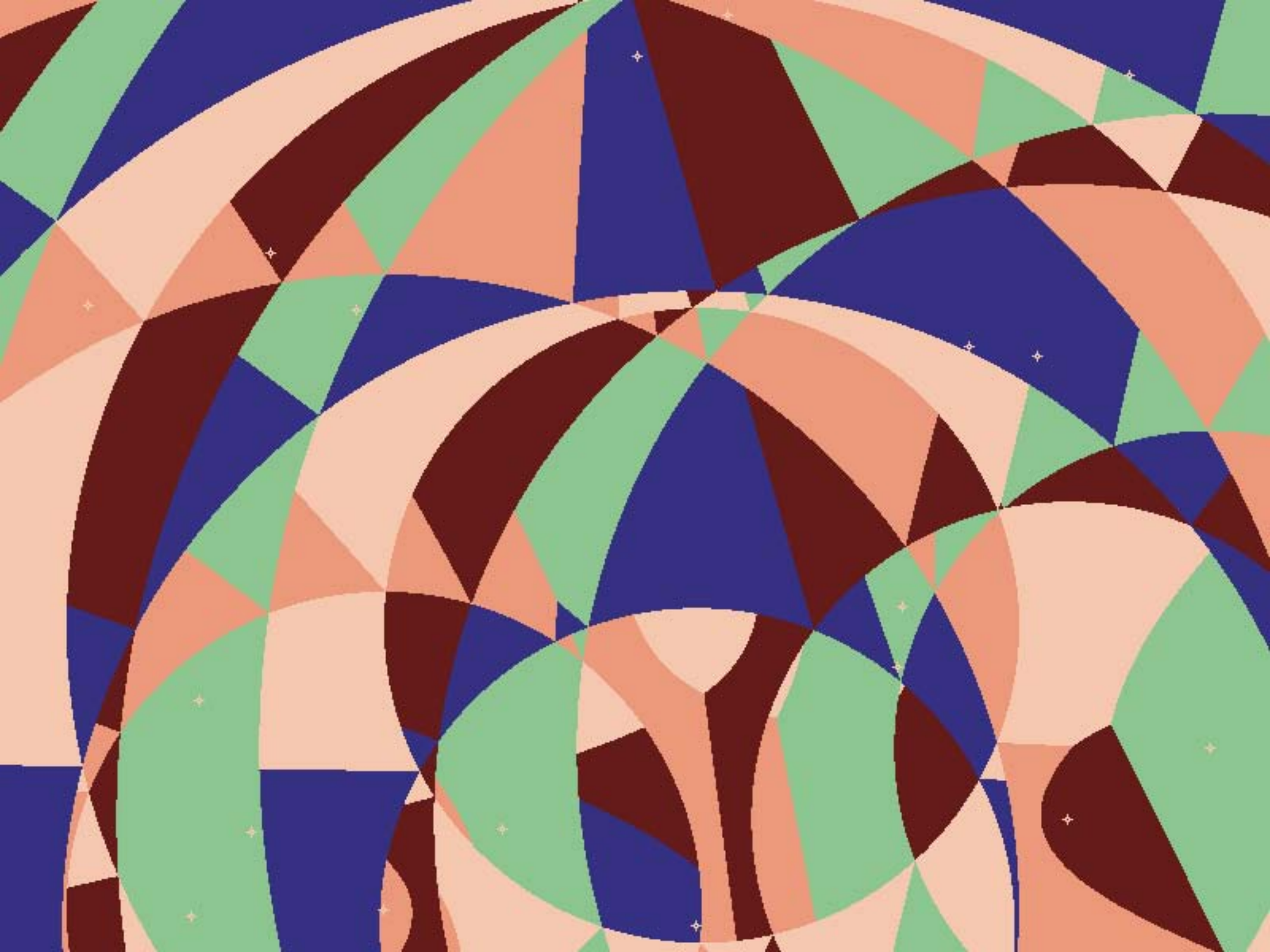


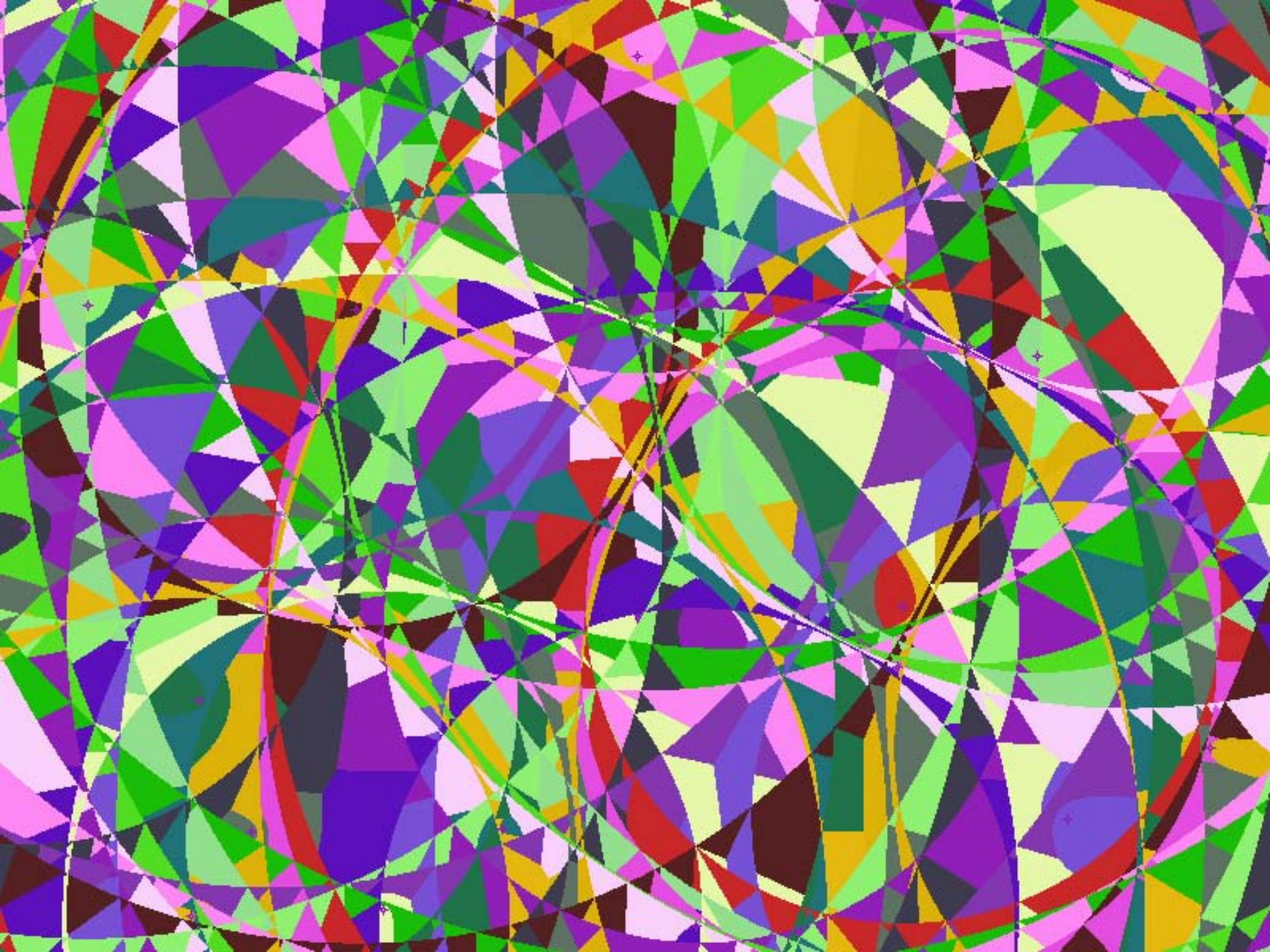


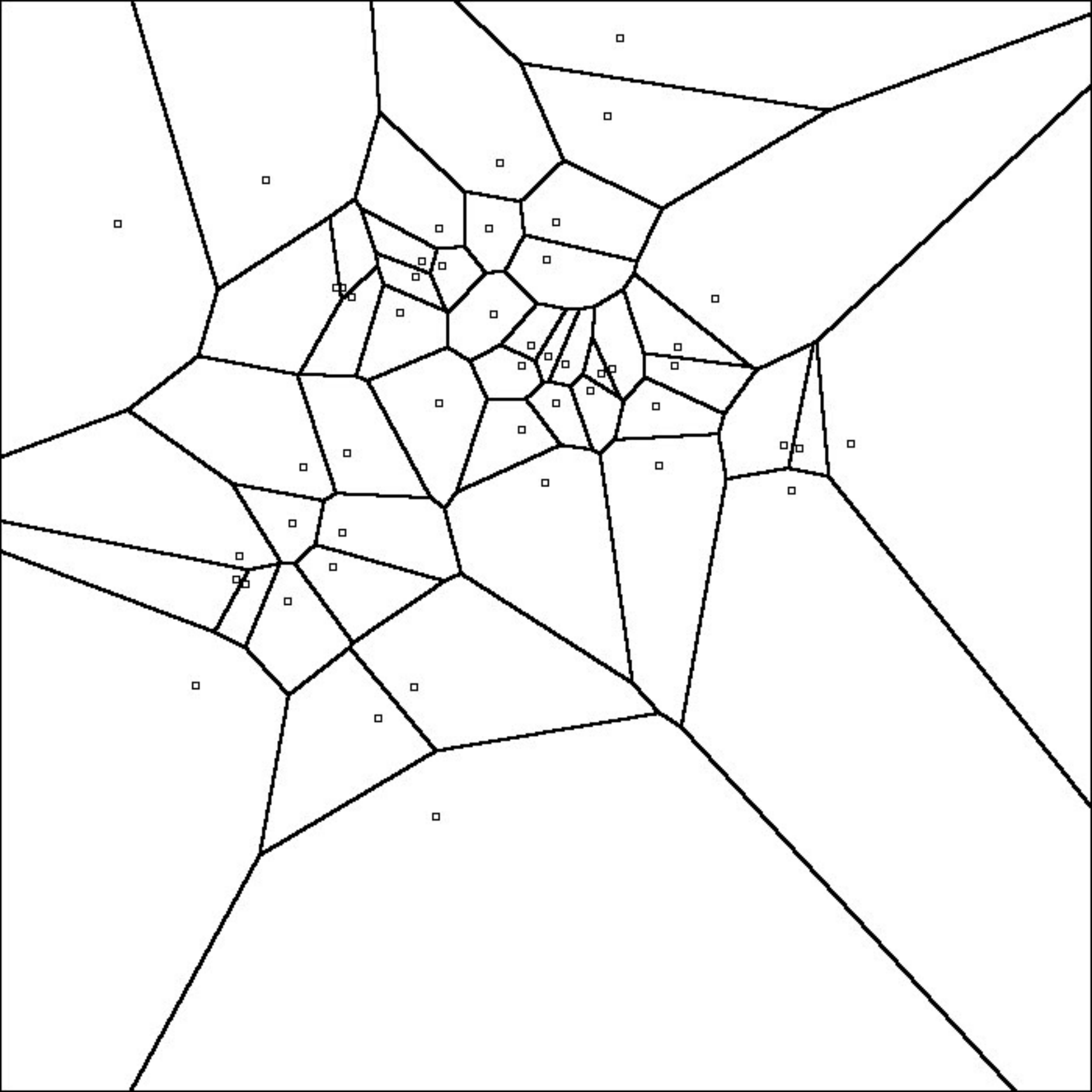












Shear Flow for Pedestrians

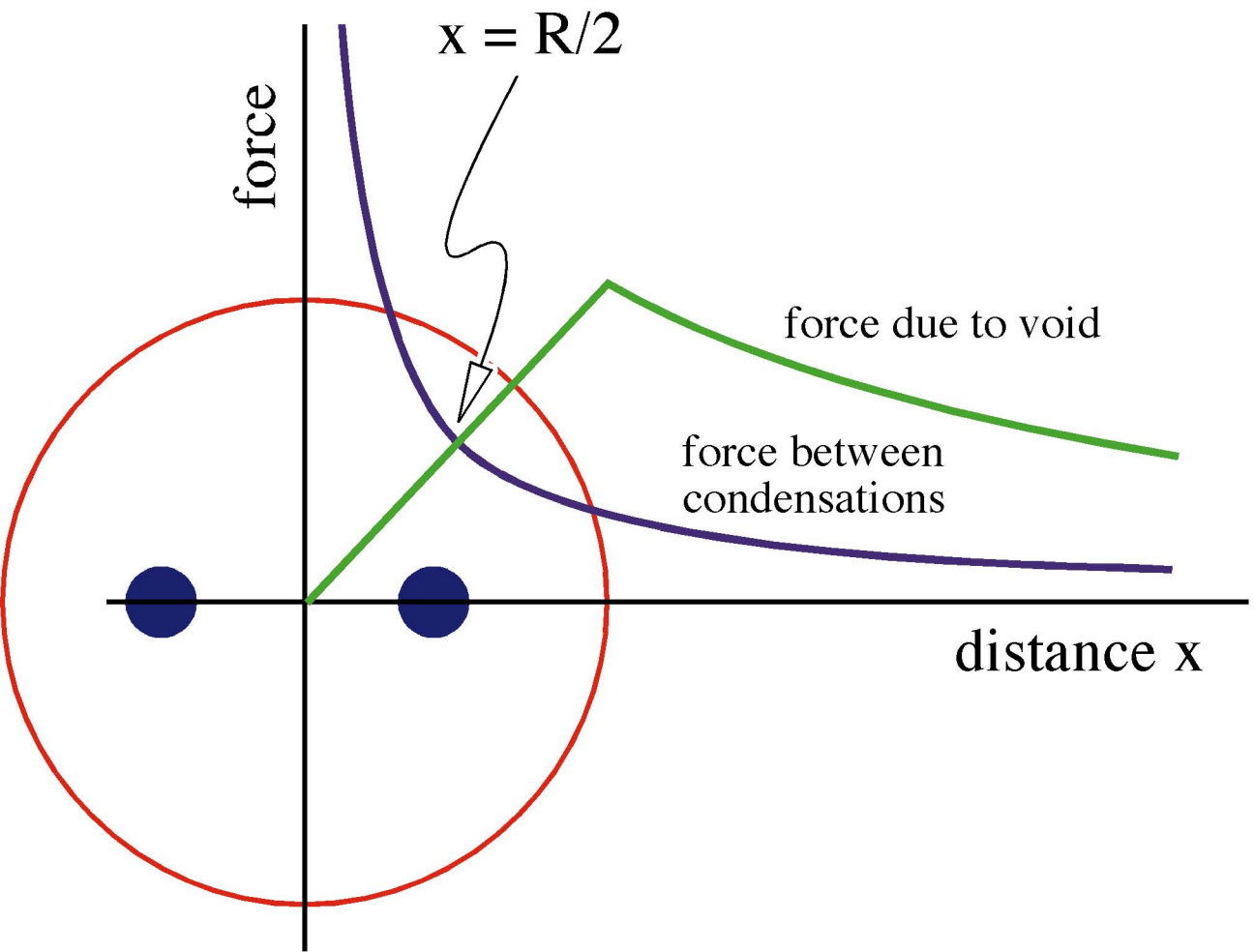
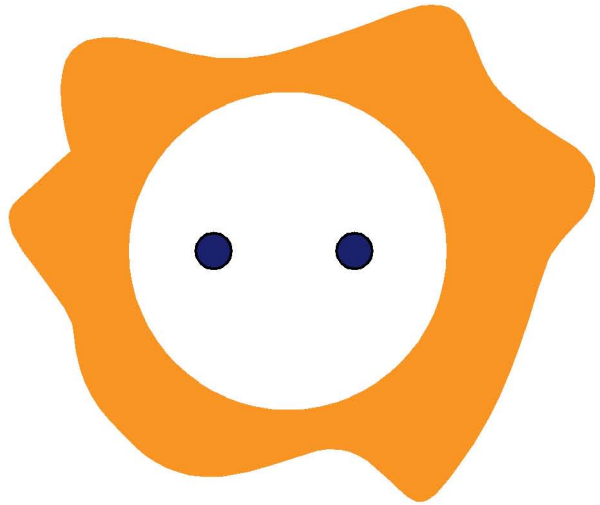
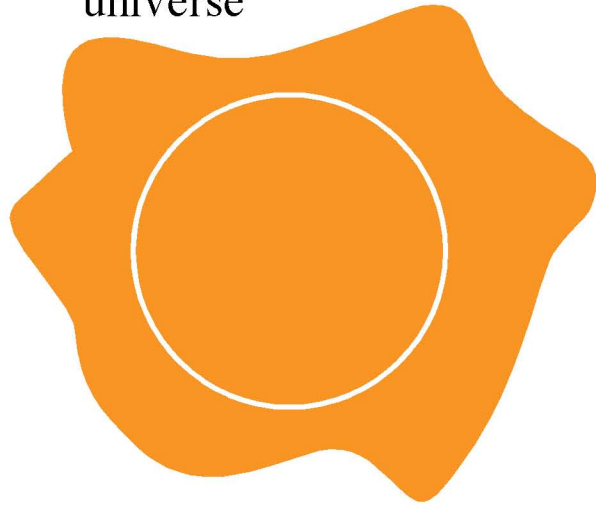
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Bubble Theorem

forces in a spherical void

unperturbed universe

inhomogeneous universe



The potential Φ near any point (x, y, z) of a self-gravitating medium can be written as

$$\Phi = \sum_{ijk} a_{ijk} x^i y^j z^k$$

Near a density maximum one has

$$\Phi = Ax^2 + By^2 + Cz^2 + \dots$$

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Neglecting terms of higher than second order, this is the potential of a homogeneous ellipsoid with axes (a, b, c) .

They evolve according to $(aX(t), bY(t), cZ(t))$, and the density ρ obeys

$$\rho(t) = \rho_0 / XYZ$$

If $a > b > c$, it follows from Poisson's Equation $\Delta\Phi = 4\pi G\rho$ and the equations of motion that

$$-\frac{1}{X} \frac{d^2 X}{dt^2} < -\frac{1}{Y} \frac{d^2 Y}{dt^2} < -\frac{1}{Z} \frac{d^2 Z}{dt^2}$$

Consequently, the axial ratios $a : b : c$ always increase with time. *Slight initial asphericities are amplified during the collapse.*

Now consider the evolution of the *low*-density regions. These are the progenitors of the voids.

A void is a region of *negative density* in a uniform background, so the voids expand as the overdense regions collapse. The sense of the deformation is reversed:

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*Slight asphericities decrease
as the voids become larger
("Bubble Theorem", Icke 1984).*

**Low-density fluctuations
form a packing of 'super-
Hubble bubbles':**

Voronoi Foam

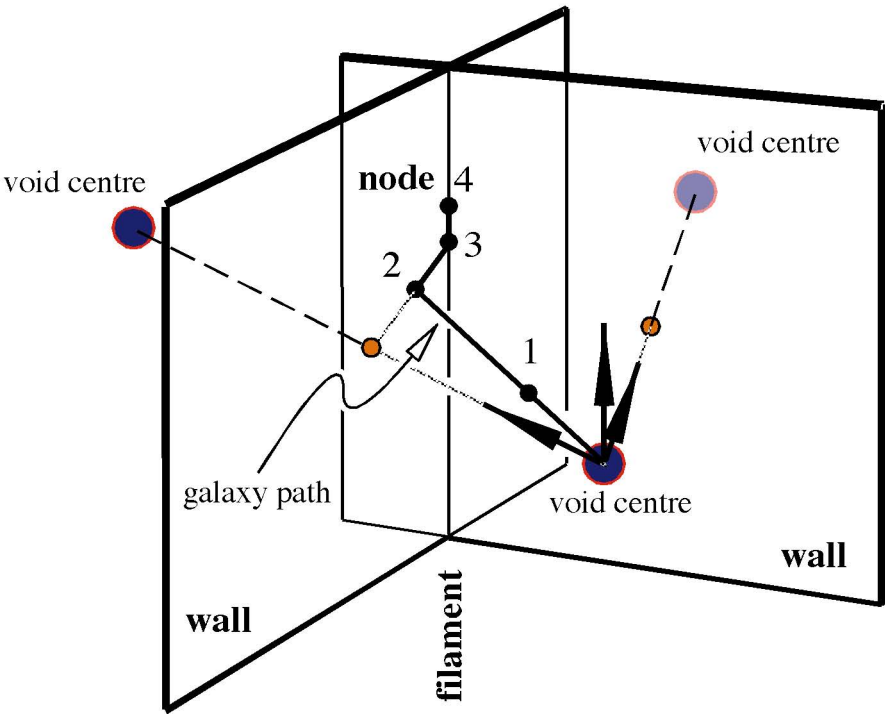
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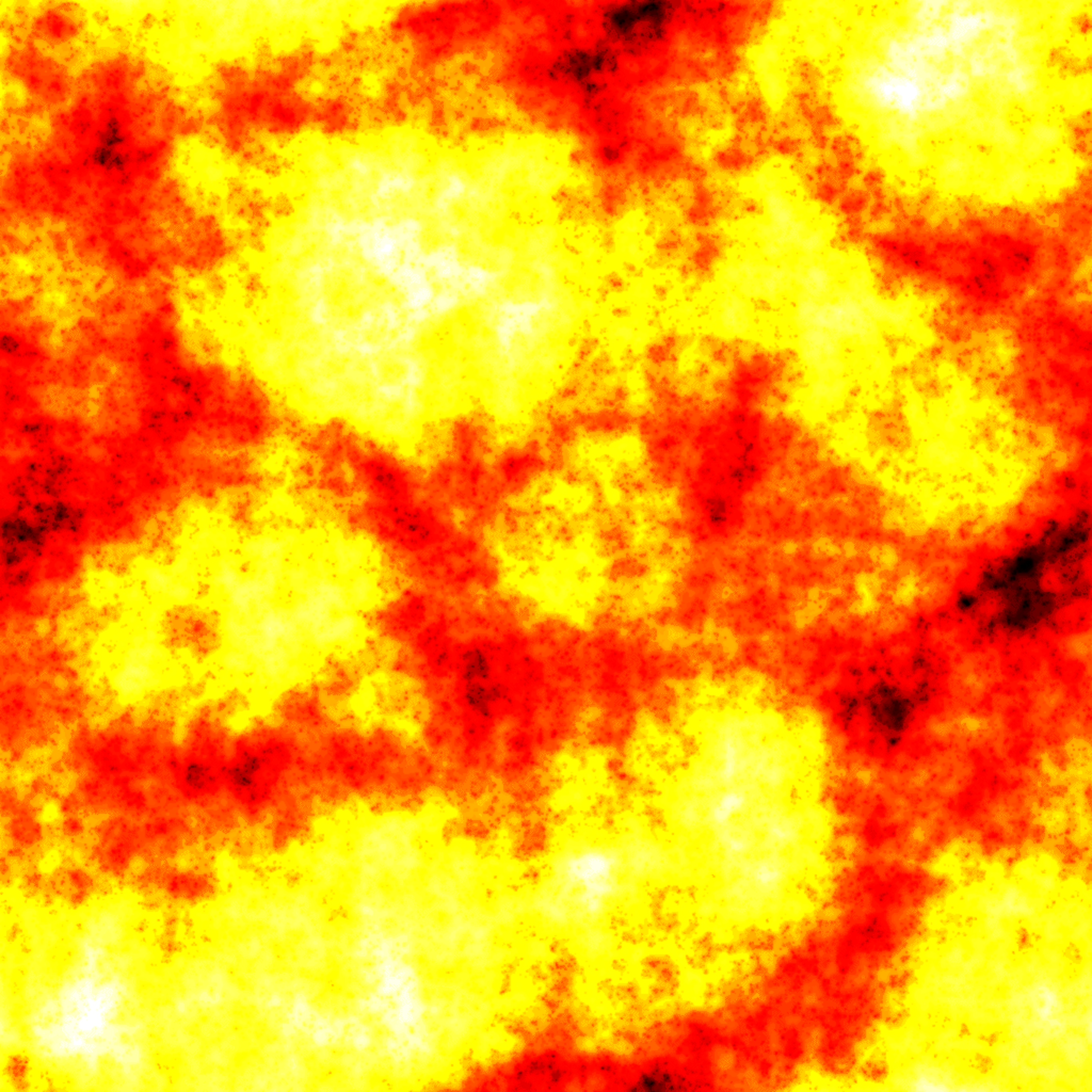
Evolution of Voronoi Features

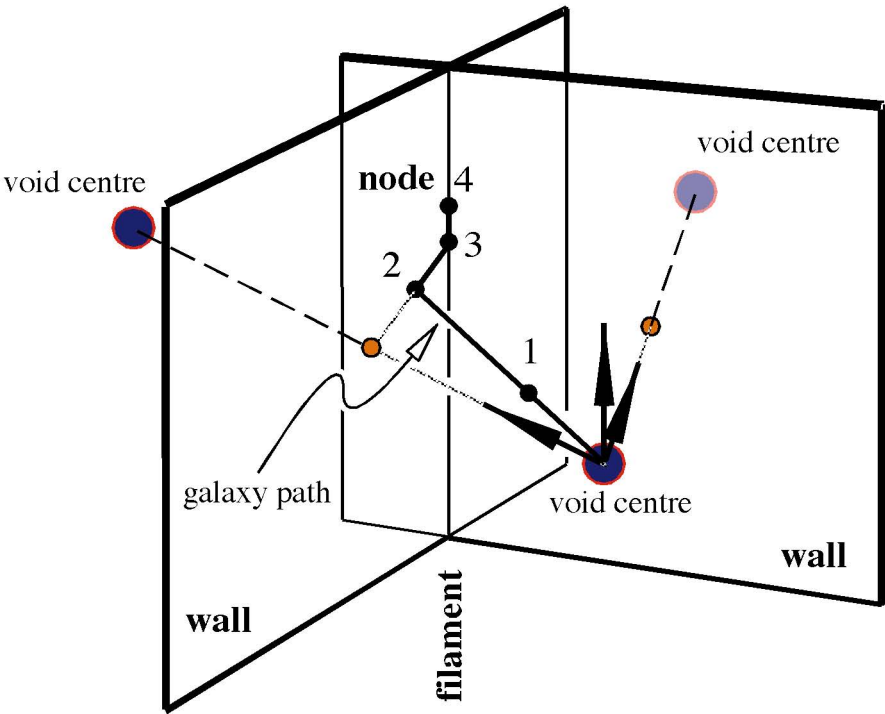
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Growth Laws for Voronoi Foam

Underdensities become more and more spherical when they expand, which they do a little faster than the Hubble expansion. Consequently, they form convex voids, from which the mass flows towards the zones separating the low-density regions. The asymptotic shape of such a mass distribution is a Voronoi foam.







Let m_v, m_w, m_f, m_n be the mass in voids, walls, filaments, and nodes, respectively. The first three features lose mass in a Voronoi “cascade”:

$$\text{void} \implies \text{wall} \implies \text{filament} \implies \text{node}$$

The mass loss term in N dimensions has the form

$$\text{mass change} = -NmH^* dt \quad (1.1)$$

where H^* is the excess Hubble expansion and dt is the time increment.

One may absorb H^* into the time, so that

$$\theta \propto t^{2/3} \quad (1.2)$$

The mass loss in N dimensions is then

$$dm = -Nm d\theta \quad (1.3)$$

The components parallel and perpendicular to the wall are

$$cH^* \cos \theta = aH^* \quad (1.4)$$

$$cH^* \sin \theta = bH^* \quad (1.5)$$

Thus, *the excess velocity in any Voronoi feature is simply found by multiplying H^* with the length along the feature.* This allows us to use the above formula for $N = 3, 2, 1,$ and 0 .

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The mass gain is found simply by reversing the sign of the loss term of the feature higher in the hierarchy. This gives the following equations:

$$\frac{dm_v}{d\theta} = -3m_v \quad (1.6)$$

$$\frac{dm_w}{d\theta} = 3m_v - 2m_w \quad (1.7)$$

$$\frac{dm_f}{d\theta} = 2m_w - m_f \quad (1.8)$$

$$\frac{dm_n}{d\theta} = m_f \quad (1.9)$$

These equations are simply solved by noting that the N -dimensional mass loss equation has a solution of the type $\psi \exp(-N\theta)$:

$$m_v = e^{-3\theta} \quad (1.10)$$

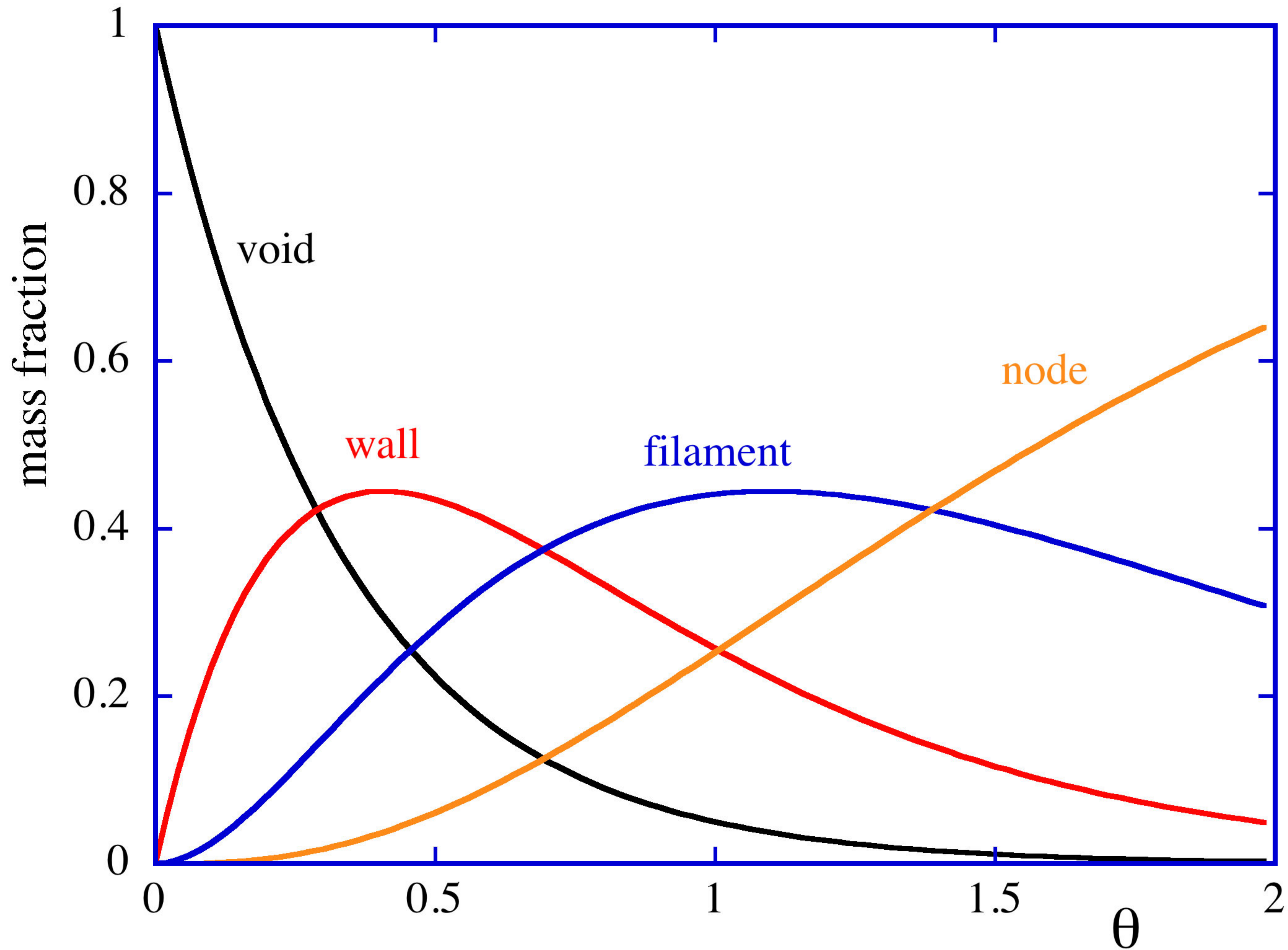
$$m_w = 3e^{-2\theta} (1 - e^{-\theta}) \quad (1.11)$$

$$m_f = 3e^{-\theta} (1 - e^{-\theta})^2 \quad (1.12)$$

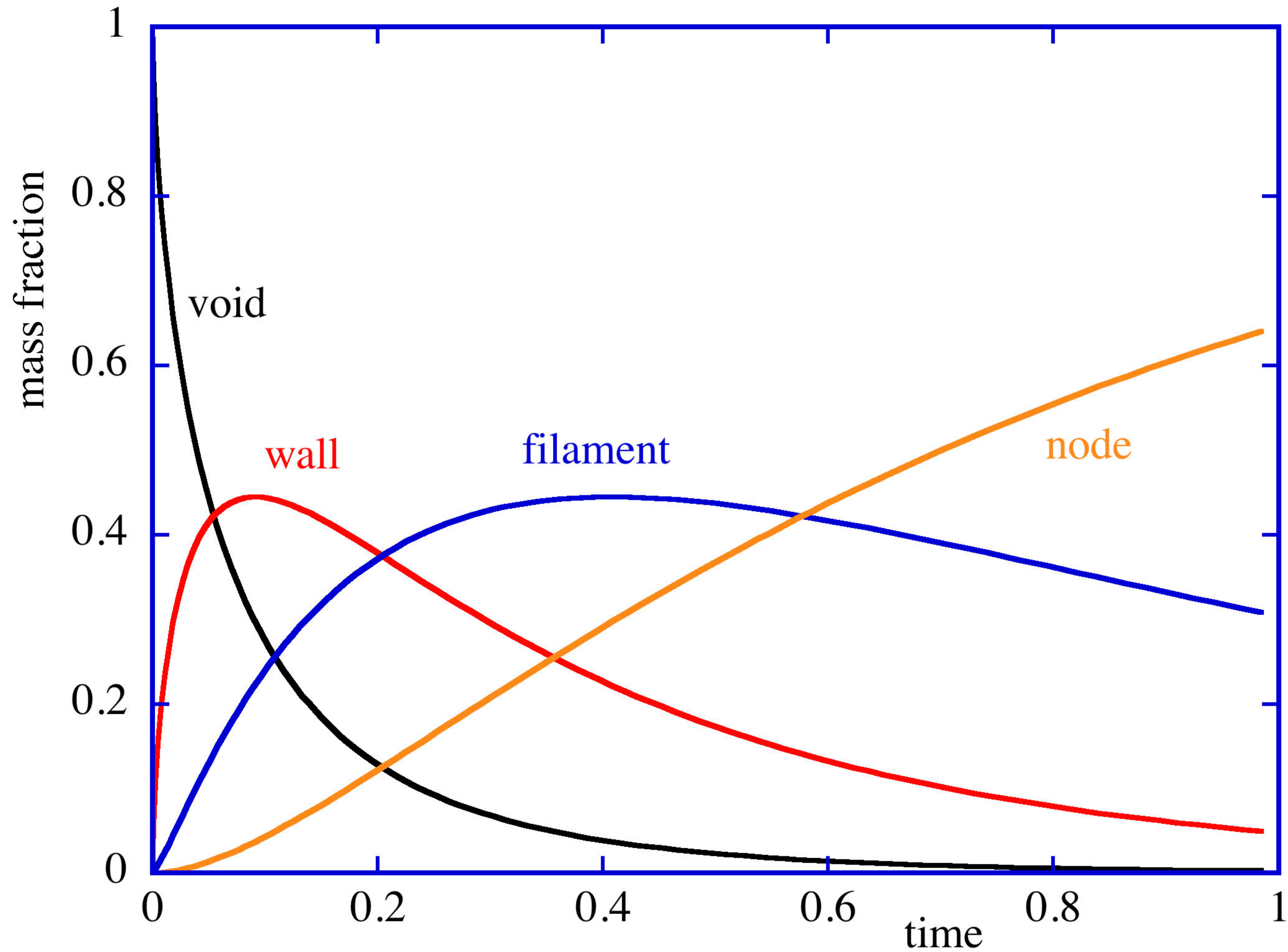
$$m_n = (1 - e^{-\theta})^3 \quad (1.13)$$

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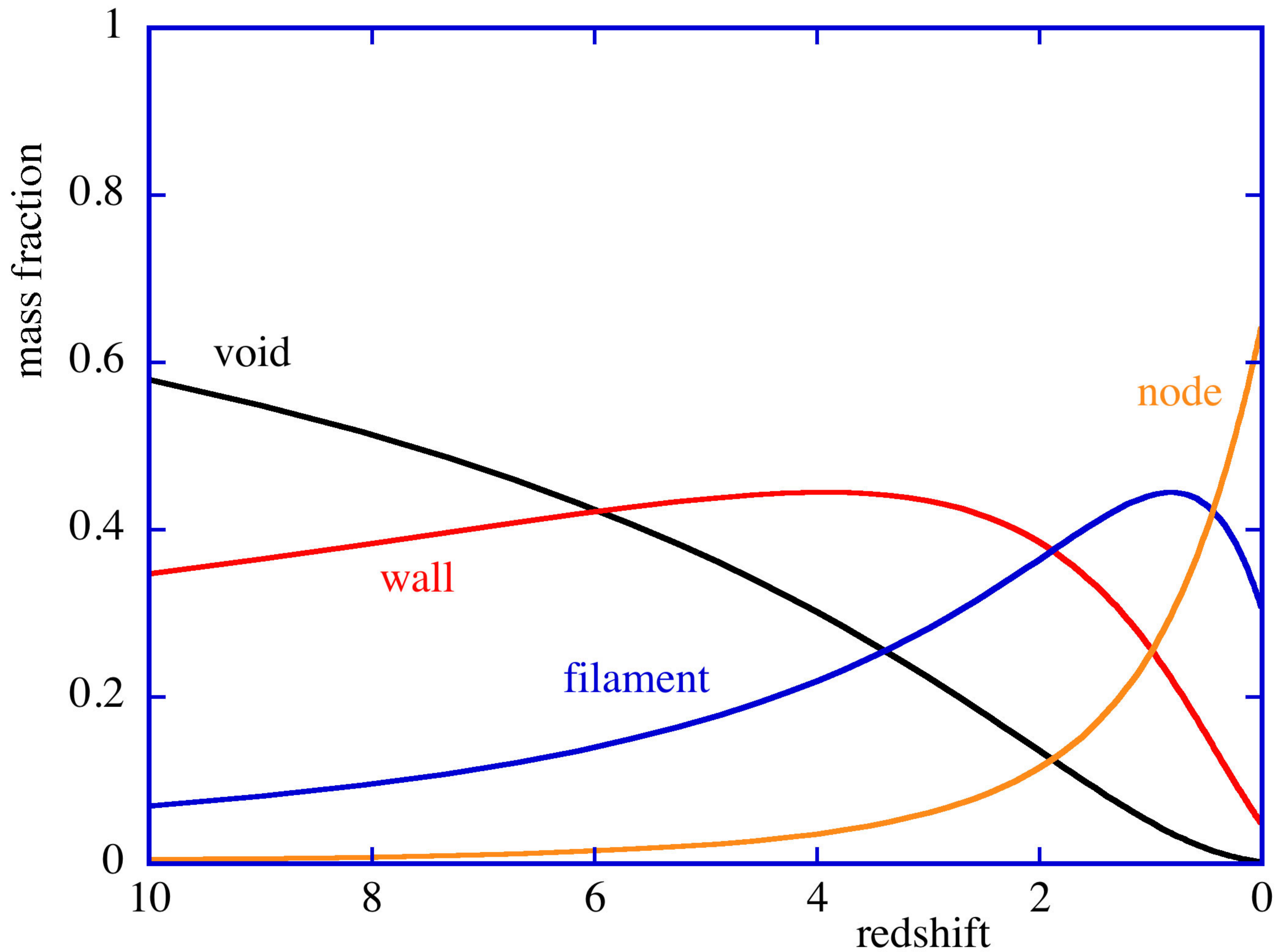
Mass of Voronoi Features



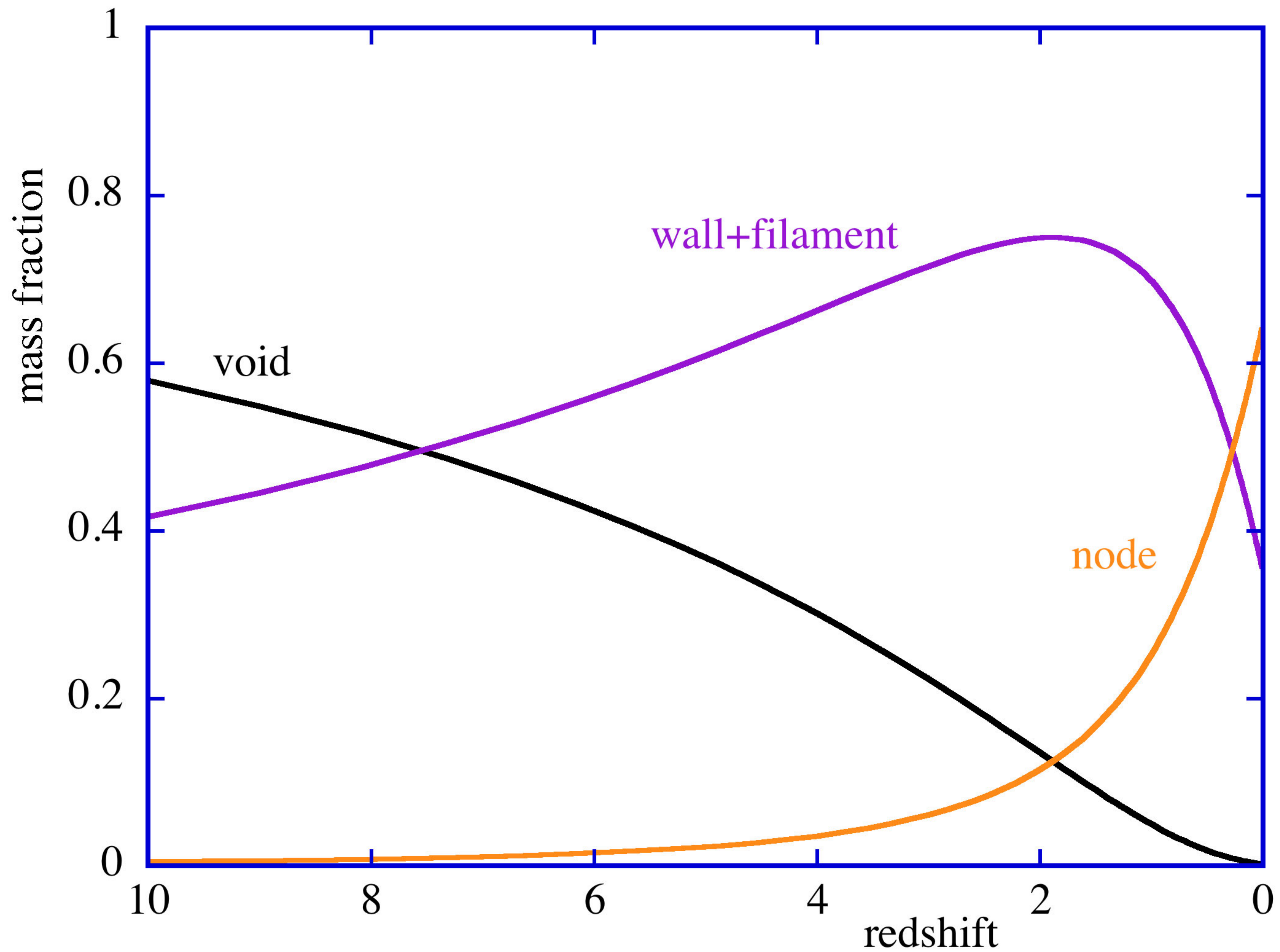
Mass of Voronoi Features



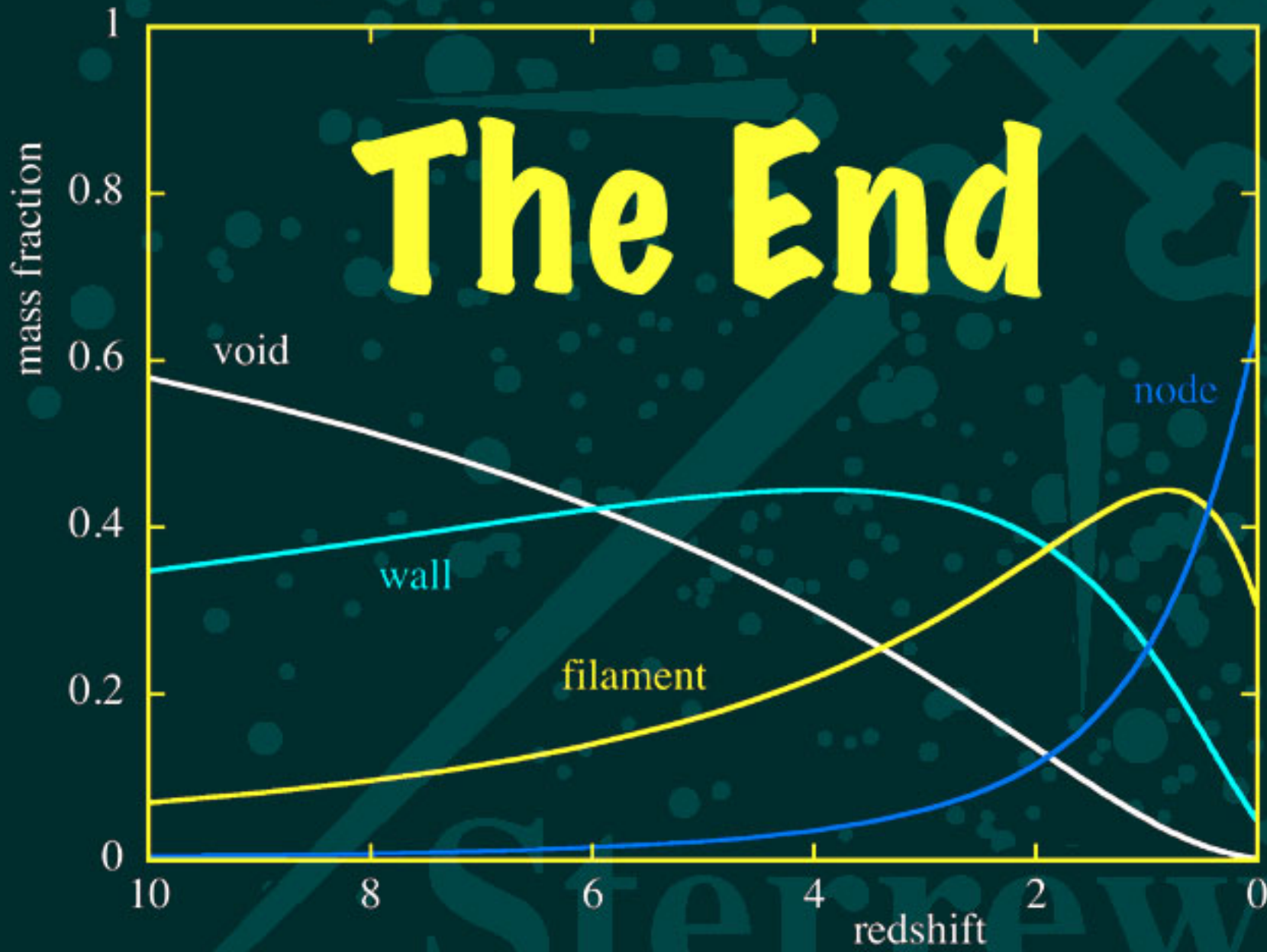
Mass of Voronoi Features



Mass of Voronoi Features



Mass of Voronoi Features



The End