Abstract

This thesis presents the work I have done as part of a project to build a radio telescope to observe the CMB at 11 GHz. It covers some basic radio astronomy theory and the weather conditions in Groningen. After this it focusses on controlling the telescope and its measuring equipment through a Raspberry Pi. It will conclude with observations of the Sun and a satellite to verify the beam size. This was found to be $12.07 \pm 0.13^\circ$ using the Sun and $12.61 \pm 0.19^\circ$ using a satellite, in agreement with expectations.
## Contents

1 Introduction .................................................................................. 3

2 Radio Astronomy Theory .............................................................. 6
   2.1 Antenna Temperature .............................................................. 6
   2.2 Noise and System Temperature ............................................... 7
       2.2.1 Johnson-Nyquist Noise ................................................... 7
       2.2.2 Propagation of Noise ....................................................... 8
   2.3 Sensitivity and Integration Time .............................................. 9
   2.4 Applying the Theory to the Kapteyn Radio Telescope ............... 10

3 Weather Conditions in Groningen ................................................. 15
   3.1 Data Acquisition ................................................................... 15
   3.2 Atmospheric Temperature .................................................... 15
   3.3 Cloud Coverage ................................................................. 18
   3.4 Situation in Groningen .......................................................... 23

4 Controlling the Telescope ............................................................ 24
   4.1 Computer Control ............................................................... 24
       4.1.1 Arduino vs Raspberry Pi ................................................. 24
   4.2 Controlling the Motor .......................................................... 25
       4.2.1 Stepper Motors .............................................................. 25
       4.2.2 Communicating with the Motor ....................................... 27
   4.3 Measurement Hardware and Control ..................................... 29
       4.3.1 Power Meter ................................................................. 30
       4.3.2 Temperature Monitor ...................................................... 30
       4.3.3 Communicating with the Devices .................................... 32
   4.4 Operating the Telescope ........................................................ 33
       4.4.1 Connecting to the Telescope .......................................... 35
       4.4.2 Observing with the Telescope ....................................... 35

5 Measurements of the Sun ............................................................ 37
   5.1 Solar Emission ..................................................................... 37
   5.2 Observing the Sun ............................................................... 38
   5.3 Verifying the Beam Pattern ................................................... 40
   5.4 Measuring the Brightness Temperature of the Sun ................... 43

6 Measurements of other Sources .................................................. 44
   6.1 Galactic Synchrotron Radiation .............................................. 44
   6.2 Astrophysical Radio Sources ............................................... 44
   6.3 Non-Astrophysical Sources ................................................... 45
7 Future Improvements and Conclusion 47
   7.1 Code Efficiency and Speed .............................................. 47
   7.2 Reading Results from the Power Meter .............................. 47
   7.3 Motor Power and Control .............................................. 47
   7.4 More Powerful Hardware .............................................. 48
   7.5 Conclusion ................................................................. 48

8 Acknowledgements 49
   A Average Temperatures for June ........................................... 52
   B Average Cloud Coverage for June ....................................... 57
   C Telescope Server ............................................................. 62
   D GPIB Communication ......................................................... 65
   E Stepper Motor Control ....................................................... 68
   F DNS Configuration ......................................................... 71
1 Introduction

A Short History of Radio Astronomy

Electromagnetic radiation can exist at any frequency. We are able to see signals coming from high energy gamma or X-ray radiation all the way down to low energy radiation from radio waves. Over time radio astronomy has come a long way and underwent a lot of changes.

The first astronomical radio observation was made by Karl Jansky. He used a rotating array of antennas that was sensitive in the horizontal direction. By rotating the array he was able to measure the intensity of the signal as well as the horizontal component of the direction that the signal came from. Nothing could be said about the vertical direction however. Jansky made his measurements at a wavelength of 14.6 meters \cite{Jansky1982}. This translates to a frequency of 20.5 MHz. He concluded that the radiation came from a region of the Milky Way \cite{Jansky1982}. His observation was made in 1933.


Interested by Jansky’s discovery, Grote Reber also started working in the field of radio astronomy. Reber built a much more advanced telescope to make his observations with. He built a 9.5 meter parabolic dish telescope \cite{Reber1944}. Reber’s telescope had three receivers operating at 3300 MHz, 900 MHz and 160 MHz. With the 160 MHz receiver he detected radiation coming from the center of the galaxy \cite{Reber1944}.

Nowadays we are still building new radio telescopes to further improve our observations by collecting more signal or by having even higher angular resolution for more detailed images. Both signal strength and angular resolution depend on the size of the telescope. Inventions range from large parabolic reflector dishes like the 100 meter Effelsberg Radio Telescope or the

\footnote{Using }{c = 299792458 \text{ m s}^{-1}}
300 meter Arecibo telescope to large arrays of antennas like the Low Frequency Array (LOFAR) or the upcoming Square Kilometer Array (SKA).

**The Cosmic Microwave Background**

After the Big Bang the Universe was hot and dense. Matter and photons were strongly coupled and light could not escape. Minutes after the Big Bang atomic nuclei formed, but it was still too hot for the electrons to recombine with the nuclei. The Universe was a hot plasma of charged particles and photons and would remain so for approximately 380000 years. After this the temperature dropped sufficiently for the electrons to recombine with the nuclei and atoms to form: the Epoch of Recombination. Before this, photons could not escape because of Thomson scattering, but now that the free electrons were captured the photons were less likely to be scattered and they could escape. It is these first photons that we see when we look at the Cosmic Microwave Background (CMB).

**Antennas**

When charged particles are accelerated they emit electro magnetic (EM) radiation. This is the way radio signals are broadcasted. A current in the emitting antenna causes the electrons to be accelerated and emit radiation with a power proportional to the acceleration squared. The reverse can also happen. EM radiation can accelerate particles. This is what happens in a receiving antenna. The EM wave causes electrons in the antenna to start moving and create a current. This current can then be recorded and is a measure of the power received by the antenna. With high angular resolution a source could be resolved and the radiation at various positions could be measured to create e.g. an image of the source in radio emission.
The Project

The ultimate goal of this bachelor project is to measure the temperature of the CMB. To do this we built a computer controlled horn antenna. Measurements will involve making sweeps of the sky to measure the power as a function of elevation. The project was done by four students, each of them focusing on some aspect of the telescope. The design of the horn was done by Bram Lap [Lap, 2015]. Next Maik Zandvliet [Zandvliet, 2015] was in charge of constructing the mount and mechanics of the telescope and for designing the back end where the electronics are. I am responsible for the control of the telescope. I wrote the software necessary to both move the telescope and provide a means to collect the data and I carried out observations of the Sun. Finally Willeke Mulder [Mulder, 2015] focussed on calibration of the telescope to look into the stability of the system and receiver temperature, which is crucial to know if we are going to measure the CMB for which she also did the observations.

Summary

This thesis will present my contribution to the project. The first two chapters will cover material that is not directly related to controlling the telescope, since during this time it was still under construction. The material presented here is general information that is useful to know. In Chapter 2, I will first look into some basic theory behind radio astronomy to see what kind of signal we are actually receiving when doing observations and how the error in this signal comes to be. Chapter 3 will next look into the weather conditions in Groningen and the effect observing conditions have on the measurements. The fourth chapter will be a more practical section about the electronics and the software for operating the telescope. Chapters 5 and 6 will present observations of the Sun and discuss the possibility of detecting other sources respectively at 11 GHz. Lastly Chapter 7 will have a discussion about possible improvements for the control of the telescope and about the results of the observations.
2 Radio Astronomy Theory

This chapter will give an explanation of what signal the telescope will receive and how it will propagate through the system. The signal that comes in will be changed both due to atmospheric effects and noise effects within the receiver. For the best detection we want to minimize these effects as much as possible.

2.1 Antenna Temperature

Measuring a received power is quickly done, however the data is unusable until it is calibrated to a standard scale. In radio astronomy the signals are often calibrated against a temperature. From a body with a known physical temperature we can measure the power it emits and use it to set a scale for other measurements. The power received can now be related to a temperature, the so-called antenna temperature. The antenna temperature is thus not the physical temperature of the antenna.

The received power comes from all over the sky including other sources and even ground emission. To get the total power we need to account for all of these contributions by taking into account the power pattern of the antenna and the observed brightness of the sky. The power pattern is a measure of how sensitive the antenna is to power from a certain direction. To get the total power that is received we should thus convolve the observed brightness with the power pattern:

\[ P(\nu) = \frac{A_e}{2} \int_{4\pi} B(\theta, \phi) P_n(\theta, \phi) \, d\Omega. \]

Here \( P(\nu) \) is the received spectral power, \( A_e \) the effective area of the telescope, \( B(\theta, \phi) \) the Planck function, \( P_n(\theta, \phi) \) the normalized power pattern and \( \Omega \) the solid angle. The integration is over the whole sky. As a temperature this becomes

\[ T_A = \frac{1}{\Omega_A} \int_{4\pi} T_b(\theta, \phi) P_n(\theta, \phi) \, d\Omega \text{ K}, \]

which follows from the fact that \( P(\nu) = kT_b \) (see next section) and the Rayleigh-Jeans approximation of the Planck function \( B(\nu) \). Here \( T_A \) is now the antenna temperature, \( \Omega_A \) the beam solid angle, \( T_b \) the brightness temperature in some direction \((\theta, \phi)\) and \( P_n \) the normalized power pattern in that direction. Finally, using the definition for the beam solid angle we get the following expression for the antenna temperature [Wilson et al., 2013]:

\[ T_A = \frac{\int_{4\pi} T_b(\theta, \phi) P_n(\theta, \phi) \, d\Omega}{\int_{4\pi} P_n(\theta, \phi) \, d\Omega}. \] (1)

The antenna temperature can be interpreted as the weighted average of the power received from all over the sky. The sky brightness distribution is contained in \( T_b \), the brightness temperature. The power received from a certain position on the sky is now multiplied by the value of the
normalized power pattern $P_n$ at that position. The normalized power pattern is 1 at its maximum and decreases for the sidelobes. Sidelobes are positions outside of the main beam where the response of the antenna is non-zero. The maximum is usually located in the center of the beam. Therefore most of the power will come from the main beam. It is not possible to know from what direction what power was received unless the observer knows there is a source at a certain position. Therefore ideally we would want all of the power to come in through the main beam so that we know exactly where the power is coming from. Hence reducing the sidelobes as much as possible is a major part of designing a radio telescope to ensure the antenna temperature incorporates only the signal of the source we want to measure and nothing else. For measuring the CMB we will consider three components for the antenna temperature: the CMB itself, the atmosphere and possible ground pickup. This results in a total antenna temperature of

$$T_A = T_{CMB} e^{-\tau} + T_{atm}(1 - e^{-\tau}) + T_{ground} \text{ [Mulder, 2015]}. \quad (2)$$

2.2 Noise and System Temperature

We consider all measured sources as noise sources. The telescope will thus measure a total noise power $P_\nu$ with associated noise temperature $T_{sys}$. Not all of the noise however comes from external sources. Besides the power received by the antenna, there is also noise power generated in the receiver itself with noise temperature $T_{Rx}$. The total noise temperature will then be the receiver, $T_{Rx}$, temperature plus the antenna temperature,

$$T_{sys} = T_A + T_{Rx} \quad (3)$$

Here $T_A$ can include contributions from any noise source one can think of.

$$T_{sys} = T_{CMB} + T_{atm} + T_{source} + T_{ground} + T_{galaxy} + T_{sun} + \cdots + T_{Rx}.$$ 

As we will see later on, for our observations we will only consider the CMB and the atmosphere contribution of Eqn. 2. To eliminate the ground contribution the horn was designed with minimal sidelobes. Even if we were to pick up some ground, its contribution is negligibly small in our case [Mulder, 2015]. In Chapter 6 we will see that the galaxy is also negligible. To relate this system temperature to the measured power it is insightful to look at the origin of the receiver noise, in particular that of the amplifiers.

2.2.1 Johnson-Nyquist Noise

The noise of the amplifiers in the back end comes from random thermal motions of electrons inside the components. The random nature of this motion means that the average displacement is zero, but the mean square displacement is not. Moving charges will create currents so the motion of the electrons causes a small current $I$ with $\langle I \rangle = 0$, but $\langle I^2 \rangle \neq 0$ to start flowing. Higher temperatures will make the electrons move faster and thus the noise goes up. Cooling the amplifiers properly can therefore reduce the produced noise. Because of its nature, this noise is called thermal noise or Johnson-Nyquist noise after the two people that played an important role in discovering and explaining this type of noise.
To find a relation between the measured power and the corresponding noise temperature let's look at a resistor. Theoretically we can replace the noisy resistor with a noise source followed by a noiseless resistor. The noise power can be determined by loading a resistor $R$ with another resistor $R_L$. The power that $R$ puts into $R_L$ is then given by

$$P_\nu = \langle I^2 \rangle R = 4kT \frac{RR_L}{(R+R_L)^2}$$ \[van Schooneveld, 1990\].

This power is maximum when the loads are matched, i.e. $R = R_L$. Now all factors except for $kT$ cancel out. We have now found a relation between power and temperature:

$$P_\nu = kT_e.$$ \hspace{1cm} (4)

Here $T_e$ is the effective temperature of the component. This can be equal to its physical temperature, but generally is not. It is important to note that this relation only holds for the spectral power $P_\nu$. The power that the telescope will measure will be changed by both the bandwidth and the gain. To convert this into a temperature we will need to correct for these before using the above relation.

### 2.2.2 Propagation of Noise

Our telescope uses two amplifiers \[Zandvliet, 2015\]. Adding more amplifiers naturally adds more noise which can get hard account for as the number of amplifiers increases. However, we only have to consider the noise coming from the first amplifier. This has to do with where in the chain the amplification happens. Consider a chain of $n$ amplifiers with effective temperatures $T_{e,n}$ and gains $G_n$. The output of an amplifier is $T_{out} = G(T_{in} + T_e)$ \[van Schooneveld, 1990\]. The total noise temperature for $n$ amplifiers is given by Friis formula \[Wilson et al., 2013\],

$$T_e = T_{e,1} + \frac{T_{e,2}}{G_1} + \frac{T_{e,3}}{G_1G_2} + \cdots + \frac{T_{e,n}}{G_1G_2\cdots G_{n-1}}.$$ \hspace{1cm} (5)

For a system consisting of two amplifiers this gives an effective temperature of

$$T_{e,sys} = T_{e,1} + \frac{T_{e,2}}{G_1}.$$  

Physically this means the following. The noise added by later amplifiers is with respect to an already amplified signal. This implies that the true noise added by this amplifier is its noise temperature divided by the gain of the previous amplifiers. From Eqn. (5) it becomes clear that this rapidly makes the noise added by later amplifiers negligible. If amplification in the first stage is high enough we can even consider the noise of the second amplifier to be negligibly small, because of this only the noise of the first amplifier is important.
2.3 Sensitivity and Integration Time

In the previous section we saw that a radio telescope measures a total temperature $T_{sys}$ at its output given by Eqn. [3] This temperature has a RMS uncertainty, $\sigma_T$, associated with it given by

$$\sigma_T = \frac{T_{sys}}{\sqrt{\Delta \nu t}}$$

(6)

The quantity $\sigma_T$ is called the sensitivity of the telescope, $\Delta \nu$ is the receiver bandwidth and $t$ the integration time, i.e. how long signals are collected for a measurement. The smallest signal that can be detected is a signal that is $\sigma_T$ higher than these RMS fluctuations. Sources of interest are often weaker than the noise sources, so good sensitivity is a major plus for an observation. Looking at Eqn. [6] we see that we can increase the sensitivity in three ways.

- **Decrease System Temperature**: Creating a system that has a low noise level will lower $T_{sys}$ and make the system more precise.

- **Increased Bandwidth**: With a large bandwidth the telescope will receive power from a wider spectrum and therefore sensitivity is increased.

- **Longer Integration Time**: By observing the source for a longer time more signal is collected, increasing precision.

To get an understanding of how Eqn. [6] comes to be we look at what $\sigma_T$ actually is. The system temperature $T_{sys}$ can be seen as a noise temperature with all signals to be coming from noise sources. The random nature of this signal results that we are in fact measuring the variance of a random signal with noise. Measuring the variance $\sigma^2$ of random samples has an uncertainty of $\sqrt{2}\sigma$. This gives the first part of the equation,

$$\sigma_T = \sqrt{2}T_{sys}$$

For the denominator we need the Shannon sampling theorem and the Central Limit Theorem. The sampling theorem provides the minimum frequency one needs to sample a signal at in order to completely reconstruct it. It states that if a signal has a frequency $f$ then one can reconstruct this signal completely by sampling it at a rate $2f$ [Jerri 1977]. Suppose an incoming signal with a bandwidth $\Delta \nu$. The theorem then implies this signal needs to be sampled at $2\Delta \nu$. After a time $t$ we will have collected a number of samples $N = 2\Delta \nu t$. Finally the Central Limit Theorem tells us that the uncertainty in $T_{sys}$ drops as $\sqrt{N}$, with $N$ being the number of measurements. This will give

$$\sigma_T = \frac{\sqrt{2}T_{sys}}{\sqrt{2\Delta \nu t}}$$

thus yielding Eqn. [6]
2.4 Applying the Theory to the Kapteyn Radio Telescope

The horn is a Pickett-Potter horn design about 30 cm in length. The receiver is built to receive signals at 11 GHz with a bandwidth of 1.05 GHz, going from 10.45 to 11.5 GHz. This means we are looking at signals with a wavelength of 2.7 cm \[\text{Lap, 2015; Zandvliet, 2015}\]. This section will apply the theory discussed above to the telescope to get a feeling for the quantities we will be dealing with.

Antenna Temperature and CMB Brightness

The main purpose of this telescope was to be able to measure the temperature of the CMB. The CMB is a near perfect black body radiating at a temperature of $T_{\text{CMB}} = 2.72584 \pm 0.00057$ K \[Fixsen, 2009\]. For designing the back end it is important to know what order of power we are receiving. Black body radiation is characterized by the Planck spectrum,

$$B_\nu(T) = \frac{2h\nu^3}{c^2} \frac{1}{\exp(h\nu/kT) - 1},$$

(7)

where $h$ is Planck’s constant; $\nu$ is the observing frequency; $c$ is the speed of light; $k$ is the Boltzmann constant and $T$ is the temperature of the black body. Figure 3 shows the spectrum of the CMB. The peak of the signal lies at a frequency of $\nu = 160$ GHz. At this frequency the CMB has a brightness of $B_\nu = 3.88 \cdot 10^{-16}$ W Hz$^{-1}$ m$^{-2}$ sr$^{-1}$. The available hardware however will limit us to a range of frequencies between 4 and 11 GHz. In Fig. 3 we see that this range is at the faint end of the CMB spectrum. Therefore, we chose the highest possible frequency of 11 GHz. This gives a brightness of $B_\nu = 9.29 \cdot 10^{-18}$ W Hz$^{-1}$ m$^{-2}$ sr$^{-1}$. The atmosphere will both attenuate the CMB signal and add to the signal \[Mulder, 2015\]. This will result in an antenna temperature of

$$T_A = T_{\text{CMB}} e^{-\tau} + T_{\text{atm}} (1 - e^{-\tau}).$$

(8)

Assuming $T_{\text{CMB}} = 2.72584$ K, $T_{\text{atm}} = 283$ K and $\tau = 0.05$ (see \[Mulder, 2015\]) when looking at zenith we can make an estimate of the antenna temperature and the received power. Using Eqn. 8 the antenna temperature is $T_A \approx 16.4$ K. The incoming power follows from multiplying Eqn. 4 with the bandwidth. The effective temperature now is a combination of the antenna temperature and the receiver temperature.

$$P = k(T_A + T_{\text{Rx}})\Delta \nu$$

If we assume a receiver temperature fo $T_{\text{Rx}} = 200$ K this gives an incoming power of $P = 3.13 \cdot 10^{-12}$ W or $P = -85.03$ dBm using $k = 1.3807 \cdot 10^{-23}$ J K$^{-1}$ and $\Delta \nu = 1.05$ GHz. The measuring range of the power meter extends from $-70$ dBm to $+20$ dBm \[Agilent Technologies, 2013a\]. Our system uses two amplifiers of around 30 dB each \[Zandvliet, 2015\]. If we assume the system gain comes only from the amplifiers and is 60 dB this results in a signal of approximately $-25.03$ dBm.

If the optical depth decreases to $\tau = 0.01$ the antenna temperature drops to 5.5 K. Using the same value for the receiver temperature and gain the power becomes $P = -25.26$ dBm. The
Figure 3: Top: the brightness $B_\nu$ of the CMB as given by the Planck function for $T_{CMB} = 2.72584$ K. Bottom: a zoom of the available frequency range, showing the expected $B_\nu \propto \nu^2$ property of the Rayleigh-Jeans law.

The final signal will be somewhat different because of effects of the other components, cables, the horn and changes in both atmospheric and receiver temperature, but it will be around these values.

**Noise and System Temperature**

We assumed before that the noise produced by the second amplifier is negligible. Here we show this is not the case. The noise temperature of the second amplifier, ALMA2, has an average value of 127.35 K [Zandvliet 2015]. Using Eqn. 5 its final contribution would then be 4.25 K. The noise temperature of the first amplifier, MITEQ, is 110.90 K. Since the CMB has a temperature of 2.7 K, a noise of 4 K on 110 K is in this case not negligible. This is all taken in however in the measurement of the receiver temperature. The noise temperature of the back end will thus lie close to the noise of the two amplifiers combined. The noise temperature of the back end is measured by [Zandvliet 2015].
Sensitivity and Integration Time

The noise temperature of the system allows us to make a prediction of our sensitivity. If we want to be able to measure the temperature of the CMB to 10% accuracy, the upper limit on the sensitivity will be as follows. Thus, the limit is set on $0.2 \text{ K}$. Eqn 6 depends linearly on $T_{sys}$. As the receiver temperature fluctuates, so will the sensitivity. We measured the receiver temperatures ranging from around 180 K to around 200 K. Looking at Fig. 4 and Fig. 5 we see that for a sensitivity of 0.2 K the integration time will be of the order of milliseconds. Even in bad conditions with high optical depths this is still the case. The low integration times tell us that our measurements will not be limited by the system noise. Instead the limiting factors will be how well we can determine the other quantities like the atmospheric temperature, the opacity and the receiver temperature.
Figure 4: The sensitivity as function of integration time for a receiver temperature of 180 K (left) and 200 K (right).
Figure 5: The integration time as a function of Zenith Angle $z$ for sensitivity $s_r = 0.2$ and different $T_r$. The integration time for different optical depths for a receiver temperature of 180 K (left) and 200 K (right).
3 Weather Conditions in Groningen

Despite being able to reach our requirements even in relatively bad circumstances, we still want the observing conditions to be as good as possible. In Groningen we have two main concerns: the atmospheric temperature and the cloudiness. In this chapter an estimate of what is approximately the best observing moment in Groningen will be made. It is not insightful to pinpoint an exact date, so we will limit the investigation to find an approximate best time of the year and time of day to observe.

3.1 Data Acquisition

To be able to get a good overview of the weather in Groningen we need data that spans over multiple years. To do this we use data from the KNMI, the institute that monitors the weather throughout the Netherlands. These data are publicly available on the website\(^2\). It is also possible to request the data through a script, which is the way it was done here. With the available data it is possible to go all the way back to the previous century. In this case we will use the data from 2007 till 2015. Since there is no weather station in Groningen we will use the data from the nearby station in Eelde [KNMI].

The data from the KNMI is available in either daily measurements or hourly measurements. We used the hourly measurements to have the greatest flexibility. The KNMI tracks all sorts of data such as windspeed, sunshine, precipitation, cloud coverage and temperatures. For us only the cloud coverage and temperature data are important. When reading data from their database it should be taken into account that all the temperatures are presented in units of \(0.1\) °C and hence needs to be converted to units of \(1\) °C first. Cloud coverage is expressed in special units of oktas or octants. This unit indicates how many eights of the sky are covered in clouds. It ranges from 0 to 8, with 0 octants meaning clear skies and 8 octants meaning a sky completely covered in clouds. In some situations the sky may not be visible at all, such as with dense fog. In this case the cloud coverage gets a 9.

3.2 Atmospheric Temperature

Equation 8 for the antenna temperature shows that the atmosphere increases the signal with unwanted noise and does so inverse exponentially. With \(e^{-\tau} = 0.95\) for \(\tau = 0.05\) the contribution from the atmosphere may seem small, but because the atmospheric temperature is high relative to the CMB and other sources, the \(T_{atm}(1 - e^{-\tau})\) term still contributes significantly, typically ranging from 3-10 K [Mulder 2015]. We therefore aim for an atmospheric temperature that is as low as possible. In Fig. 6 we have plotted the minimum, average and maximum temperature for every day of the year. The distribution looks roughly the same every year. We can see that the highest and lowest temperatures are during the summer and winter months respectively. For a more clear view we calculated the minimum, average and maximum temper-

\(^2\)http://www.knmi.nl/klimatologie/uurgegevens/selectie.cgi (hourly) or http://www.knmi.nl/klimatologie/daggegevens/download.html (daily).
Figure 6: The minimum, average and maximum temperature as measured at Eelde from 2007 till 2015. The month labels do not exactly separate the data points by month, but give a good indication.
Figure 7: Minimum (blue crosses), average (yellow dots) and maximum (red crosses) temperatures for each month separated by year. The errorbars give the standard deviation of the average.
Table 1: Minimum, average and maximum temperatures in June per year in °C.

<table>
<thead>
<tr>
<th></th>
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<td>31.3</td>
<td>25.6</td>
<td>28.4</td>
<td>26.6</td>
</tr>
</tbody>
</table>

3.3 Cloud Coverage

Even more important than the atmospheric temperature is the cloud coverage. When there is a cloud in the field of view of the telescope the opacity of that area can increase dramatically. For larger optical depth the factor $e^{-\tau}$ becomes increasingly smaller. This reduces the signal from the CMB and increases that of the atmosphere. For comparison, the contribution of the atmosphere is 2.97 K for $\tau = 0.01$ while it is 14.53 K for $\tau = 0.05$, assuming $T_{atm} = 298$ K. Besides increasing the opacity, clouds will also make it non-uniform over the area of the beam. This adds another complication in accounting for the contribution from the atmosphere.

In Fig. 9 we see the cloud coverage as measured by the station in Eelde. We see that the data is spread out through the plot filling the entire 0 to 9 scale. It does however seem to tend to the higher cloud coverage values on average.

To get a better view of the data we took the average per month. The results of this are shown in Fig. 10. This confirms the earlier statement of higher overall cloud coverage. Even in the months that were good in terms of atmospheric temperature, the sky is still covered in clouds for the most part. This is not completely unexpected. Even if there would be days of little to no cloud coverage these are averaged out by days of high cloud coverage. It would be quite rare for an entire month to have a low cloud coverage.

Finally we split the data by part of the day like for the temperature to see if it changes significantly. Figure 11 demonstrates this, but the results are inconclusive. Appendix B has the plots for the other years but they do not give usable information either.
Figure 8: Temperatures of June 2014 separated by part of the day. From top to bottom: night, morning, afternoon and evening.
Figure 9: Cloud coverage in octants for all days of the year. The month labels are not exact but give a good indication. Red: maximum; Yellow: average; Blue: minimum.
Figure 10: The average cloud coverage per month. The errorbars give the standard deviation of the average. Red: maximum; Yellow: average; Blue: minimum.
Figure 11: Cloud coverage of June 2014 separated by part of the day. Red: maximum; Yellow: average; Blue: minimum.
3.4 Situation in Groningen

The two most important factors of the weather are atmospheric temperature and cloud coverage. It turns out we have limited control over the first one. We can make a rough estimate of when the temperatures will be lowest and hence best for observing. This is in the winter months January, February and March. Further, breaking it down and looking at the time of day, we also saw that night time has the lowest temperatures. The best observing date and time would therefore be sometime in the winter during the night (i.e. 00:00 - 06:00). We will be observing in June, but the night time will still have the lowest temperatures. Cloud coverage on the other hand is beyond our control. Due to its changing nature we cannot make an estimate of a best observing moment. The best moment for this would need to be decided by the observer by checking the weather conditions.

Although this analysis has given some insight in the weather conditions, it still remains difficult to find reliable correlations between observing time and cloud coverage. For the best observing time it comes down to checking the weather forecast or the weather outside directly to see if conditions are and will stay suitable.
4 Controlling the Telescope

4.1 Computer Control

Our measurements will involve moving the telescope over a strip in the sky in elevation. To make a reliable measurement we decided to mount the horn on a frame and control the telescope with a computer. The telescope moves with a motor and power is read digitally from a powermeter. This gives us a couple of advantages. First of all the mount will ensure that we start and end in the same strip of the sky. This way we only have to worry about the received power as function of elevation, easing the data reduction later on. Secondly we will control the speed factor. Using a motor controlled by a computer ensures that each measurement will be taken in the same amount of time. This way data collection is consistent for every measurement. Lastly it makes determining the received power at a certain angle easier and more reliable. A full computer would be too large to fit into our build so we opted to go for something smaller. This gives us the choice between an Arduino and a Raspberry Pi.

4.1.1 Arduino vs Raspberry Pi

The telescope will be rotated by a stepper motor. Controlling the motor is a simple task and can be done by a simple system, but maybe expansions are desired in the future needing a more advanced system. To make a decision we compared the two devices with each other. Table 2 lists the specifications of the Arduino Uno and Raspberry Pi B+ respectively. It immediately shows that a Raspberry Pi (from here on "Pi") is much more powerful than an Arduino. This is because there is a key difference in how they operate. The Arduino has a microcontroller. It has no OS, it just runs its code when told to do so. The Pi however is a full computer complete with OS.

Programming the Arduino would (in principle) require programs in C and the programming would be indirect. One would write the code and upload it to the microcontroller. It would then be able to do the one thing the program was written for. Programming the Pi can be done in Python on the machine itself. This has multiple advantages. For one, Python is a flexible and easy to use programming language, making developing with it much more convenient. Being able to program on the device itself is also a plus for quickly making small changes later on.

<table>
<thead>
<tr>
<th>Arduino Uno</th>
<th>Raspberry Pi B+</th>
</tr>
</thead>
<tbody>
<tr>
<td>Processor</td>
<td>Processor</td>
</tr>
<tr>
<td>ATmega328</td>
<td>Broadcom BCM2835 SoC</td>
</tr>
<tr>
<td>Voltage</td>
<td>Voltage</td>
</tr>
<tr>
<td>5V/7-12V</td>
<td>5V</td>
</tr>
<tr>
<td>CPU Speed</td>
<td>CPU Speed</td>
</tr>
<tr>
<td>16 MHz</td>
<td>700 MHz</td>
</tr>
<tr>
<td>SRAM</td>
<td>SDRAM</td>
</tr>
<tr>
<td>2 kB</td>
<td>512 MB</td>
</tr>
</tbody>
</table>

Table 2: Arduino Uno and Raspberry Pi B+ specifications.

---

Finally the Pi is the most future proof should upgrades, e.g. more sensors, be desired in the future. It is for these reasons that the telescope will be controlled by a Raspberry Pi B+.

### 4.2 Controlling the Motor

#### 4.2.1 Stepper Motors

To control the elevation of the horn we used a stepper motor. A stepper motor divides a full rotation into a number of steps of equal size. Applying an electrical pulse moves the motor one step. This allows for the precise control of the position of the motor. Since it moves in discrete steps we can determine its position simply by counting the number of pulses we sent to it. This eliminates the need of external tools for measuring the position. There will be an error associated with the angle moved due to the possibility of missed steps. A great property of stepper motors is that this error is non-cumulative. One step will therefore have the same error as a million steps [Ericsson]. We will not account for this error in the position. The reason for this is that we tested the motor multiple times by letting it move to various angles and then return to the park position many times. It had not missed any steps during the tests, so the error is negligible.

A stepper motor consists of two parts: a stator and a rotor. Figure 12 shows this basic structure. The stator is the outer ring consisting of electromagnets with teeth. Each pair of opposing poles of electromagnets is called a phase. Most common is the 2-phase stepper motor. Adding more phases will give a higher resolution, i.e. a smaller angle per step. Recently, 5-phase stepper motors have achieved higher resolutions and more steps per revolution. A 2-phase motor has four poles per phase, while a 5-phase motor, for example, has two poles per phase.

In Fig. 12 we see a 2-phase and a 5-phase motor with the phases labeled. It might seem that the 2-phase motor actually has four phases, but this is because it are actually two motors together. For 2-phase motors the coils usually have a center tap, which allows for easy reversing of the magnetic poles without switching the current direction. This is called a unipolar motor.

The rotor is the inner disk. This disk has magnetized teeth, all coming in north-south pairs. Once a phase is turned on some of these teeth on the rotor will be out of alignment with those on the stator. Magnetic forces will now pull the rotor in alignment with the stator: the step. To move the motor the phases are turned on and off in a certain sequence [Orientalmotor 2015]. These sequences are determined by so-called drive methods. There are three main drive methods: wave drive or 1-phase full step, 2-phase full step and half step [Ericsson].

- **Wave Drive:** Wave drive has one phase on at a time. It is the simplest way to drive a stepper motor and the coil activation scheme would be $A \rightarrow B \rightarrow \bar{A} \rightarrow \bar{B}$. The downside is that for a unipolar motor this way does not give maximum torque because only a quarter of the coils is used at any time.

- **2-phase Full Step:** For maximum torque the 2-phase full step scheme is better. This way two phases are on at all times, using half of the available coils. The coil activation scheme for this drive method is $AB \rightarrow B\bar{A} \rightarrow \bar{A}\bar{B} \rightarrow \bar{B}A$. 

- **Half Step:** Half stepping combines wave drive and 2-phase driving. This way the motor can rotate half the angle it would with the other two methods. The coil activation scheme for half stepping is $A \rightarrow AB \rightarrow \bar{A} \rightarrow \bar{A}B \rightarrow \bar{B} \rightarrow \bar{BA}$

![Stepper Motor Schematic](http://www.orientalmotor.com/technology/articles/2phase-v-5phase.html)

**Table 3: General properties of the Astrosyn Y129 stepper motor.**

<table>
<thead>
<tr>
<th>Stepper Motor Properties</th>
<th></th>
</tr>
</thead>
<tbody>
<tr>
<td>Type</td>
<td>Unipolar</td>
</tr>
<tr>
<td>Phases</td>
<td>2</td>
</tr>
<tr>
<td>Step Angle</td>
<td>1.8°</td>
</tr>
<tr>
<td>Holding Torque</td>
<td>9 N cm</td>
</tr>
<tr>
<td>Operating Voltage</td>
<td>12 V</td>
</tr>
<tr>
<td>Steps per Rev.</td>
<td>200</td>
</tr>
</tbody>
</table>

The motor driving the telescope is an Astrosyn Y129 2-phase stepper motor. Table 3 lists some general information about this motor. The most important is that it takes 200 steps for the motor to make one revolution. We operate the motor in half step mode. With each step it will move by 1.8°. To mount the telescope the motor is connected to a larger rotatable disk. Our supervisor told us that in this setup it takes 72 revolutions of the stepper motor for the disk to make one. We verified this by marking the disk and then sending $200 \cdot 72$ pulses to the motor. This means with each step our telescope will move $0.025°$ or 90 arcsec on the sky. In between the pulses there is a small delay of 2 ms or 1.3 ms depending on the speed setting. If the pulses
arrive too fast the motor will not have time to respond and it will hang. It takes the horn 80 seconds to do a measurement with a 1 degree step and 32 seconds with a 13 degree step both at the fastest mode.

### 4.2.2 Communicating with the Motor

To control the motor we use the *General-Purpose Input/Output* (GPIO) pins on the Raspberry Pi. The stepper motor controller we use has a 15-pin connector, but we only need three pins: ground, step and direction. These pins are connected to a ground pin and two GPIO pins on the Pi respectively. We cannot however directly connect the motor controller to the Pi, as the GPIO pins run at 3.3 V while the controller needs 5 V. If the 5 V would flow into these pins the Pi could possibly get damaged. To prevent this from happening we made a special connector with Zener diodes connecting the step and direction pins to ground. A Zener diode allows current to flow backwards once a certain breakdown voltage is exceeded. The connector has two 3.3 V Zener diodes. If the voltage over the Zener diode now exceeds 3.3 V the current will flow towards ground. The schematic for the socket is shown in Fig. 13. These three wires are then connected to the following pins on the Pi: 29, 31 and 30 for direction, step and ground respectively. The pins on the Pi can be set either high (3.3 V) or low (0 V). For movement, we set the direction pin either high or low to move forwards or backwards. Next the step pin is continuously switched between high and low. Every time the step pin is powered the motor will move one step. After testing with a prototype of loose wires, a custom cable was made which can be seen in Fig. 14.

![Figure 13: Schematic for the connector connecting the motor controller to the Pi. The connector is a 15 pin female D-Sub connector.](image)
Figure 14: The cable connecting the motor controller box to the Raspberry Pi. Left: the inside of the connector showing the wires and the Zener diodes. Right: the complete cable with aluminium case and connector for the GPIO pins.
4.3 Measurement Hardware and Control

To measure the power received by our telescope we use an Agilent E4412A power sensor connected to a Agilent E4418B power meter. These devices are shown in Fig. 15. To measure temperatures we use a Lakeshore 218S Temperature Monitor for reading out the sensor, see Fig. 16. To save the measurements we have to set up a connection between the Pi and the devices so we can read out values and store them.

Figure 15: The equipment used for measuring the power. Left: power meter. Right: power sensor. Image credit: http://www.sglabs.it/ and http://www.us-instrument.com/

Figure 16: The Lakeshore 218S Temperature Monitor used for reading out the sensor. Image credit: http://www.lakeshore.com/products/cryogenic-temperature-monitors/model-218/pages/Overview.aspx
Table 4: The integration time in seconds as for all combinations of speed and number of averages per measurement.

<table>
<thead>
<tr>
<th>Speed</th>
<th>Readings Averaged</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>1</td>
</tr>
<tr>
<td>20</td>
<td>0.05</td>
</tr>
<tr>
<td>40</td>
<td>0.025</td>
</tr>
<tr>
<td>200</td>
<td>0.005</td>
</tr>
</tbody>
</table>

4.3.1 Power Meter

The used power meter has a wide detection range capable of detecting signals with frequencies between 100 kHz up to 110 GHz and powers between -70 dBm and +44 dBm [Agilent Technologies, 2013a]. The power sensor falls nicely within this range, being able to measure frequencies between 10 MHz and 18 GHz and powers between -70 dBm and +20 dBm. The power meter is a key component in the sensitivity of our system as given by Eqn. 6. The power meter is set to take a certain number of readings and average them together to produce a measurement. The measurement speed (amount of readings per taken per second) and the number of readings to average are separately adjustable. These quantities now determine what our integration time will be according to Eqn. 9:

\[
\tau = \frac{\text{readings averaged}}{\text{measurement speed}}
\]  

The available options are limited however. Table 4 lists the possible integration times as function of the measurement speed and the number of readings averaged for a measurement. Averaging more samples will give lower noise levels improving the precision. For an absolute measurement it is ±0.02 dB when measuring in dBm or 0.5% when measuring in W [Agilent Technologies, 2013a]. This error can be brought down by taking multiple measurements, i.e. changing the number of readings that are averaged. The error will go down as the square root of the number of readings.

4.3.2 Temperature Monitor

To measure the atmospheric temperature and the temperature of the hotload we use a Pt1000 sensor. It operates between -50°C and +150°C with tolerance class 2B. This means the sensor has an accuracy given by

\[
\sigma = \pm (0.6 + 0.005|t|) \quad [\text{Heraeus}, 2013].
\]  

Note that \(t\) should be in °C. Outside this range the sensor still functions but the uncertainty becomes larger and it might damage the sensor due to material stresses. To connect the sensor to the temperature monitor another cable was made which can be seen in Fig. 17.

The temperature is determined by measuring the resistance of the sensor. This introduces a problem: wires have resistance. If we were to construct a circuit as in the left of Fig. 18...
we would not only measure the resistance of the sensor, but that of the wires as well. This is a direct error in the measurement. The solution to this problem is to separate the current and voltage measurements. The current will be the same since it is a series circuit. The voltage however drops by Ohm’s law $V = IR$. To eliminate the wire resistance one should only measure the voltage drop over the sensor as in the right picture of Fig. 18. This is known as four-wire sensing or Kelvin sensing. The wires to the voltmeter only carry a tiny current. The voltage drop caused by them can therefore be neglected. The measured voltage is now only the voltage drop caused by the sensor.

Figure 17: The Pt1000 temperature sensors with the cable to connect it to the Lakeshore temperature monitor.

Figure 18: Measuring the resistance of the sensor (left) vs. measuring the voltage drop over the sensor (right).
4.3.3 Communicating with the Devices

To communicate with the power meter and temperature monitor we had two options: RS232 or GPIB. The Pi does not have an RS232 interface, but there are USB converters so that would be no problem. Our supervisors advised against RS232 for two main reasons. First of all RS232 would be a difficult protocol to work with. One would have to account for e.g. start/stop bits, the number of data bits sent or the number of bits transmitted per second. A second problem would arise when using an RS232 to USB converter. The name by which the computer addresses the USB device depends on the order in which all devices are plugged in. The names are not persistent, so if multiple devices will be used we cannot be sure that the next time devices are plugged they will have the same names again.

GPIB or IEEE-488 was developed in 1974 by HP for devices that can automatically test other devices or perform measurements automatically. It was an attempt to standardize communication between computers and instruments. With the General Purpose Interface Bus (GPIB) all instruments would have the same connector. Later with the introduction of the Standard Commands for Programmable Instruments (SCPI) all devices obtained a common basic programming command set along with a set of guidelines for manufacturers should they wish to include new commands. Using GPIB multiple measuring devices can easily be used and is as simple as daisy chaining them together. This is because each device has its own address for communication. Two limitations should be considered however. Each device can have an address between 0 and 30, but a driver can only drive 14 devices at a time. Second, the cable length should not exceed 20 m in total or 2 m between devices for maximum data transfer rates. [ICS Electronics] [2009]

For this project it turned out that GPIB connection would be the easiest to work with. The Pi does not have a GPIB connector either, so to communicate with the equipment we use a GPIB-Ethernet controller shown in Fig. 19. This acts as a bridge converting signals received over ethernet to GPIB and vice versa. To read out results one connects to the controller via an...
ethernet cable. Before commands can be sent we need to set up a connection with the controller. This is done by setting up a TCP connection at port 1234 at the controllers IP address [Prologix, 2013]. In Python this is easily done using sockets in the `socket` module. The details can be seen in Appendix C. The equipment can now be read out by sending the appropriate commands over the socket. When talking to a device two things should be accounted for. Commands should always be terminated by the newline character `\n`. This indicates you are done sending the command and that it can start processing it. The second is the way sockets work. In the Python documentation they warn the user that when sending or receiving a message, you may not send or receive the entire message in one call. It is the users responsibility to make sure everything is sent or received. Sending the entire message can be done with the socket module itself by using the `socket.sendall` method. Commands are then sent to the controller which processes them and queries the power meter for a response. However there is no such method for receiving. To make sure we receive everything the power meter sends us we keep querying it for more data until it returns `None` indicating there is nothing more to send. The code governing communications with the GPIB controller and the powermeter can be found in `GPIBComm.py` in Appendix D. The fact that we can use ethernet influenced the decision on how to operate the telescope. This is discussed in the next section.

### 4.4 Operating the Telescope

There are many question as to how the telescope can be operated. Should the measurements be automatic to eliminate user error? Should the operator have input? If so, to what extent? How will the user interface look like? What preliminary knowledge does the user need to have? These are all questions to consider when thinking about the interface to the telescope. The telescope will possibly be used for the astronomy bachelor course *Radio Astronomy* in the third year. Thus it will mainly be students operating the telescope.

The Pi has four USB ports and an HDMI port so the first idea was to simply have a monitor, keyboard and mouse. During the frame design however (see [Zandvliet, 2015]) it became clear this would not be ideal. It has too many downsides. Having many peripherals would mean one would need to access the Pi frequently for connecting and disconnecting and we would need a way to handle all the cable work. It also meant carrying many things alongside the telescope. This would compromise portability a lot. Furthermore the whole setup would become just too bulky. All these situations make the setup less user friendly, so we discarded this idea.

Thinking about the portability led us to consider the operator controlling the telescope from his or her own laptop. The only problem was then how the communication between the Pi and the laptop would work. The only ethernet port was taken by the GPIB controller and there are no other connections. Connecting everything by ethernet does solve the problem of not knowing what device has what name, since everything gets an IP address. We thought of setting up a small LAN network depicted in Fig. 20. The network would consist of the following components:

- Raspberry Pi
- GPIB-Ethernet controller
Figure 20: The LAN network connecting all devices in the telescope with each other.
• User laptop
• Network switch or router

The switch or router will provide the ethernet ports to connect everything together. The Pi will take care of communication between the user’s laptop and the system.

4.4.1 Connecting to the Telescope

To operate the telescope the user will connect a laptop to the router or switch. To deal with the IP addresses we have two choices: configure the Pi to handle them or use a router. While the router might save us some work, we chose to configure the Pi to act as a DHCP server and hand out the IP addresses. To do this we used the `dnsmasq` tool, a free DNS forwarder and DHCP server. The configuration file is in Appendix F. The network was set up in such a way that the Pi and the GPIB controller have a static IP so that we know for sure they are at those addresses. The address of the laptop is not important as we do not need to address it. The user should now establish an SSH connection to the Pi. The IP addresses are set as 192.168.0.15 for the Pi and 192.168.0.142 for the GPIB controller. The Pi now automatically assigns addresses between 192.168.0.20 and 192.168.0.100 to other connecting devices. All scripts are executed on the Pi. They should therefore either be written on the Pi or be transferred to it beforehand. Normal observers should connect to the Pi with the `observer` account with password `observer`.

4.4.2 Observing with the Telescope

Once logged in the user should change directories to the Desktop folder where the main script for controlling the telescope is located: `ScopeServer`. This script should be run with the `sudo` command otherwise control through the GPIO pins is not possible. When run, the script will try to have all GPIB connected devices identify themselves to make sure all the required devices are there. If you see a `None` response it means a command did not receive a response and the connection should be checked. Once the identification finishes you will see a prompt starting with `>>>`. From here the user can either control the telescope directly or run a script. The commands for direct control are listed in Table 5.

When making a sweep of the sky it is not ideal to simultaneously measure the power and the atmospheric temperature. Measuring simultaneously will interrupt often the sweep and therefore cause the motor to stop or even miss a step. When the motor misses a step the angle of the horn will be off and the data becomes unusable. Instead it should be measured at one of the loads or set the temperature monitor to log the data and retrieve it later on.
Table 5: A list of commands available to the observer for manual control.

<table>
<thead>
<tr>
<th>Command</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>!calib</td>
<td>Forces recalibration of the zero position.</td>
</tr>
<tr>
<td>park</td>
<td>Returns the telescope to its zero position.</td>
</tr>
<tr>
<td>prep</td>
<td>Points the telescope to the cold load.</td>
</tr>
<tr>
<td>rot (±ang)</td>
<td>Rotates the telescope by the given amount of degrees. Will give an error when attempting to move out of bounds (i.e. below 0 or above 180 degrees).</td>
</tr>
<tr>
<td>!rot (±ang)</td>
<td>The same as rot but it ignores the limits. Use this with caution.</td>
</tr>
<tr>
<td>run (script)</td>
<td>Executes the script. Name should be given without the .py extension.</td>
</tr>
<tr>
<td>speed (rps) ⟨avgs⟩</td>
<td>Changes the setting of the power meter. rps sets the readings per second. avgs sets the number of readings averaged for a measurement.</td>
</tr>
<tr>
<td>speed?</td>
<td>Queries the current speed setting.</td>
</tr>
<tr>
<td>turbo (0</td>
<td>1)</td>
</tr>
</tbody>
</table>
5 Measurements of the Sun

The ultimate goal of this telescope is to measure the temperature of the CMB. It is nice to check though if we could observe more sources with it. The Sun for example is a near, bright source so it would be nice to see if we are able to measure it with our telescope. We also look into the possibility of detecting some of the brightest radio sources in the sky besides the Sun.

5.1 Solar Emission

Emission from the Sun has multiple origins. Charged particles accelerated in the Sun’s magnetic field will emit radiation via mechanisms such as synchrotron radiation, free free emission or a form of gyro radiation. This is mainly important for sunspots however. The conditions in these areas allow for gyro-resonant emission which is an efficient emission mechanism coming from electrons rotating around the magnetic field lines [Lee, 2007]. Another part of the emission comes from thermal radiation. If we consider the Sun to be a black body with a temperature equal to its surface temperature, it will radiate according to Planck’s law (Eqn. [7] with $T = 5778$ K).

Sunspots are part of the active Sun, while the thermal emission comes from the quiet Sun. Fig. 21 shows spectra of different sources. We see that at 11 GHz the contribution of the active Sun and quiet Sun are nearly equal and becomes a straight line. The plot is in log scale so a straight line means a power law relation. In the Rayleigh-Jeans limit ($h\nu \ll kT$) the Planck function reduces to a power law resulting in a line with slope 2,

$$B_\nu(T) = \frac{2kT\nu^2}{c^2}$$  \hspace{1cm} (11)

The radiation we measure with our telescope is thus mostly that of the quiet Sun. Based on its temperature we can estimate the brightness of the Sun and the power that we will receive. To find the power we receive from an object, we multiply the brightness ($B_\nu$) with the solid angle of the object ($\Delta\Omega$), the effective area of the telescope ($A_e$) and the bandwidth ($\Delta\nu$) giving

$$P = B_\nu\Delta\Omega A_e\Delta\nu.$$  \hspace{1cm} (12)

The estimated brightness is $B_\nu = 2.15 \times 10^{-16}$ W Hz$^{-1}$ m$^{-2}$ sr$^{-1}$. The solid angle of the Sun can vary slightly, so we use its size at a distance of 1 AU: 1919 arcsec.\footnote{NASA Sun Fact Sheet: http://nssdc.gsfc.nasa.gov/planetary/factsheet/sunfact.html} This gives a solid angle of $\Delta\Omega = 6.8 \times 10^{-5}$ sr. The effective area of our telescope is 0.023516 m$^2$ \footnote{Lap, 2015} and our bandwidth is 1.05 GHz. This gives an expected power of $P_\odot = 3.61 \times 10^{-13}$ W or $P_\odot = -94.42$ dBm. Combined with the CMB and the atmosphere mentioned earlier in Chapter 2, we expect a total power of $P = -84.55$ dBm arriving at the antenna when pointing at the Sun with an optical deph of $\tau = 0.05$. The gain of the telescope is around 60 dBm so this will be well within the observable range and we are able to measure the Sun. By measuring the Sun we can do two interesting experiments. We can try to measure properties of the main beam and we can try to determine the brightness temperature of the Sun.
Figure 21: Spectral distributions of various radio sources. Image credit: Tools of Radio Astronomy.

### 5.2 Observing the Sun

We had an unofficial first light at June 8, 2015 to see if everything worked. During this observation we tried to measure the Sun. This was in a phase where the hot and cold load were still under construction so the full range of 0 to 180 degrees was measured. It also means that the system is uncalibrated. Therefore we cannot draw conclusions from the received power itself. The beam size however is independent of calibration. Two data sets of this observation can be seen in Fig. 22.
Figure 22: Data from the unofficial first light. Top: data from 13:08 with no clouds. The little peak around 20 degrees is our supervisor John McKean standing in the beam. Bottom: data from 13:12.
5.3 Verifying the Beam Pattern

As mentioned before, we can determine the main beam size through observations of the Sun. This is possible due to the fact that we are able to consider the Sun as a bright and isolated point source. In the middle of the main beam, the point source would give the highest power peak and when the source is moving out of the main beam, the signal will decrease. Therefore, moving the beam of the antenna over the point source, the Sun, will map out the main beam. Determining the beam size requires a few steps:

1. Starting with the whole sky observation of 13:08, which is the data set where there were no clouds present. We only need the data where the radiation from the Sun is measured. The Sun is represented by the second peak in Fig. [22] Therefore, from now on we only consider the data in the range of 100 < angle < 150.

2. The data representing the Sun also involves a moving baseline and background radiation. To remove all the data except for the peak caused by the Sun, a function can be fitted on the background data. This is done by using the SciPy function `scipy.optimize.curve_fit`. The peak extends over the region from 115° to 145°. Fitting a polynomial to the data on the angles 100 to 115 and 145 to 150, will induce the green plot in Fig. [23a].

3. Subtracting the fit for the background from the original data will leave the data representing only the Sun. The main beam has got a Gaussian shape. Therefore, to correct for possible side lobes and interferences, a Gaussian model is fitted for the data. This model will represent the main beam of the antenna. Figure [23b] shows this fit, again in green.

4. To say something about the resolution of the antenna, we need to determine the Full Width Half Maximum (FWHM) of the main beam, also known as the Half Power Beam Width (HPBW). The FWHM is specified to be the limit of the antenna beam width and therefore equals the resolution. It can be determined by plotting a line at half of power of the maximum power in the Gaussian model and intersect with the model itself. The difference in angles at these intersection points represent the FWHM. The relation between the FWHM and the standard deviation $\sigma$ of a Gaussian function is $\text{FWHM} = 2\sqrt{2\ln 2}\sigma$. Looking at Fig. [23c] we can state that the FWHM of the antenna equals 11.99 ± 0.094°. This is somewhat off from the FWHM of 12.78° measured by [Lap][2015].

For a better estimate we need to take more measurements. We observed the Sun again at June 30 multiple times. The FWHM was then calculated the same way as above for each measurement. The results are listed in Table [6]. Averaging these values yields a beam size of 12.07 ± 0.13°. The error is calculated according to

$$\sigma = \sqrt{\frac{1}{N} \sum_{i=1}^{N} (x_i - \bar{x})^2}.$$ 

This is still off from the measured value in the lab. Most likely this is because pointing the telescope such that the sun passes exactly through the middle of the beam is difficult to do, causing us to measure a smaller FWHM. This may also explain the spread in the values themselves.
(a) The signal of the Sun enlarged. In green is the fit for the background.

(b) The Gaussian fit for the main beam (in green).

(c) Indication of the FWHM of the main beam. The value comes down to 11.99°.
<table>
<thead>
<tr>
<th>Time (hh:mm:ss)</th>
<th>HPBW (deg)</th>
<th>Error (deg)</th>
</tr>
</thead>
<tbody>
<tr>
<td>12:01:32</td>
<td>12.09</td>
<td>0.071</td>
</tr>
<tr>
<td>12:03:18</td>
<td>12.23</td>
<td>0.073</td>
</tr>
<tr>
<td>12:04:58</td>
<td>12.08</td>
<td>0.073</td>
</tr>
<tr>
<td>12:07:01</td>
<td>12.25</td>
<td>0.085</td>
</tr>
<tr>
<td>12:15:58</td>
<td>11.99</td>
<td>0.082</td>
</tr>
<tr>
<td>12:17:53</td>
<td>11.92</td>
<td>0.087</td>
</tr>
<tr>
<td>12:19:52</td>
<td>11.91</td>
<td>0.087</td>
</tr>
</tbody>
</table>
5.4 Measuring the Brightness Temperature of the Sun

If we have calibrated data, we can relate the power from the Sun to a brightness temperature. This can be determined by calculating the antenna temperature of the Sun. This is most easily done with the system temperature,

\[
T_{\text{sys}} = T_{\text{CMB}} e^{-\tau} + T_{\text{atm}} (1 - e^{-\tau}) + T_{\odot} e^{-\tau}. \tag{13}
\]

To determine the brightness temperature we thus need to know the gain of the telescope to convert the incoming power to a temperature and the optical depth of the atmosphere such that we can invert the above relation to obtain

\[
T_{\odot} = \left[ T_{\text{sys}} - T_{\text{CMB}} e^{-\tau} - T_{\text{atm}} (1 - e^{-\tau}) \right] e^\tau. \tag{14}
\]

Due to unforeseen complications with the cold load that emerged at the final stages, we unfortunately were not able to correctly calibrate the measurements. Hence it is not possible to determine the brightness temperature of the sun at this moment.
6 Measurements of other Sources

6.1 Galactic Synchrotron Radiation

Electrons in the Milky Way and cosmic ray electrons are accelerated in the magnetic field thus producing synchrotron radiation. To see if the galaxy’s synchrotron radiation will affect our measurements we need to know how strong it is with respect to the CMB at 11 GHz. Since this radiation is frequency dependent we need to know how it scales with $\nu$, which in turn depends on how the energy is distributed among the electrons. The energy spectrum of cosmic ray electrons follows a power law relation, giving the energy distribution given by

$$N(E) dE \propto E^{-\delta} dE. \quad (15)$$

This means the brightness temperature will follow the power law given by

$$T(\nu) \propto \nu^{-\beta} \quad (16)$$

with $\beta = (\delta + 3)/2$ being the spectral index for synchrotron radiation. In the range of $1.42 - 14.9$ GHz this spectral index is $\beta = 3.0$ [Platania et al., 1998]. The relative intensity of the galactic synchrotron radiation with respect to the intensity of the CMB is given by

$$\frac{I_{\text{sync}}}{I_{\text{CMB}}} \propto \nu^{\beta-3} (e^{h\nu/kT_{\text{CMB}}} - 1) \quad [\text{Readhead and Lawrence, 1992}]. \quad (17)$$

For a spectral index of 3.0 only the exponential part remains. This results in $\frac{I_{\text{sync}}}{I_{\text{CMB}}} \propto 0.21$. This relation is proportional to, but we can calculate an approximate contribution to the antenna temperature assuming this ratio and a gain of 60 dB. The power from just the CMB would be $P_{\text{CMB}} = -85.03$ dBm as seen in chapter 2. If the synchrotron radiation contributes 20% of this power this changes to $P_{\text{CMBS}} = -84.19$ dBm. This will result in a change of antenna temperature of $\Delta T_A = 0.046$ K. Contribution of synchrotron radiation is therefore most likely negligible in our case.

6.2 Astrophysical Radio Sources

Besides the Sun the sky houses many other radio sources. I have chosen four sources to investigate: Cassiopeia A, Cygnus A, Taurus A and Virgo A. Cas A is a well known supernova remnant and is one of the brightest radio sources in the sky. Cygnus A is also a strong radio source. It is a radio galaxy with two large lobes on either side. Taurus A or the Crab Nebula is one of the most well known supernova remnants. In its center is a pulsar emitting radio waves. Lastly Virgo A or M87 is an elliptical galaxy with an radio emitting AGN at its center. Table 8 lists the flux density and the resulting contribution to the antenna temperature of these sources. The flux densities are calculated according to [Baars et al., 1977] who use

$$\log_{10} S [\text{Jy}] = a + b \log_{10} \nu [\text{MHz}] + c \log_{10}^2 \nu [\text{MHz}] \quad (18)$$

for calculating the flux density of these sources. The parameters for this equation of each source are listed in Table 7.
Table 7: Selected sources with their parameters for the flux density equation.

<table>
<thead>
<tr>
<th>Source</th>
<th>a</th>
<th>b</th>
<th>c</th>
</tr>
</thead>
<tbody>
<tr>
<td>Cassiopeia A</td>
<td>5.745</td>
<td>-0.770</td>
<td>0</td>
</tr>
<tr>
<td>Cygnus A</td>
<td>7.161</td>
<td>-1.244</td>
<td>0</td>
</tr>
<tr>
<td>Taurus A</td>
<td>3.915</td>
<td>-0.299</td>
<td>0</td>
</tr>
<tr>
<td>Virgo A</td>
<td>5.023</td>
<td>-0.856</td>
<td>0</td>
</tr>
</tbody>
</table>

Table 8: Flux density of the selected sources at 11 GHz and antenna temperature for $\tau_0 = 0.01$.

<table>
<thead>
<tr>
<th>Source</th>
<th>Flux Density (Jy)</th>
<th>Antenna Temperature (K)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Cassiopeia A</td>
<td>322.29</td>
<td>0.0054</td>
</tr>
<tr>
<td>Cygnus A</td>
<td>135.99</td>
<td>0.0023</td>
</tr>
<tr>
<td>Taurus A</td>
<td>508.89</td>
<td>0.0086</td>
</tr>
<tr>
<td>Virgo A</td>
<td>36.61</td>
<td>0.00062</td>
</tr>
</tbody>
</table>

To be able to measure these sources the fluctuations in $T_{sys}$ need to be smaller than their respective contribution. This means that we need a sensitivity of 0.0086 K or lower to detect Taurus A, the brightest of the sources. From the available power meter settings we need to use an integration time of 0.8 seconds giving a sensitivity of 0.0069 K. With an integration time of 1.6 seconds we should be able to detect Cassiopeia A with a sensitivity of 0.0049 K. To detect Cygnus A we need to use the 12.8 second integration time to obtain a sensitivity of 0.0017 K. To detect Virgo A we need a much longer integration time. The longest integration time possible is 51.2 seconds. This gives a sensitivity of $8.6 \cdot 10^{-4}$ K. This is larger than the signal of Virgo A and fluctuations of this order are in the regime of fluctuations caused by the Allan time. Therefore, we are not able to detect this source with our telescope. As for the other sources, any measurement of the brightness temperature will most likely not possible. Theoretically it should be possible to make a detection of these sources by making a lot of measurements and then stacking them. This way the noise will be reduced compared to the signal of the source. However, this is based on the radiometer equation which in our case does not give the limiting sensitivity. Detection of other sources is therefore most likely not possible.

6.3 Non-Astrophysical Sources

The sky is also home to man made objects such as satellites. We discovered a satellite by accident when trying to find a clear piece of sky. When viewing the data we saw a large peak that was approximately 3.5 dBm higher than the background. This rules out the Sun as it is much
Weaker. It is also unlikely to be an astronomical source. The peak is located at 150° meaning it is 30° above the horizon. Checking the sky for possible sources at this location revealed that there are geostationary satellites at this location. These satellites are used for communications and emit high power radio signals. As the satellite is an even stronger point source than the Sun it is an ideal object to check the size of the beam. We use the measurements from June 30 to do this. The results are listed in Table 9. From these measurements we obtain an average HPBW of 12.61 ± 0.19°.

Figure 24: The mysterious signal. The left peak is the signal from the Sun and the right peak is from the satellite.
7 Future Improvements and Conclusion

Performance of the telescope is not yet optimal on the software and measuring side. Possible points of improvement are:

- Improve the overall speed and efficiency of the code.
- Change the way results are read from the power meter and temperature monitor.
- Use a more powerful motor or implement more advanced motor control.
- Use a more powerful version of the Pi or split the system into multiple components.

7.1 Code Efficiency and Speed

The code to control the telescope is self written. Something that can almost certainly be improved is the code itself. Certain routines can be optimized or special libraries designed for speed could be used. For even better performance we could switch to a compiled language such as C instead of Python, an interpreted language. This will however require a complete rewrite of the control code and it will reduce the flexibility of possible measurements. The performance gain from this switch will not be of great importance either as the other components in the system become the limiting factors. The code can still be optimized, but it will not give a significant improvement.

7.2 Reading Results from the Power Meter

The major bottleneck as of now is reading the power from the power meter. Due to the way it is done now this can take up to a second. This interrupt causes the telescope to wait after each step instead of moving continuously. Attempts have been made to get around this interrupt problem by running movement and measurement in parallel. It turned out however that the CPU on the Pi is not fast enough to handle this. The only way to change this is to use a different method of reading from the device. There is a specific setup for the GPIB controller and the power meter that makes it respond as soon as it has something to respond when asked for it. With this, response times should drop to nearly instantaneously.

7.3 Motor Power and Control

A stepper motor, like any motor, has a maximum load it can drive. Stepper motors in particular are limited by the amount of torque they can produce. The highest torque is gotten from the lowest rotation speeds. The faster it spins, the less torque it can generate. This means it might not be able to start at full speed. More advanced control schemes with, for example, acceleration towards maximum speed can solve this. A more powerful motor can generate more torque and would therefore be able to drive the horn at higher speeds right from the start.

Another option would be to implement more advanced control of the motor. The motor has a recommended operating voltage, but it can endure much higher voltages when moving.
Applying a higher voltage gives the motor an increased torque allowing it to start faster. The downside to this is that higher voltages also induce higher currents and hence increase heat output. If the motor becomes too hot it can potentially damage or kill it. This can be circumvented by having the voltage gradually increase as the motor gains speed and decrease as it slows down.

### 7.4 More Powerful Hardware

The last and most profound changes can be made at hardware level. The Pi is now the main unit that both controls all devices and processes all signals. To reduce the load on the system the control could be split up across multiple units. An example would be to use separate units for moving the motor and for reading the power meter and temperature sensors. These would then still be controlled from the Pi to coordinate everything, but the load would be spread over multiple devices allowing for smoother operation.

### 7.5 Conclusion

During this project I was responsible for control of and acquiring data from the telescope. While there is still room for improvement, it is in a well usable state for observations. The observations that I carried out had varying results. The brightness temperature of the Sun could not be determined due to unforeseen complications, however verifying the beam size was successful with both the Sun and the satellite. The Sun yields a beam size of $12.07^\circ$ and the satellite one of $12.61^\circ$. The latter is closest to the beam size determined in the lab. This is not surprising since the satellite was a much stronger source than the Sun.
8 Acknowledgements

I would like to thank John McKean for providing this project. It was a great opportunity to actually apply some of the theory learned in the Radio Astronomy course to a real life problem of building a radio telescope. My share in this project would not have been possible without all the help from the electronics department of SRON and the help of other people. I would like to thank Andrey Barychev, Andrey Kudchenko, Ronald Hesper, Henk Ode and Duc Nguyen for their help in constructing the electronics, acquiring measurement devices, trouble shooting and overall help.
References

KNMI. Data aquisition from a script. Last referenced 06-06-2015.


A Average Temperatures for June
Figure 25: Temperatures of June 2007 separated by part of the day.

Figure 26: Temperatures of June 2008 separated by part of the day.
Figure 27: Temperatures of June 2009 separated by part of the day.

Figure 28: Temperatures of June 2010 separated by part of the day.
Figure 29: Temperatures of June 2011 separated by part of the day.

Figure 30: Temperatures of June 2012 separated by part of the day.
Figure 31: Temperatures of June 2013 separated by part of the day.
B Average Cloud Coverage for June
Figure 32: Temperatures of June 2007 separated by part of the day.

Figure 33: Cloud coverage of June 2008 separated by part of the day.
Figure 34: Cloud coverage of June 2009 separated by part of the day.

Figure 35: Cloud coverage of June 2010 separated by part of the day.
Figure 36: Cloud coverage of June 2011 separated by part of the day.

Figure 37: Cloud coverage of June 2012 separated by part of the day.
Figure 38: Cloud coverage of June 2013 separated by part of the day.
C Telescope Server

#!/usr/bin/env python

from __future__ import division
import numexpr as ne
import datetime as dt
import errno
import os
import readline
import sys

try:
    import RPi.GPIO as GPIO
except RuntimeError:
    print 'Error importing GPi.GPIO. Try running as root.'

import StepperMotor as sm
import GPIBComm as gc

YELLOW_BOX_IP = '192.168.0.142'
PORT = 1234

label_m = ('calib', 'rot', 'park', 'prep', 'turbo')
funct_m = (m.calibrate, m.rot, m.park, m.prep, m.set_turbo)

label_c = ('measure', 'speed', 'speed?', 'send', 'unit')
funct_c = (c.measure, c.set_measurespeed, c.get_measurespeed, c.send, c.units)

CMD_CTRL = dict(zip(label_c, funct_c))
CMD_MOTR = dict(zip(label_m, funct_m))

def call(func, device, *args):
    ''' Processes the commands given to the server prompt. 
    Call functions specific to a certain device. Currently available devices are: controller, motor. 
    
    Args:
        func - the function to execute.
        device - name string of which type of device to address.
        misc args - additional arguments to the function.
    Returns:
        None
    
    if device == 'controller':
        if args is not None:
            CMD_CTRL[func](*args)
        else:
            CMD_CTRL[func]()
    elif device == 'motor':
        if args is not None:
            CMD_MOTR[func](*args)
        else:
            CMD_MOTR[func]()
    else:
        print 'Unknown device.'

# Telescope server #
# Server startup. Clear the terminal and check if everything is ok.

```python
os.system('clear')
print '=== Kapteyn Radio Telescope Server ==='
try:
    # See if the yellow box is alive.
    c = gc.GPIB_Controller(YELLOW_BOX_IP, PORT)
    c.send('++ver')
    c.read()
    # Ask the power meter to identify itself.
    c.send('++addr 14'); c.read(print=False)
    c.send('idn?')
    c.read()
    c.send('SENS:AVER:COUN:AUTO OFF')
    c.read(print=False)
    # Ask the temperature monitor to identify itself (it responds by itself).
    c.send('++addr 10'); c.read()
    # Default to powermeter.
    c.send('++addr 14'); c.read(print=False)
    c.units('dbm')
except:
    print 'An error occured. Is the Raspberry Pi connected to the switch and are the devices turned on?'
    sys.exit(-1)

# Initiate a motor and calibrate it to zero.
m = sm.StepperMotor(31, 29, 200, 72)
m.calibrate()

# Set up a directory on the users desktop to store measurements in.
try:
    os.makedirs(os.path.expanduser('/home/' + os.environ['SUDO_USER'] + '/Desktop/' + dt.datetime.now().isoformat()[:10].replace('-', '')))
except OSError as e:
    if e.errno != errno.EEXIST:
        raise
try:
    # Enter the main loop.
    while 1:
        # Ready the input.
        inp = raw_input('>>> ').lower()
        cm = inp.split()
        if len(cm) > 1:
            cmd = cm[0]
            args = cm[1::]
        elif len(cm) == 1:
            cmd = cm[0]
            args = []
        else:
            cmd = inp
        # Process input.
        if cmd == 'quit':
            c.quit()
            m.quit()
            break
        elif cmd == 'run':
            # User wants to run his/her own script.
            f = args[0]
            try:
                execfile('./' + f + '.py')
            except Exception as e:
                print e
                print 'Failed to execute script'
        elif cmd in CMD_CTRL:
            # Additional command
```

63
if len(args) > 0:
    call(cmd, 'controller', *args)
else:
    call(cmd, 'controller')
elif cmd in CMD_MOTR:
    if len(args) > 0:
        #call(cmd, 'motor', *[float(a) for a in args])
        call(cmd, 'motor', *[float(a) if a.strip().isdigit() else a for a in args])
    else:
        call(cmd, 'motor')
    # Special commands starting with !.
    elif cmd.startswith('!rot'):
        m.rot(float(cm[1]), override=True)
    elif cmd.startswith('!calib'):
        m.calibrate(override=True)
else:
    c.send(inp)
    c.read()
except KeyboardInterrupt:
    pass
finally:
    GPIO.cleanup()
D  GPIB Communication

#!/usr/bin/env python
import readline
import socket

class GPIB_Controller(object):
    ''' GPIB_Ethernet Controller.'''
    def __init__(self, ip, port):
        ''' Initialize a new connection to a GPIB-Ethernet Controller.
        Sets up a TCP connection to the controller at the specified ip address and port.
        '''
        self.ip = ip
        self.port = port
        self.sock = socket.socket(socket.AF_INET, socket.SOCK_STREAM, socket.IPPROTO_TCP)
        self.sock.connect((ip, port))
        self.sock.settimeout(1)
        # Initialize to the power meter.
        self.addr = 14
        self.send('++addr %d
' % addr); self.read(prnt=False)
        
    def measure(self):
        ''' Query the power meter to return a measurement. Assumes the address is set to the power meter.
        Args:
        None
        Return:
        str res - the answer of the power meter.
        OR
        None - if the address is wrong.
        '''
        self.send('FETC?
')
        res = self.read(prnt=False)
        return res

    def measure_temp(self, sens):
        ''' Query the temperature monitor to return a measurement. Assumes the address is set to the temperature
        monitor.
        Args:
        str sens - the sensor to read out
        Returns:
        str res - the answer of the temperature monitor.
        OR
        None - if the address is wrong.
        '''
        sensors = {'hot':1, 'cold':2, 'atm':4}
        self.send('KRDG? %d\n%sensors[sens])
        res = self.read(prnt=False)
        return res

    def read(self, prnt=True, maxiter=2):
        ''' Read responses from the buffer. Keep looping until we receive None to make sure we have the response.
        Args:
        None
        Return:
        None
        OR
        str res - the answer of the response.
        OR
        None - if the address is wrong.
        '''
        pass
bool prnt - print the output to the console (default is True).
int maxiter - number of secondary attempts to make when response is None (default is 2).
Returns:
str data or None - response of the GPIB controller.

```python
data = None
itr = 0
# Keep reading until we have everything.
while (data is None) and (itr < maxiter):
try:
    # Query the yellow box to speak to us.
    self.sock.send('++read \n')
    # Read from the socket; strip all whitespace (spaces, LF, CR etc.) on the right end.
    data = self.sock.recv(4096).rstrip()
except socket.timeout:
    # No response or empty lines.
    pass
finally:
    itr += 1
if prnt:
    print data
return data
```
def send(self, cmd):
    ''' Send a command to the GPIB controller.
    Args:
        str cmd - command to send.
    Returns:
        bool True/False - success or failure of sending the command.
    '''
    try:
        self.sock.sendall(cmd+'\n')
        return True
    except:
        return False
def set_measurespeed(self, rps=20, avg=1):
    ''' Sets the integration time of the telescope by specifying the readings per second and the number of averages taken for a measurement.
    Args:
        int rps - readings per second. Available: 20, 40 or 200.
        int avg - number of averages for a measurement. Available: 1, 2, 4, 8, 16, 32, 64, 128, 256, 512, 1024.
    Return:
        None
    '''
    rp = [20, 40, 200]
av = [2**i for i in range(11)]
rps = int(rps); avg = int(avg)
if (rps not in rp) or (avg not in av):
    print 'Invalid value in measure speed.'
else:
    self.send('SENS:AVER:COUN %d' % avg)
    self.send('SENS:SPE %d' % rps)

def get_measurespeed(self, p=True):
    ''' Requests the averaging rate and readings per second.
    Args:
        bool p - print the values to the terminal or not.
    Returns:
        int read - readings per second.
        int aver - the number of averages taken for a measurement.
    '''
    self.send('SENS:SPE?')
aver = self.read(prnt=p)
self.send('SENS:AVER:COUN?')
read = self.read(prnt=p)
return read, aver

def units(self, un):
    ''' Set the units of measurement. Possible units are 'W' or 'dBm'.
    Args:
        str un - unit to return.
    Return:
        None
    '''
    u = un.lower()
    if u == 'w':
        self.send('UNIT:POW W')
        self.read(prnt=False)
    elif u == 'dbm':
        self.send('UNIT:POW DBM')
        self.read(prnt=False)
    else:
        print 'Invalid unit.'

def quit(self):
    ''' Correctly close the connection to the controller.
    Args:
        None
    Return:
        None
    '''
    self.sock.close()
E Stepper Motor Control

#!/usr/bin/env python
from __future__ import division
import numexpr as ne
import readline
import time

try:
    import RPi.GPIO as GPIO
except RuntimeError:
    print 'Error importing GPi.GPIO. Try running as root.'

class StepperMotor(object):
    ''' Stepper motor.
    '''
    def __init__(self, pdir, pstep, stprev, revdisk):
        # Raspberry Pi GPIO pin setup #
        self.PIN_DIR = pdir
        self.PIN_STEP = pstep
        # Set pin numbering as on the pi
        GPIO.setmode(GPIO.BOARD)
        # Set pins as output pins.
        GPIO.setup(pdir, GPIO.OUT)
        GPIO.setup(pstep, GPIO.OUT)

        # Motor constants #
        self.REV_DISK = revdisk # Number of stepper revolutions required for one
disk revolution.
        self.STEP_STPM = stprev # Number of steps per stepper motor revolution.
        self.STEP_DISK = self.STEP_STPM * self.REV_DISK # Number of steps per disk revolution.
        self.SLEEP_TIME = 0.002 # Number of seconds to sleep between steps.
        self.TURBO = False # Shorter sleep time hence faster movement.
        self.PARK_POS = 0 # Parking position.
        self.ANGLE = 0 # Tracks telescope position.
        self.CALIBRATED = False # Check for calibration.
        self._rotating = False

    def calibrate(self, override=False):
        ''' Calibrate the telescope to the parking posistion.
        Args:
            float cal_ang - set the current angle to the parking angle.
            bool override - override the existing calibration.
        Returns:
            None
        ...
        if (not self.CALIBRATED) or (self.CALIBRATED and override):
            # Set the current angle to this position.
            self.ANGLE = 0
            self.CALIBRATED = True
        else:
            print 'System already calibrated.'

    @property
    def is_rotating(self):
        ''' Returns wether the motor is rotating or not.
        ...
        return self._rotating

68
def park(self):
    ''' Return the telescope to park position.
    Args:
    None
    Returns:
    None
    '''
    self.rot(-self.ANGLE)

def prep(self):
    ''' Prepare the telescope for calibration starting from zero point.
    This points the horn to the cold load at -15 degrees from the horizontal.
    Args:
    None
    Returns:
    None
    '''
    self.rot(-15, override=True)

def rot(self, ang, override=False):
    ''' Rotate the disk by a given amount of degrees.
    There are built in safety limits so that the telescope will not move below 0 or over 180 degrees. This
    protection should not be overridden unless you know what you are doing. In general it should not be
    necessary to do this.
    Args:
    float ang - rotation angle in degrees.
    bool override - override the safety limits. Use carefully!
    Return:
    None
    '''
    self._rotating = True
    # Update the telescope position.
    if (((self.ANGLE + ang) > 180) or ((self.ANGLE + ang) < 0)) and not override:
        print 'Going too far! Aborting.'
        return
    self.ANGLE += ang
    # Set forwards or backwards.
    if ang < 0:
        self.set_direction(True)
    else:
        self.set_direction(False)
    # Calculate the number teps needed.
    steps = (abs(ang) / 360) * self.STEP_DISK
    istep = 0
    while istep < steps:
        self.step()
        istep += 1
    self._rotating = False

def set_direction(self, direction):
    ''' Set the movement forwards or backwards.
    Args:
    bool direction - True: backwards, False: forwards.
    Return:
    None
    '''
    GPIO.output(self.PIN_DIR, direction)

def set_turbo(self, x):
    ''' Change the delay between switching pin states.
    If turbo is on the torque produced by the motor will be lower.
    Args:
int x - 0 or 1 for turbo on or off.
Return: None
    if x:
        #self.TURBO = True
        self.SLEEP_TIME = 0.0013
    else:
        #self.TURBO = False
        self.SLEEP_TIME = 0.002

def step(self):
    """Make the stepper motor do one step.
    Sleep occurs _after_ state change. This means the following cycle: high, sleep, low, sleep. This means
    movement is slower by a factor of two w.r.t to the sleep time.
    ""
    Args:
        None
    Return:
        None
    GPIO.output(self.PIN_STEP, True)
    time.sleep(self.SLEEP_TIME)
    GPIO.output(self.PIN_STEP, False)
    time.sleep(self.SLEEP_TIME)

def quit(self):
    """Properly exit the motor controls.
    ""
    Args:
        None
    Return:
        None
    GPIO.cleanup()

if __name__ == '__main__':
    mot = StepperMotor(38, 40, 200, 72)
    while True:
        cmd = raw_input('>>> ').lower()
        com = cmd.split(' ')
        if len(com) == 2:
            command = com[0].lower()
            arg = com[1]
            if command == 'rot':
                deg = ne.evaluate(arg)
                mot.rot(deg)
            if command == '!rot':
                deg = ne.evaluate(arg)
                mot.rot(deg, override=True)
        elif cmd == 'calib':
            mot.calibrate()
        elif cmd == '!calib':
            mot.calibrate(override=True)
        elif cmd == 'park':
            mot.park()
        elif cmd == 'quit':
            break
        else:
            print 'Invalid command.'
    GPIO.cleanup()
F DNS Configuration

# Set up for the ethernet port of the Pi.
interface=eth0
# Give out IP addresses between 0 and 100 with 1 hour lease time.
dhcp-range=192.168.0.20,192.168.0.100,255.255.255.0,1h

# Logs
log-facility=/var/log/dnsmasq.log
log-async
log-dhcp