

# SCHRIFTELIJK TENTAMEN

## COSMOLOGY

### 3<sup>rd</sup> term 2003/2004

#### Question 1.

The Friedman equations correspond to an adiabatically expanding medium. On the basis of this observation one can work out the temperature history of the Universe.

- Start by giving the full Friedmann equations (i.e. including cosmological constant, and including pressure). Describe and explain the various quantities in the equations.
- A Newtonian equivalent of these equations can be derived too. Describe how. Subsequently specify the fundamental differences between the Newtonian equations and the genuine relativistic one (there are 3 terms that should be discussed).
- From the Friedmann equations one can derive the equation for the evolution of  $\rho(t)$ , the change of the energy density of the expanding Universe. Derive this equation.
- Show that the derived equation in (c) implies an adiabatically expanding medium.
- For an adiabatic expanding medium we know that  $TV^{\gamma-1} = cst$ , with  $T$  the temperature of the medium and  $V$  the volume. Derive the temperature change of a uniform radiation field in an expanding Universe as a function of expansion factor  $a_{exp}$ .
- Given the fact that in an expanding Universe the frequency  $\nu$  of radiation is redshifted, so that the frequency of a photon of current frequency  $\nu_o$  has a frequency  $\nu(z) = \nu_o(1+z)$  at redshift  $z$ , show that when the radiation field is blackbody at a particular cosmic epoch, it will remain blackbody ! (hint: combine temperature and frequency evolution of photons).

#### Question 2.

To reconstruct the thermal history of the Early Universe, one needs to combine the knowledge of the temperature evolution of the Universe, that of the interaction rate  $\Gamma$  of the various relevant physical processes and the dynamical timescale of the Universe (...). Physical transitions occur when physical processes get out of equilibrium.

- What is the dynamical timescale of the Universe ? Write down the criterion for a physical process being in equilibrium and thus for when it runs out of equilibrium.
- What makes the Universe such a very special physical system ? In this, take into account that the temperature is continuously decreasing (affecting the reaction rates and timescales !) and about the number of photons available (a fundamental number of cosmology !!!!).
- Subsequently, describe in some detail the thermal history of the Early Universe. Mention at least five crucial transitions in the pre-recombination Universe, describe what happened at that transition, and approximately at what temperature/energy and redshift. Amongst those five, elaborate specifically on the epoch of primordial nucleosynthesis and that of the recombination/decoupling epoch.

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### Question 3.

- a) Einstein's general theory of relativity is a metric theory of gravity. Give the expression of the Einstein equation, and provide a physical explanation for the significance of this equation. Emphasize the essential difference with Newtonian gravity wrt. behaviour of spacetime in the general relativistic world !!!
- b) What does the Cosmological Principle say ? Needed is an enumeration of the (3) ingredients of the cosmological principle and a short description of what they mean. In addition, also discuss the meaning of Weyl's principle.
- c) What is the crucial significance of the cosmological principle in the context of relativistic cosmology. Describe the possible geometries allowed for a medium obeying the cosmological principle.
- d) Translate the above into a metric expression for the spacetime of an evolving universe obeying the cosmological principle and Weyl's principle:
  - give the name of this metric
  - write out the metric form,  $ds^2 = g_{\mu\nu}dx^\mu dx^\nu$ , for this spacetime
  - explain the meaning of each of the characteristic quantities. Focus in particular to those concerning the radius of curvature of spacetime, the related definition of the "scale factor"  $a_{exp}$  of the medium, and the factors expressing its curvature.
- e) In the metric we usually work in terms of the comoving distance  $r$ , while an observational cosmologist usually thinks in terms of the redshift  $z$ . How is the redshift  $z$  related to the expansion factor  $a_{exp}$  ? How can you translate comoving distance  $r$  in terms of the "redshift"  $z$  of an object (hint: I am asking for an integral expression, in which one also finds the Hubble expansion rate  $H(z) \equiv \dot{a}/a$ ).

### Question 4.

Standard FRW cosmology has three important problems or "fine-tuning" issues. One of them is the monopole problem, created approximately one within the horizon at GUT transition and which should therefore crowd the present-day Universe. They do not ! Inflationary cosmology, stating that the Universe underwent a rapid de Sitter expansion at a very early phase, tries to explain why not.

The two additional problems of standard cosmology are the "flatness" problem and the "horizon" problem. Here we will elaborate on these:

- a) For a FRW Universe:
  - define the meaning of the critical density  $\rho_{crit}$ .
  - give the expression for the critical density.
  - give an estimate of the value of the present-day critical density of the universe, in  $[g/cm^3]$  and in  $[M_\odot/Mpc^3]$ .
  - give the definition of  $\Omega(t)$ .

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What is a typical value for this radius of curvature (both for a curved space  $k \neq 0$  as well as for a flat space  $k = 0$ ).

- c) For the specific case of a curved matter-dominated Universe, work out the evolution of  $\Omega$  as a function of expansion factor  $a$  (or redshift  $z$ ) for various possible Universes. What will be the value of  $\Omega$  at  $a(t) = 0.001$  (recombination epoch) if the present  $\Omega_0 = 0.3$ . And at  $a(t) \approx 0.0001$  (around matter-radiation equivalence). In general, what do you expect  $\Omega$  to be for  $a(t) \rightarrow 0$ . Include a sketch of the  $\Omega$  evolution. Explain now what the flatness problem is.
- e) What is the definition for the (particle) horizon in a FRW Universe. Give the expression for the horizon distance  $d_{Hor}(t)$ . How does the horizon distance evolve with time  $t$  in a radiation-dominated Universe (assume  $\Omega_r = 1$ ), and in a matter-dominated Universe (assume  $\Omega_m = 1$ ). How can it be that the horizon seems to grow faster than the velocity of light? Subsequently, express these horizon distances in terms of the Hubble parameter  $H(t)$ .
- f) Work out the typical angular diameter  $\theta$  of the horizon at recombination ( $z_{dec} \approx 1000$ ).
- You may use the following expression for the angular diameter distance  $D_a$  in a matter-dominated Universe for an object at redshift  $z$ :

$$d_A = \frac{2c}{H_0 \Omega_0^2} \frac{1}{(1+z)^2} \left\{ \Omega_0 z + (\Omega_0 - 2) \sqrt{1 + \Omega_0 z} - 1 \right\} \approx \frac{2c}{H_0 \Omega_0} \frac{1}{z} \quad (\text{for } z \gg 1)$$

- if you are left with an equation including the Hubble parameter  $H(t)$ , use the expression of the evolution of  $H(z)$  as function of redshift  $z$  (from FRW equation), approximating  $(1+z) \approx z$  for  $z \gg 1$  and using

$$\Omega_0 H_0^2 = \frac{8\pi G}{3} \rho_0$$

- g) Given the impressive isotropy of the Cosmic Microwave Background, how does the derived value of  $d_{Hor}(z_{dec})$  imply the “horizon” problem?

SUCCESS !!!!

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