# Formation of dwarf galaxies (in voids)

or

# Dwarf galaxies in the Universe Why aren't there more of them?

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# Dwarf galaxies in cosmological voids?

N-Body simulations: "Yes"



Surveys: "No (???)"

Gottlöber et al. 2003

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 $L = 20 \ h^{-1} Mpc$ 

### The halo mass function

pc<sup>-3</sup>1

r°s

Nhalo (>M)

infinite (?) number of small halos



### Do we see all these halos?

Bullock, Kravtsov, Weinberg



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#### The galaxy dark matter connection

populate simulated dark matter distributions with observed galaxies



van den Bosch, Yang, Mo, 04

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**Textbook solution:** 

(Padmanabhan)



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### $nk_{\rm B}T / n^2 \Lambda V$ (GM/R<sup>3</sup>)<sup>-1/2</sup>



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(Padmanabhan)



 $nk_{\rm B}T / n^2 \Lambda V$  (GM/R<sup>3</sup>)<sup>-1/2</sup>

Bremsstrahlung:

R = 74 kpc $M = 3 \times 10^{11} \text{ M}_{\odot}$ 



# Cooling of primordial plasma



### Cosmological hydrodynamical void simulation

Diameter=16 Mpc $\Omega_{\rm M}$ =0.03Mass resolution (gas)~2×10<sup>5</sup> h<sup>-1</sup> M<sub>☉</sub>

TreeSPH Gadget2 Radiative cooling UV-heating Star formation subgrid model feedback





## Multiphase model

Yepes et al. 1997 Springel & Hernquist 2002

Aim: determine SFR  $\dot{
ho}_*(
ho,T)$ 



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### **Baryon fraction**

Halos below few times 10<sup>9</sup> M<sub>☉</sub> are *baryon-poor* 

*Characteristic mass* scale depends on redshift





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*Characteristic mass* scale depends on redshift



# Mimic the UV-background: lower T limit



Crain et al. 2006

### Redshift evolution of the baryon fraction

*Characteristic mass* scale decreases with redshift





### Redshift evolution of the baryon fraction

*Characteristic mass* scale decreases with redshift





### Redshift evolution of the baryon fraction







mass

characteristic

M<sub>c</sub> rises significantly with redshift





M<sub>c</sub> rises significantly with redshift





mass

characteristic

M<sub>c</sub> rises significantly with redshift





M<sub>c</sub> rises significantly with redshift





# **Filtering Mass**

$$\frac{\mathrm{d}^2 \delta_X}{\mathrm{d}t^2} + 2H \frac{\mathrm{d}\delta_X}{\mathrm{d}t} = 4\pi G \bar{\rho} (f_X \delta_X + f_b \delta_b)$$

$$\frac{\mathrm{d}^2 \delta_{\mathrm{b}}}{\mathrm{d}t^2} + 2H \frac{\mathrm{d}\delta_{\mathrm{b}}}{\mathrm{d}t} = 4\pi G \bar{\rho} (f_{\mathrm{X}} \delta_{\mathrm{X}} + f_{\mathrm{b}} \delta_{\mathrm{b}}) \left| -\frac{c_{\mathrm{S}}^2}{a^2} k^2 \delta_{\mathrm{b}} \right|$$

$$rac{\delta_b}{\delta_X} = 1 - rac{k^2}{k_F^2}$$





Gnedin & Hui 1997

# Filtering mass (cont.)

$$\frac{1}{k_{\rm F}} = \frac{3}{2} \Omega_0 \frac{1}{D(a)} \int_0^a da' \frac{D}{S \, a' k_{\rm J}^2} \int_{a'}^a da'' \frac{1}{a''^2 S}$$
$$S^2 = 1 + \Omega_0 (1/a - 1) + \Omega_\Lambda (a^2 - 1)$$

 $c_{
m s}^2 = rac{3}{5} rac{k_{
m B} \langle T 
angle_{
m something}}{\mu m_{
m p}}$ 

$$M_{
m F} = rac{4\pi}{3} 
ho \; \left(rac{2\pi a}{k_{
m F}}
ight)^3$$



# Filtering mass (final)





### Baryon fraction: Void + Group

In dense environments the characteristic mass corresponds to that in void regions





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# Tidal stripping with cool gas + stars



# Gas accretion, schematically





Hot halo





Internatio

• Accretion

# ... more realistically shaped



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"Cold mode" (Keres et al. 04) of galactic gas accretion: gas creeps along the equilibrium line between heating and cooling



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$$k_B T_{\rm vir} = \frac{1}{2} \mu m_p \frac{G M_{\rm vir}}{r_{\rm vir}}$$

$$\frac{M_{\rm vir}}{4/3\,\pi\,r_{\rm vir}^3} = \Delta_c(z)\,\langle\rho\rangle$$

 $T_{
m entry} \ge T_{
m vir}$ 

$$\frac{M_{\rm c}(z)}{10^{10} h^{-1} M_{\odot}} \simeq \left\{ \frac{T_{\rm entry}(z)}{3.5 \times 10^4 \,\mathrm{K}} \frac{1}{1+z} \right\}^{\frac{3}{2}} \left\{ \frac{\Delta_c(0)}{\Delta_c(z)} \right\}^{\frac{1}{2}}$$

$$\frac{M_{\rm c}(z)}{10^{10} \, h^{-1} \, M_{\odot}} = \left\{ \tau(z) \; \frac{1}{1+z} \right\}^{3/2} \; \left\{ \frac{\Delta_c(0)}{\Delta_c(z)} \right\}^{1/2}$$



Max gas temperature

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Condition for suppression

$$T_{\rm entry} \ge T_{\rm vir}$$

 $\frac{M_{\rm c}(z)}{10^{10} \, h^{-1} M_{\odot}} \simeq \begin{cases} T_{\rm entry}(z) \\ 3.5 \times 10^4 \, {\rm K} \\ 1 + 1 \end{cases}$ 

$$\left[\frac{1}{2}\right]^{\frac{3}{2}} \left\{\frac{\Delta_c(0)}{\Delta_c(z)}\right\}^{\frac{1}{2}}$$

$$\frac{M_{\rm c}(z)}{10^{10} h^{-1} M_{\odot}} = \left\{ \tau(z) \ \frac{1}{1+z} \right\}^{3/2} \left\{ \frac{\Delta_c(0)}{\Delta_c(z)} \right\}^{1/2}$$



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Measurement M<sub>c</sub>

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## Mass accretion history



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# Baryon poor small halos



derived from the characteristic mass scales

τ:

T<sub>entry</sub>: taken from the densitytemperature phase space

Good agreement in particular for the newly accreted (cooled) mass



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derived from the characteristic mass scales

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UUE

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### The characteristic mass is "robust"

even a significantly different heat input has only little effect

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### Gas accretion revisited

Total heat input by UV heating

10<sup>43</sup> -10<sup>47</sup> erg yr<sup>-1</sup>

(Very crude estimate!)





### In which reservoir does the halo cool?



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## Summary

 $\bullet$  Photoheating suppresses the condensation of gas in halos  $< M_{c}$ 

•  $T_{vir} < T_{entry}$  is a very good criteria for ongoing accretion

• Photoheating by UV-background is not sufficient to explain the paucity of dwarf galaxies

• Galactic feedback (even without winds) provides much more heat, and suppresses therefore accretion much stronger

