

STRUCTURE OF GALAXIES

Lecture 3. Distribution of photometric parameters; photometric evolution and population synthesis.

Distributions of central surface brightness and scalelength in disks.

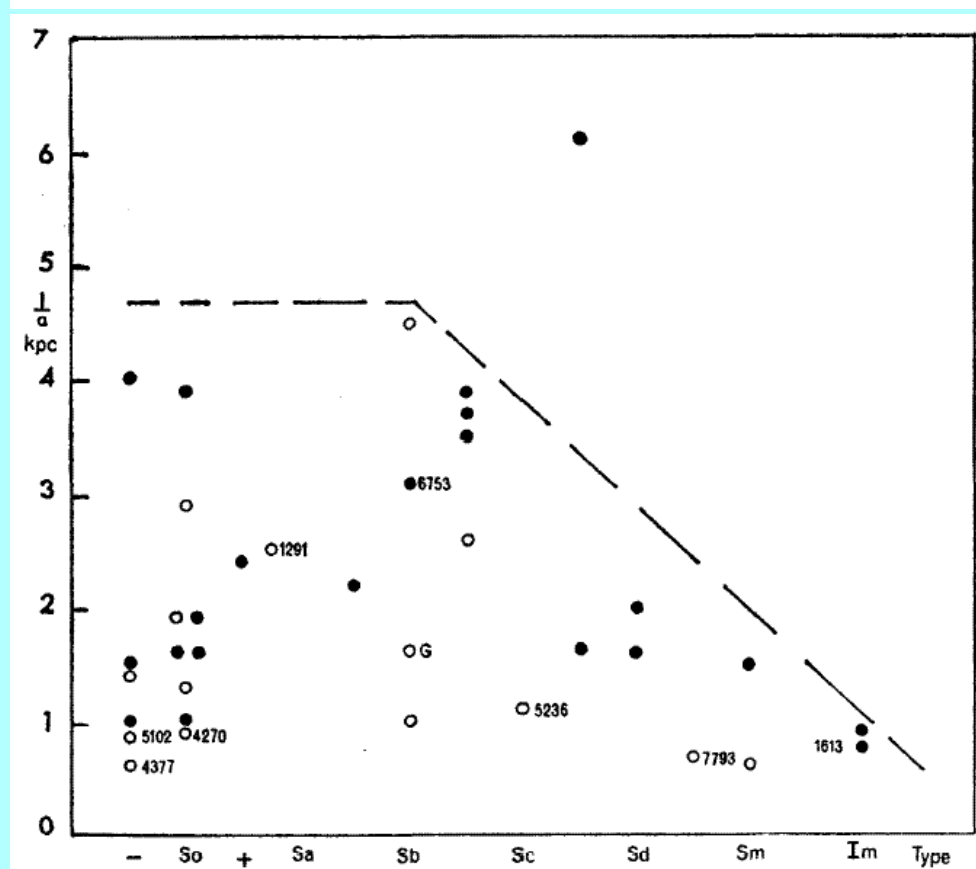
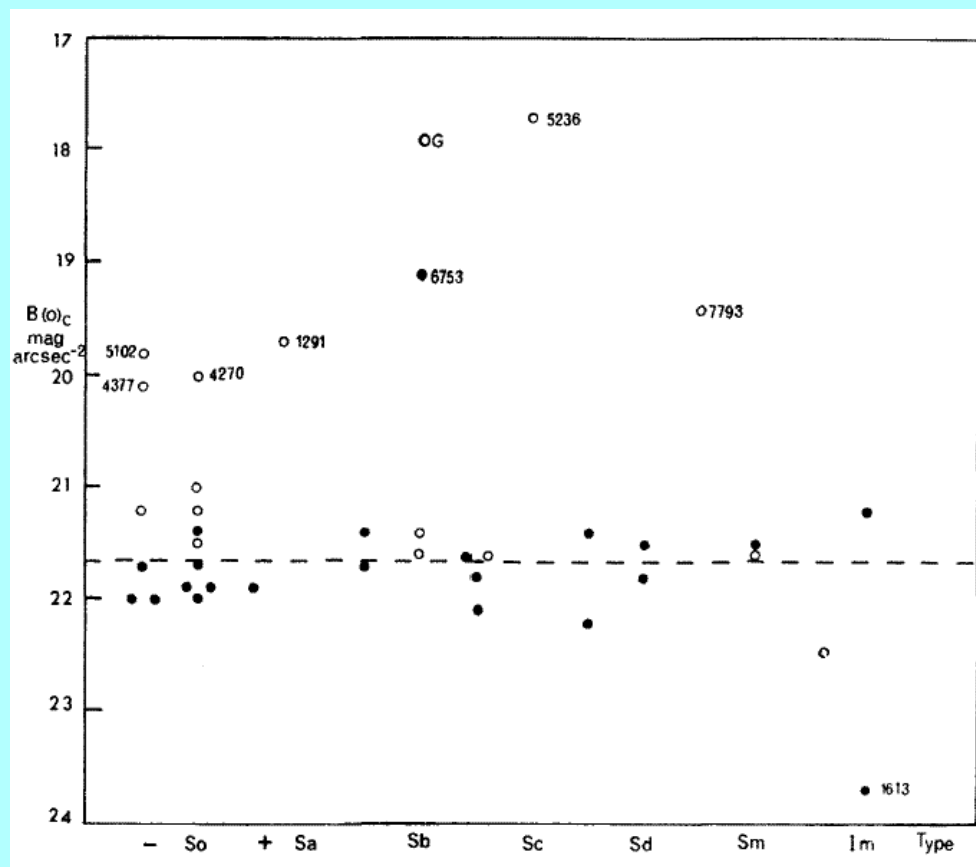
Freeman* was the first to study the distribution of properties of exponential disks.

His results are in the following two figures; the small range of (extrapolated) face-on, central surface brightness is known as “Freeman's Law”:

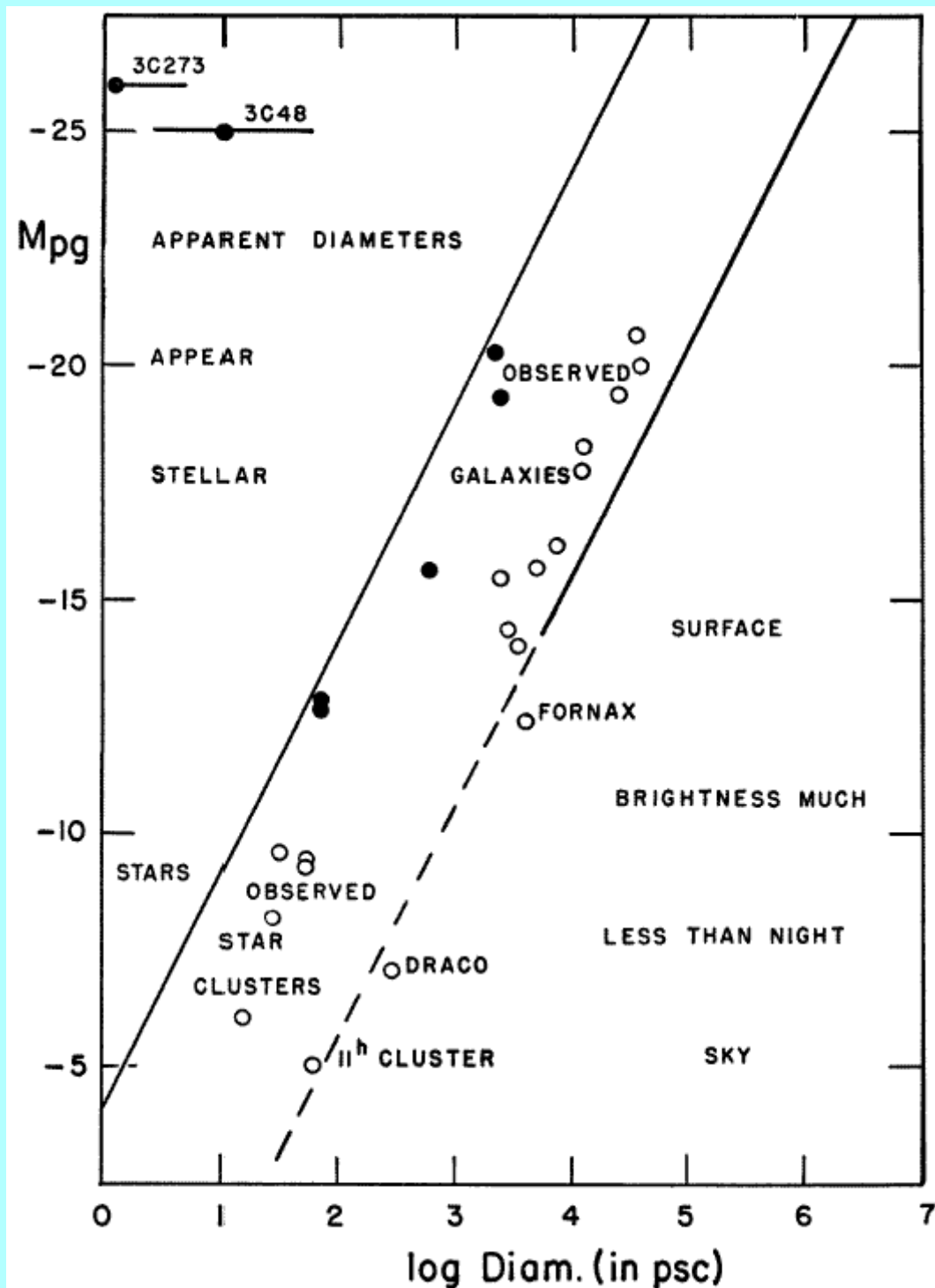
$$\mu_0 = 21.67 \pm 0.30 \text{ B - mag arcsec}^{-2}$$

This has generated considerable discussion. The problem is that samples need to be statistically complete and Freeman's sample had serious selection effects.

*Ap.J. 160, 811 (1970)



The selection was discussed first by [Arp](#) *.



*Ap.J. 142, 402 (1965).

The selection effects operating here are:

- For a particular luminosity and a faint μ_0 we get a large h , but for the most part the object is fainter than sky.
- For the same luminosity and a bright μ_0 we get small h and the object will appear star-like.

We will quantify this below.

First we will consider the V/V_{\max} -test for completeness.

For this we need to know the selection criteria of the sample. These could be for example all objects down to a certain angular diameter (at some isophotal level) or integrated apparent magnitude.

Suppose that an object has a distance R . Now shift it in distance until it drops out of the sample due to the completeness limit and call this distance R_{\max} .

Then we have V as the volume corresponding to R and V_{\max} as the volume relating to R_{\max} .

Now, in case of a uniform space distribution each object has an uniform chance to be actually located throughout the volume V_{\max} .

In otherwords, the property V/V_{\max} calculated for all objects in the sample should be distributed uniformly over the interval 0 to 1.

Note that V/V_{\max} can usually be calculated without knowing the actual distance.

In practice the test is to calculate $\langle V/V_{\max} \rangle$. For a **complete** sample it is required that

$$\langle V/V_{\max} \rangle = 0.5.$$

The error in $\langle V/V_{\max} \rangle$ is $(12 n)^{-1/2}$. This is so, because all numbers between 0 and 1 have an average of 0.5 and a dispersion of $\sqrt{12}$.

Selection and Freeman's law.

Disney* suggested that Freeman's law is the result of sample selection (and not only of incompleteness).

In the process he also addressed the equivalent for elliptical galaxies, called Fish's law.

The analysis was later extended as in the following†.

Assume luminosity-law (in linear units)

$$\sigma(R) = \sigma_0 \exp - (R/h)^{1/b}$$

This means:

$b = 1$: exponential disk

$b = 4$: $R^{1/4}$ bulge or elliptical galaxy.

We then have for the integrated luminosity:

$$L_{\text{tot}} = \int_0^\infty 2\pi R \sigma(R) dR = (2b)! \pi \sigma_0 h^2$$

*Nature 263, 573 (1975)

†Disney & Phillipps, Mon. Not. R.A.S. 205, 1253 (1983);
see also Davies, Mon. Not. R.A.S. 244, 8 (1990)

a. Diameter selection.

Suppose that a sample is complete for a radius larger than θ_{lim} arcsec at an isophote of μ_{lim} magnitudes arcsec⁻². For a radius R and a distance d the angular diameter is $\theta = R/d$ radians.

For clarity we now do the derivation only for an exponential disk.

The disk has an apparent radius

$$R_{\text{app}} = h \ln \left(\frac{\sigma_{\circ}}{\sigma_{\text{lim}}} \right)$$

In magnitudes arcsec⁻² this is

$$R_{\text{app}} = 0.4 \ln 10 \, h (\mu_{\text{lim}} - \mu_{\circ})$$

With $L = 2\pi\sigma_{\circ}h^2$ this becomes

$$R_{\text{app}} = \frac{0.4 \ln 10}{\sqrt{2\pi}} \left(\frac{L}{\sigma_{\circ}} \right)^{-1/2} (\mu_{\text{lim}} - \mu_{\circ})$$

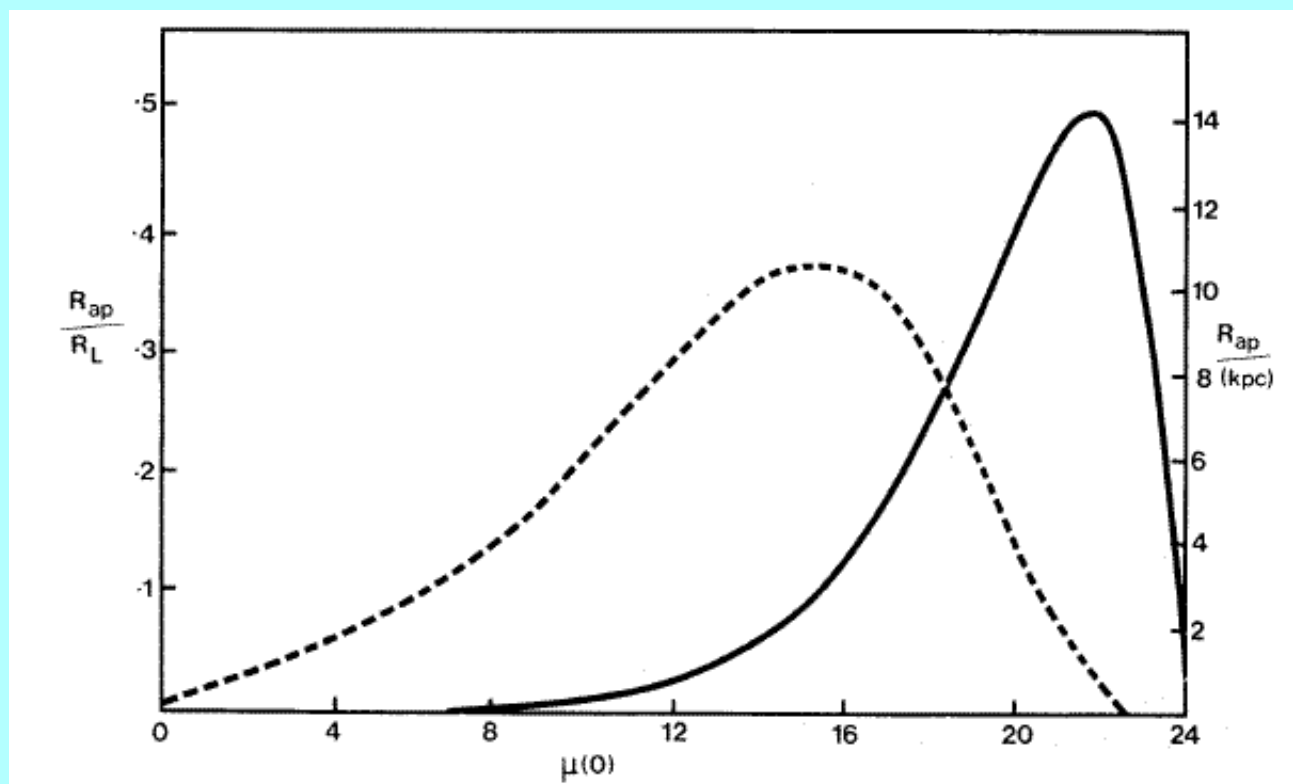
This can be rewritten as

$$R_{\text{app}} \sqrt{\frac{\pi\sigma_{\text{lim}}}{L}} = \frac{0.4 \ln 10}{\sqrt{2}} 10^{-0.2(\mu_{\text{lim}} - \mu_{\circ})} (\mu_{\text{lim}} - \mu_{\circ})$$

The square-root term on the lefthand side is a kind of fiducial radius, that Disney and Phillipps write as R_L .

The case with $\beta = 4$ for elliptical galaxies is

$$\frac{R_{app}}{R_L} = \frac{(0.4 \ln 10)^4}{\sqrt{8!}} 10^{-0.2(\mu_{lim} - \mu_o)} (\mu_{lim} - \mu_o)^4$$



In this figure we see the behavior of R_{app}/R_L as a function of the central surface brightness μ_o for the case of a diameter selection at an isophote of 24 (B-)magnitudes arcsec⁻².

The apparent diameter for exponential disks (full line) peaks at a central surface brightness of $(\mu_{\text{lim}} - \mu_{\circ}) = 2.171$; for elliptical galaxies (dashed line) this occurs at $(\mu_{\text{lim}} - \mu_{\circ}) = 8.686$.

Now when we express surface brightness μ in magnitudes arcsec⁻² and distances (such as $\sqrt{\sigma/L}$) in parsec we can derive

$$\frac{L}{\sigma_{\text{lim}}} = 10^{0.4(\mu_{\text{lim}} - M + 5)}$$

Then for the maximum distance for a galaxy to remain in the sample d in parsec and angular radius limit θ_{lim} in arcsec we get

$$d_{\text{size}} = \frac{0.4 \ln 10}{\sqrt{2\pi}} \frac{\mu_{\text{lim}} - \mu_{\circ}}{\theta_{\text{lim}}} 10^{0.2(\mu_{\circ} - M + 5)}.$$

For the general case the result is

$$d_{\text{size}} = \frac{(0.4 \ln 10)^b}{\sqrt{\pi(2b)!}} \frac{(\mu_{\text{lim}} - \mu_{\circ})^b}{\theta_{\text{lim}}} 10^{0.2(\mu_{\circ} - M + 5)}$$

The **maximum** of **d** occurs at

$$\mu_{o,\max} = \mu_{\text{lim}} - \frac{b}{0.2 \ln 10}$$

b. Integrated magnitude selection

Now the sample is supposed complete up to a limiting integrated apparent magnitude m_{lim} within an isophote μ_{lim} .

Assume that the image is overexposed at isophote μ_M to allow for photographic surveys and define

$$s = 0.4 \ln 10 (\mu_M - \mu_0) \quad ; \quad p = 0.4 \ln 10 (\mu_{\text{lim}} - \mu_0)$$

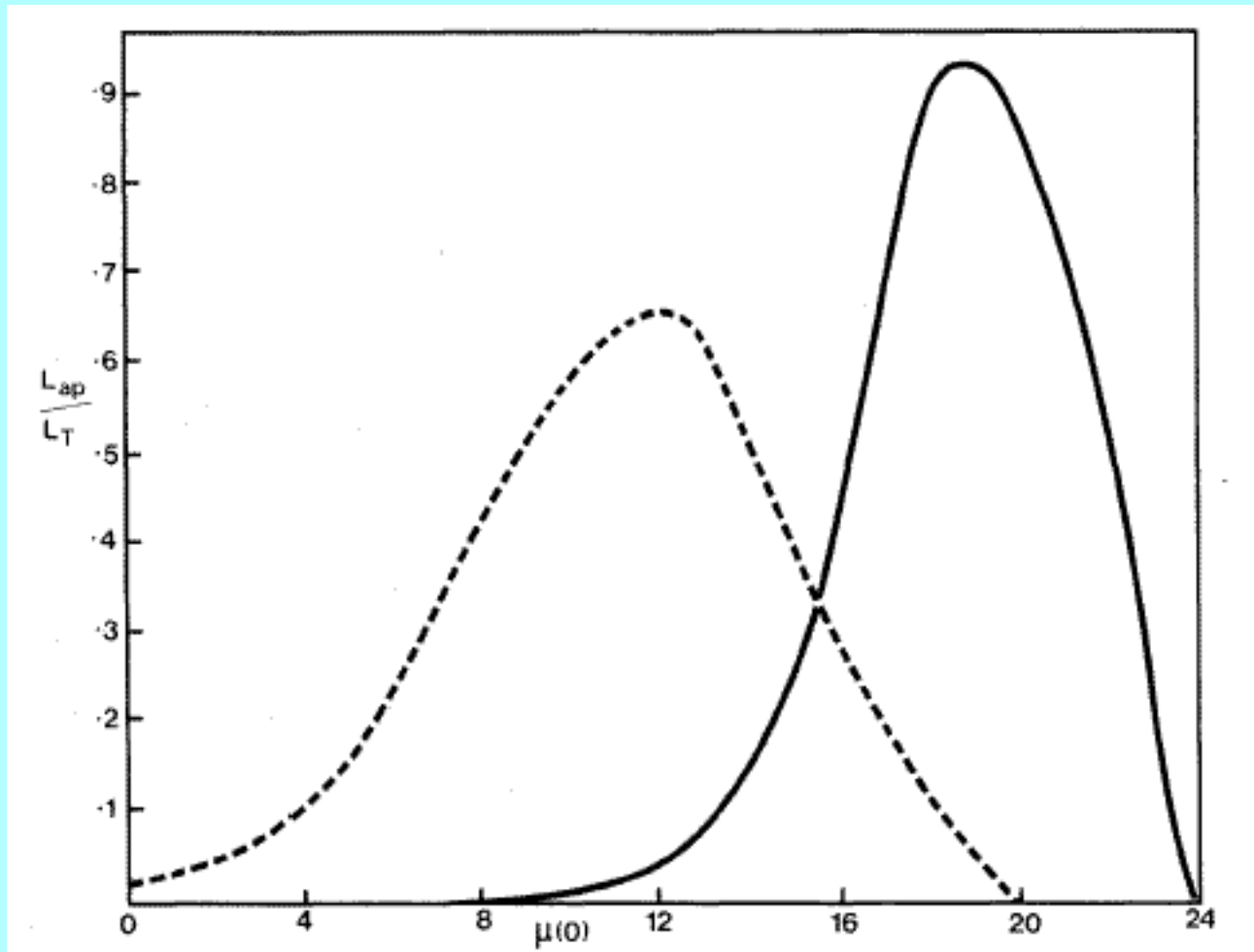
The maximum distance then comes out as

$$d_{\text{magn}} = \left[A_s e^{-s} - A_p e^{-p} \right]^{1/2} 10^{0.2(m_{\text{lim}} - M + 5)}$$

with

$$A_s = \sum_{n=0}^{n=2b} \frac{s^n}{n!} \quad ; \quad A_p = \sum_{n=0}^{n=2b-1} \frac{p^n}{n!}$$

This figure below is for a limiting isophote of 24 magnitudes arcsec^{-2} and a saturation isophote of 19 magnitudes arcsec^{-2} .



Again we see maxima as for diameter selection.

Note that both diameter and magnitude selection work in favor of disks around Freeman's surface brightness and elliptical systems near Fish's value.

Some actual values:

For Palomar Sky Survey:

$$\mu_{\text{lim}} \approx 24 \text{ B-mag arcsec}^{-2}$$

$$\mu_{\text{M}} \approx 19 \text{ B-mag arcsec}^{-2}$$

Diameter selection: d^3 peaks at:

$$- 21.8 \text{ B-mag arcsec}^{-2} \text{ for } b = 1$$

$$- 15.3 \text{ B-mag arcsec}^{-2} \text{ for } b = 4$$

Magnitude selection: d^3 peaks at:

$$- 18.5 \text{ B-mag arcsec}^{-2} \text{ for } b = 1$$

$$- 12.0 \text{ B-mag arcsec}^{-2} \text{ for } b = 4$$

Observed:

$$b = 1: 21.6 \pm 0.3 \text{ B-mag arcsec}^{-2} \text{ (Freeman's law)}$$

$$b = 4: 14.8 \pm 0.9 \text{ B-mag arcsec}^{-2} \text{ (Fish's law)}$$

In any catalogue each galaxies has a value for d according to the selection criteria.

If both diameter and magnitude selection play a role the smallest of the two values is the appropriate one.

We can then define the **visibility** as the value for d^3 for each galaxy: in an unbiased sample and a uniform distribution a value of μ_0 will occur at a frequency $\propto d^3$.

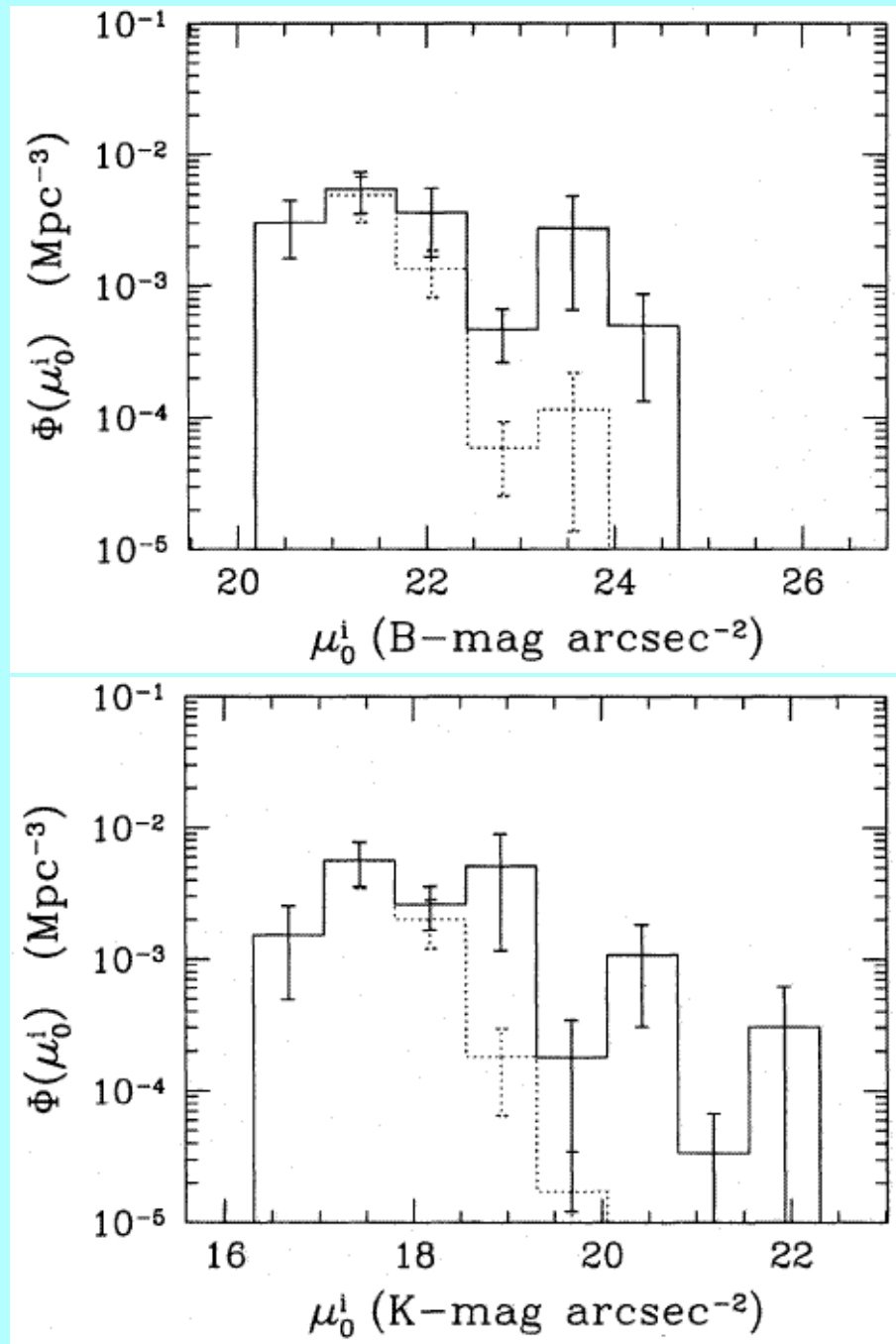
The equations for the visibility can of course also be used to correct complete sample for the volumes over which galaxies are sampled as a function of their properties in order to obtain space densities as a function of parameters.

This can be used to study the question of the origin of Freeman's law and whether it results from selection effects.

Allen & Shu* were the first to suggest that the selection only works at the faint level and that there is only a real **upper** limit to the central surface brightnesses.

*Ap.J. 227, 67, (1979)

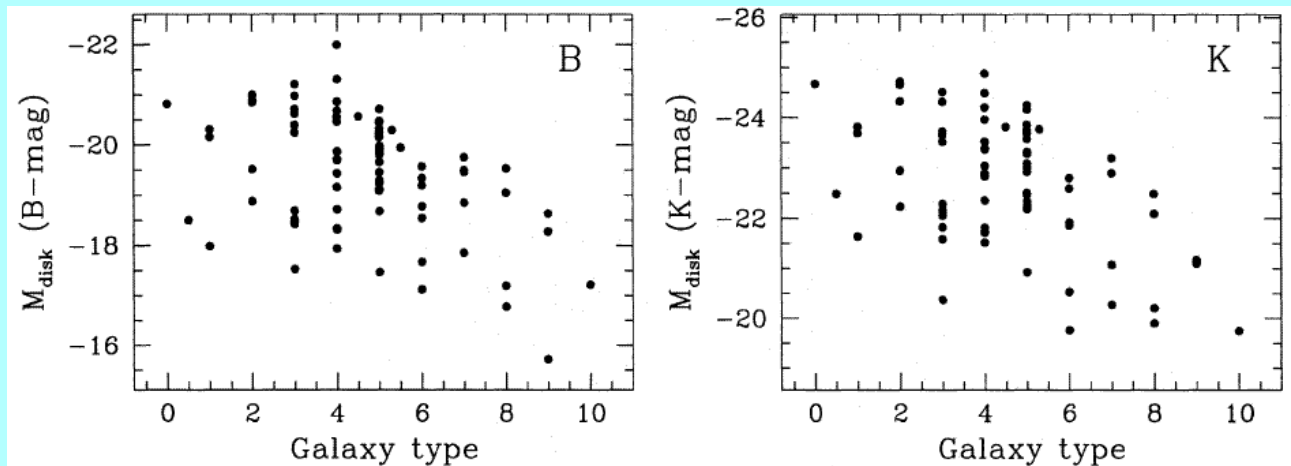
This is confirmed by [de Jong*](#), who also confirmed that the faint surface brightness disks are all of late type[†].



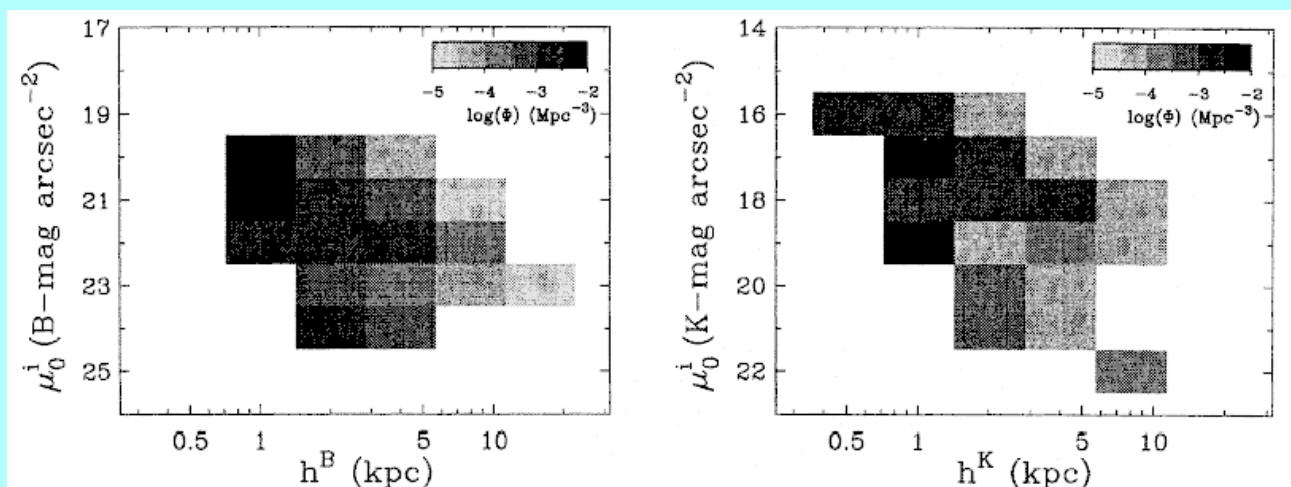
*A.&A. 313, 45 (1996)

†van der Kruit, A.&A. 173, 59 (1987)

This is related to the fact that late type galaxies generally have fainter disks.



Data can be combined in **bi-variate distribution functions**.



From a weighing with the total luminosity it can be estimated that high surface brightness galaxies probably provide the majority of the luminosity density in the universe.

Photometric evolution

Fundamental discussion is by Tinsley*.

The Initial Mass Function (IMF) is the distribution over stellar masses during star formation.

It is determined in the solar neighborhood independently for low and high mass stars:

- Low masses ($M < 1M_{\odot}$) from general distribution of masses of older stars in the disk, since these are all still present.
- High masses ($M > 1M_{\odot}$) from distribution of stellar masses in actual clusters and associations.

Normalisation is done such that the two parts join smoothly at $\approx 1M_{\odot}$ (continuity constraint).

*Fund. Cosmic Physics 5, 287 (1980)

An usefull analytic form of IMF:

$$\phi(M) = x M_L^x M^{-(1+x)} dM$$

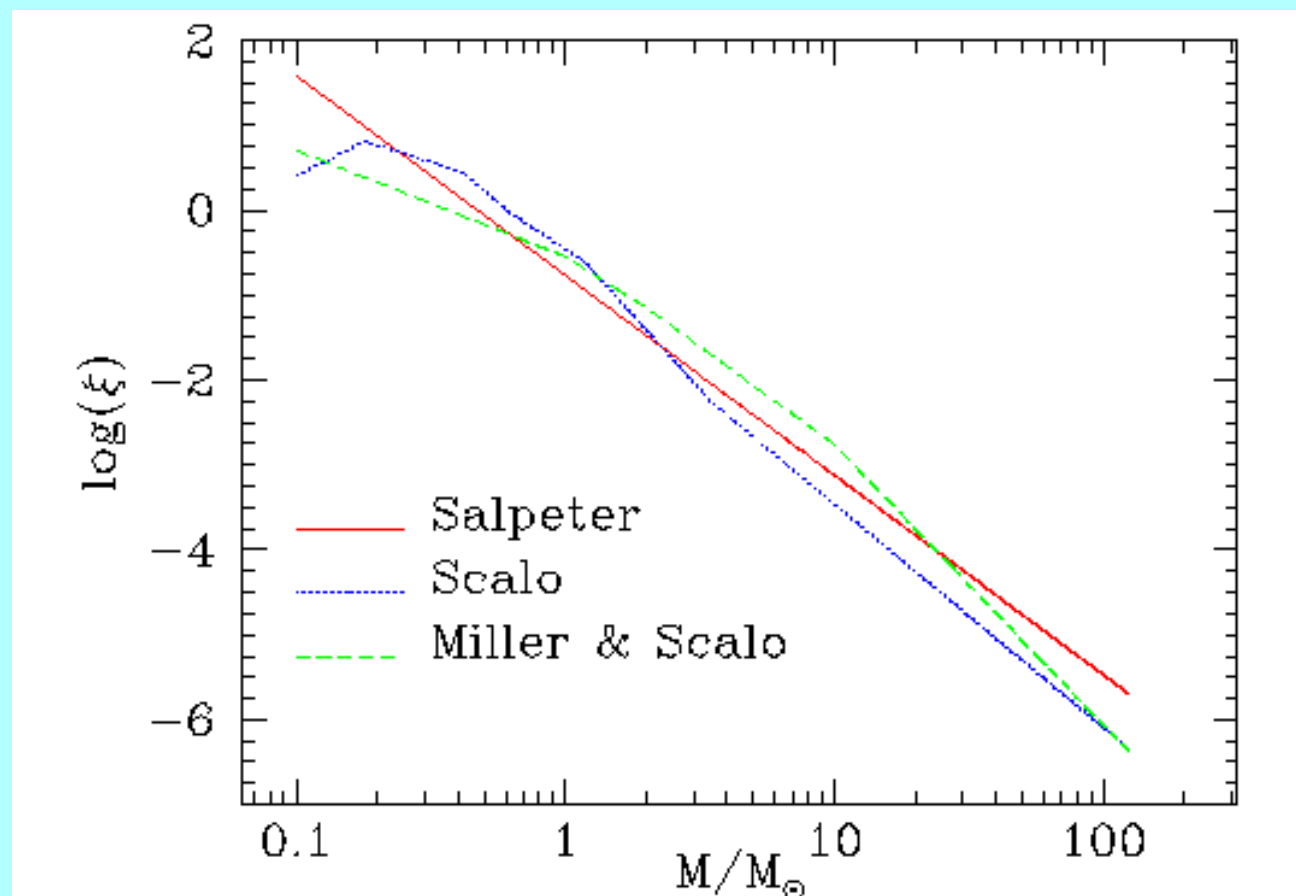
for

$$M_L < M < M_U$$

Usually $M_L = 0.1 M_\odot$ and $M_U = 50 M_\odot$.

The “Salpeter-function” has $x = 1.35$.

Here are some forms of the IMF often used.



Bi-model star formation was proposed by Larson*. It says that the two modes of star formation of high- and low-mass stars are independent and normalisation of the IMF must be done separately.

The **Star Formation Rate(SFR)** is the total mass in newly formed stars as a function of time.

In the solar neighborhood it has been approximately constant with time.

It may vary between galaxies, but is usually taken independent of position in a galaxy.

With an IMF and a SFR it is possible to calculate the luminosity and colors of galaxies as a function of time.

This is done by first calculating the photometric evolution of a star clusters by assuming the IMF and making use of stellar evolution tracks.

*Mon. Not. R.A.S. 218, 409 (1986)

In principle this needs to be done for different metal abundances.

These clusters can then be added according to the SFR (and the evolution of metal abundance with time).

Analytical model for photometric evolution.

a. Single burst.

First look at the Main Sequence; we have approximately:

$$L \propto M^{\alpha}$$

Rough values for α are 4.9 in U, 4.5 in B and 4.1 in V.

The main-sequence life-time is:

$$t_{\text{MS}} = M^{-\gamma}$$

With M in M_{\odot} the unit of time is $\approx 10^{10}$ years. A good value for γ is 3.

Assume that stars formed all at $t = 0$ and that the total mass is ψ_0 . Then

$$\begin{aligned} L_{MS}(t) &= \int_{M_L}^{M_t} \psi_0 M^\alpha \phi(M) dM \\ &= \frac{x}{\alpha - x} M_L^x \psi_0 M_t^{\alpha - x}, \end{aligned}$$

where

$$M_t = t^{1/\gamma}$$

Now look at the giants. Assume all giants have a luminosity L_G and are in that stage for a time t_G .

Reasonable values for L_G are 35 in U, 60 in B and $90 L_\odot$ in V and 0.03 for t_G .

The number of giants at time t is then

$$\begin{aligned} N_G(t) &= \psi_0 \phi(M_t) \left| \frac{dM}{dt_{MS}} \right|_{M=M_t} t_G \\ &= \psi_0 \frac{x}{\gamma} M_L^x M_t^{\gamma - x} t_G \end{aligned}$$

Now we can derive the **Single Burst** luminosity at time t :

$$L_{\text{SB}}(t) = L_{\text{MS}}(t) + N_{\text{G}}(t)L_{\text{G}}$$

Using $U_{\odot} = 5.40$, $B_{\odot} = 5.25$ and $V_{\odot} = 4.70$, and $M_{\text{L}} = 0.1M_{\odot}$, the following table can be calculated.

t	$(U - B)$	$(B - V)$	$(M/L)_{\text{B}}$
0.01	-0.34	0.12	0.15
0.03	-0.06	0.45	0.38
0.1	0.18	0.64	1.12
0.3	0.38	0.79	2.79
1	0.56	0.90	6.95
3	0.66	0.96	14.9

b. **Ongoing star formation.**

Write the SFR as $\psi(t)$. Then

$$L(t) = \int_0^t \psi(t - t') L_{\text{SB}}(t') dt'$$

For two extreme cases we get at $t = 1$:

Model	(U-B)	(B-V)	$(M/L)_B$
Single burst	0.56	0.90	7.0
Constant SFR	-0.25	0.24	1.0

This spans the range of the observed two-color diagram with the single burst corresponding to elliptical and S0 galaxies and the constant SFR for Sc and later types.

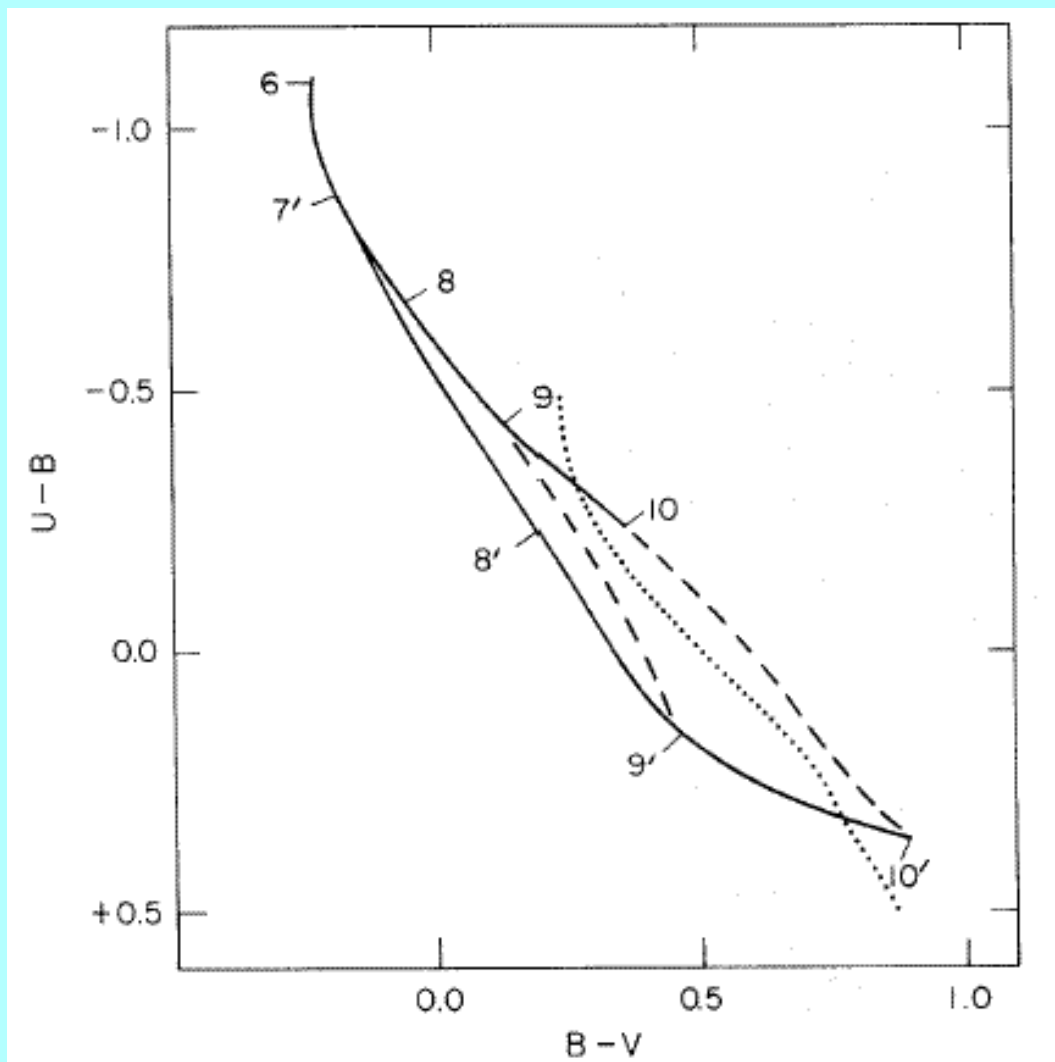
Now let us look at some more detailed studies.

Searle, Sargent & Bagnuolo* find the following luminosities and colors for a number of slopes of the IMF.

AGE (10^7 yrs)	$\alpha = 2.1$			$\alpha = 2.45$			$\alpha = 3.2$		
	M_V	$B - V$	$U - B$	M_V	$B - V$	$U - B$	M_V	$B - V$	$U - B$
0.1....	-7.2	-0.23	-1.18	-6.6	-0.22	-1.10	-5.6	-0.22	-1.10
0.3....	-7.6	-0.19	-0.96	-6.9	-0.21	-0.96	-5.9	-0.18	-0.96
1.0....	-7.4	-0.15	-0.83	-6.9	-0.18	-0.78	-6.0	-0.14	-0.78
3.0....	-6.2	-0.05	-0.60	-6.0	-0.03	-0.58	-5.6	-0.05	-0.58
10.0....	-5.0	+0.19	-0.22	-5.0	+0.19	-0.23	-5.0	+0.11	-0.23
30.0....	-3.8	+0.21	+0.03	-3.9	+0.34	0.00	-4.4	+0.27	+0.06
100.0....	-2.5	+0.44	+0.12	-2.8	+0.46	+0.16	-3.6	+0.47	+0.22
300.0....	-1.6	+0.66	+0.26	-1.9	+0.67	+0.28	-2.9	+0.68	+0.29
1000.0....	-0.9	+0.89	+0.38	-1.2	+0.90	+0.36	-2.3	+0.89	+0.36

*Ap.J. 179 ,427 (1973)

Using this they get a predicted two-color diagram with the Salpeter IMF as in the following figure. Here numbers x show the location of models of ages 10^x years old; with primes for SB models and unprimed for constant SF. All normal galaxies lie to the right of the dotted line.



Searle *et al.* conclude that the models and observations are consistent with:

- All galaxies $\approx 10^{10}$ years old.
- IMF everywhere similar to local IMF.
- Mean SFR averaged over sufficiently large area's and long times generally declines with time.
- Decay times vary among late-type galaxies; some show bursts, some show uniform SFR.

Larson & Tinsley* add:

- Precise form of SFR is not important. Important is only SFR over the last $\approx 10^8$ years to mean SFR over the life of the galaxy.
- Effects of different ages, metallicities and upper stellar masses are small.

*Ap.J. 219, 46 (1978)

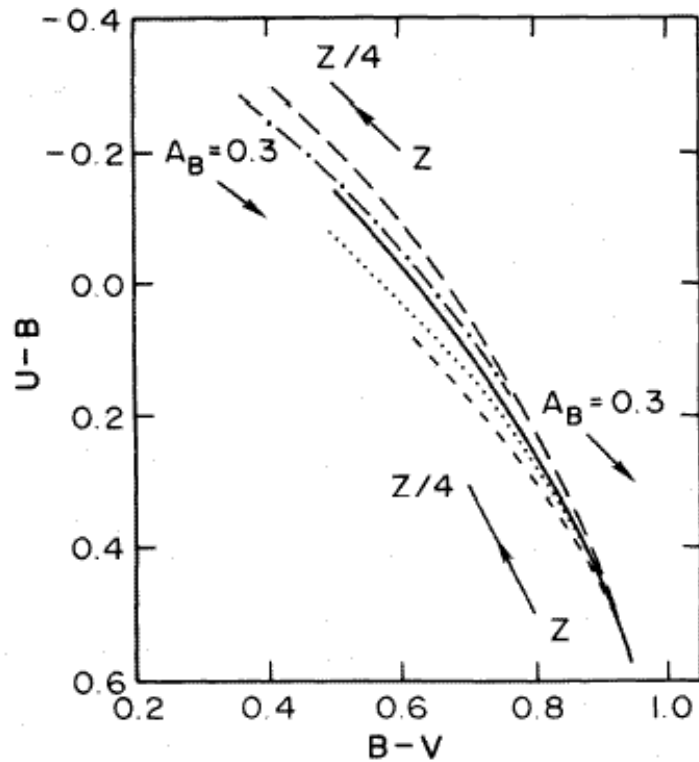
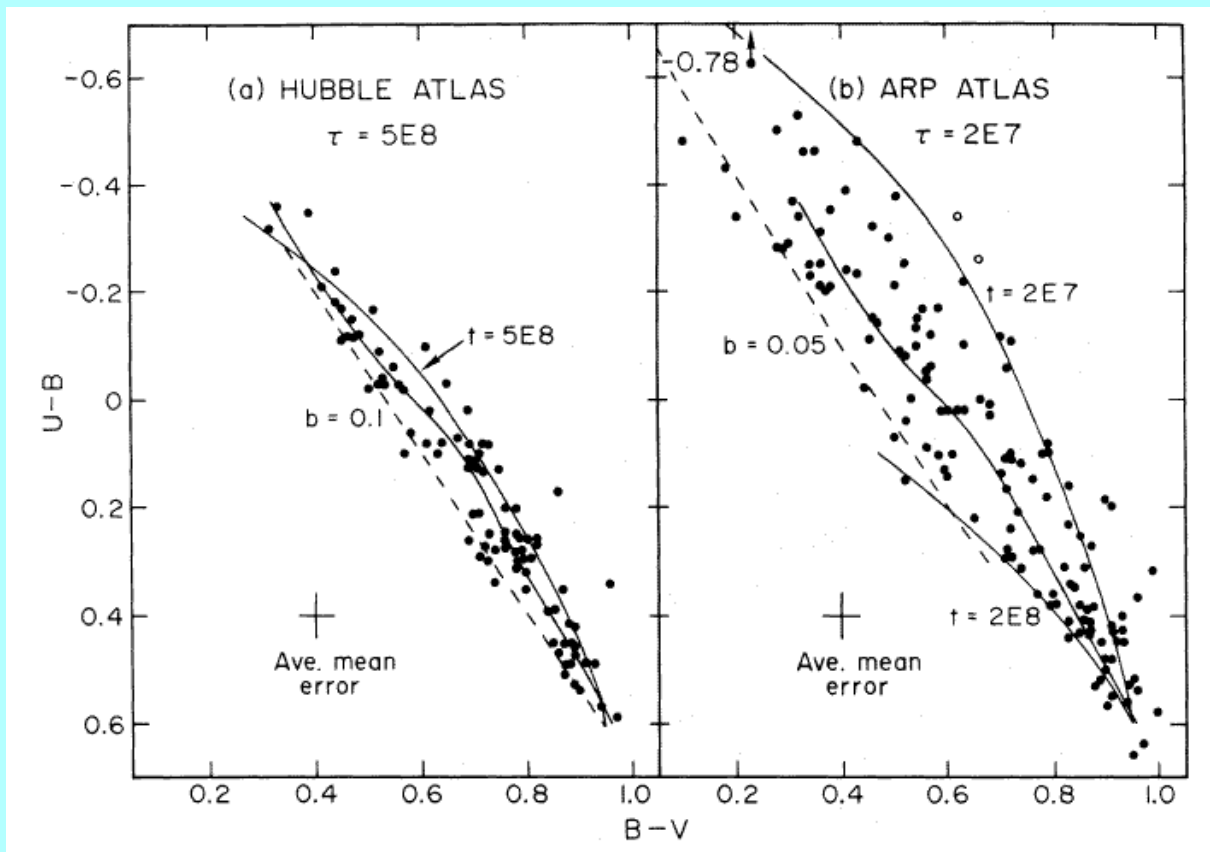


FIG. 7.—Colors of models with monotonic SFRs and age 10^{10} yr. *Heavy line*, local IMF. *Long dashes*, IMF with slope $x = 1$. *Short dashes*, $x = 2$. The foregoing use case T supergiant colors and have an upper mass limit $m_U = 30 M_\odot$. *Dot-dashes*, $x = 1$, $m_U = 30 M_\odot$, and case C supergiant colors. *Dots*, $x = 1$, case T, and $m_U = 10 M_\odot$. The reddening vectors for $A_B = 0.3$ show the RC2 formula for galactic reddening which depends on $B - V$. The other vectors indicate schematically how colors of red and blue galaxies, respectively, may change with a factor 4 reduction in metallicity.

- Interacting galaxies show more scatter in two-color diagram. This can be explained with bursts of 5% (fraction of mass to total stellar mass at time of burst; $b \sim 0.05$) and duration $\tau \approx 2 \times 10^7$ years.



Kennicutt* adds the integrated H_{α} fluxes (in the form of an equivalent width[†]), providing independent information on recent formation of heavy stars.

His most important results are the following:

*Ap.J. 272, 54 (1983)

[†]Equivalent width is the wavelength interval in the continuum that corresponds to as much flux as the line

- Slope of upper IMF roughly as Salpeter function.

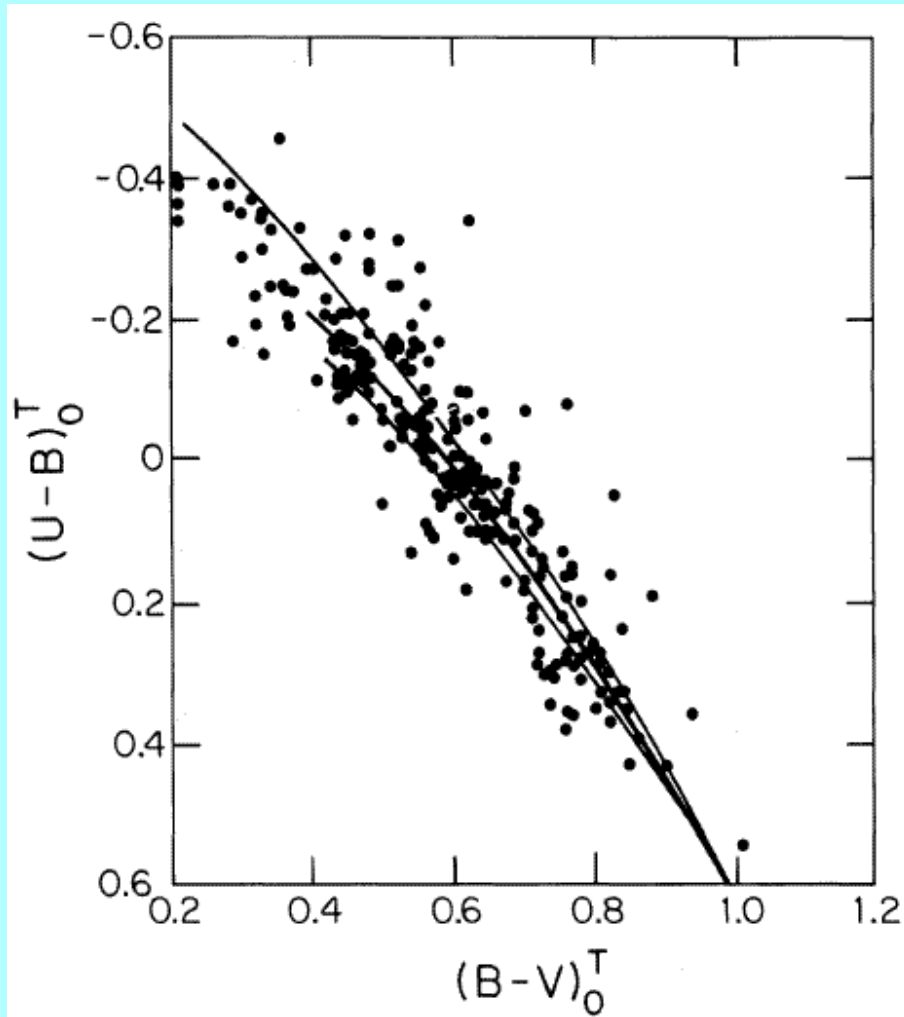
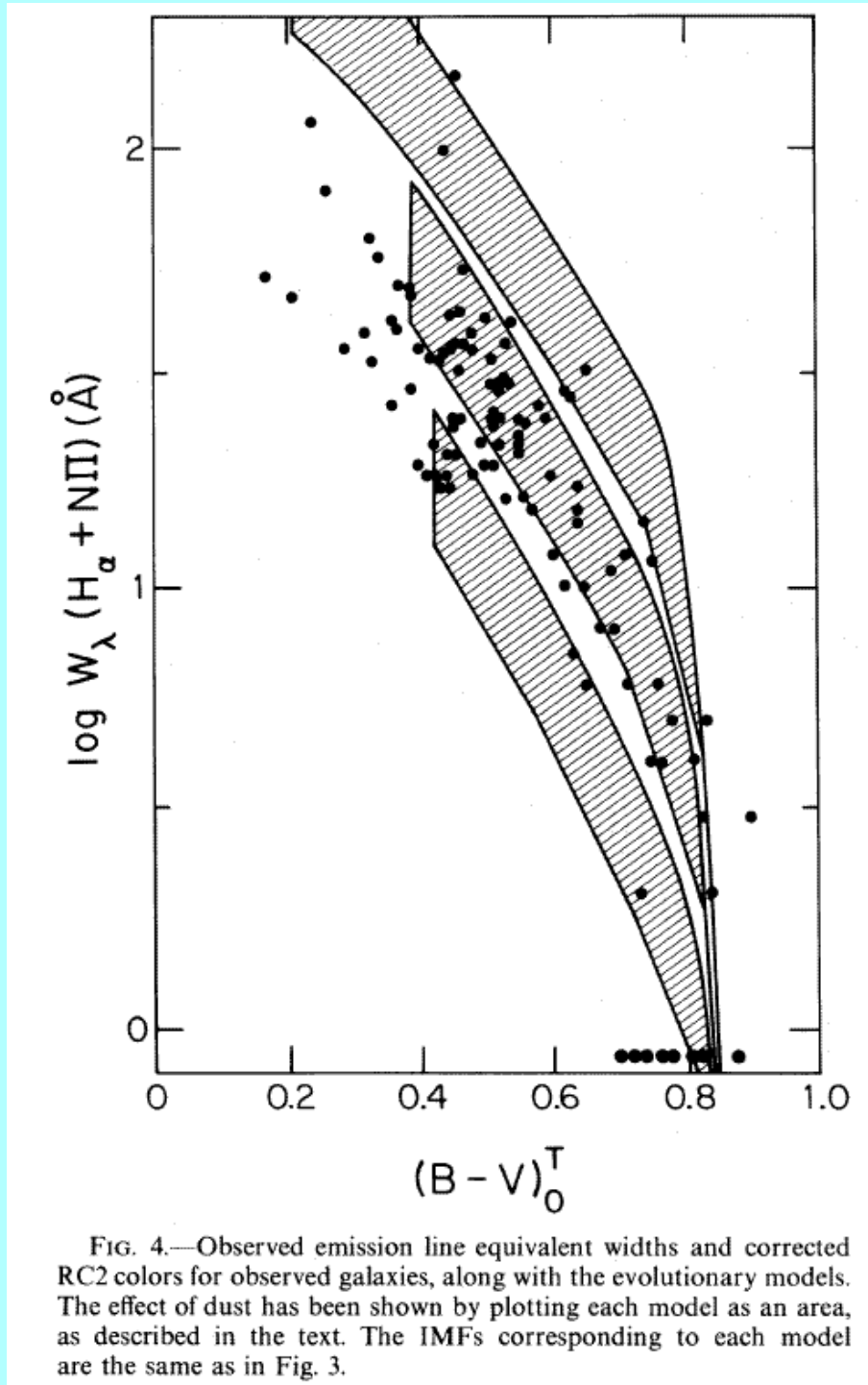
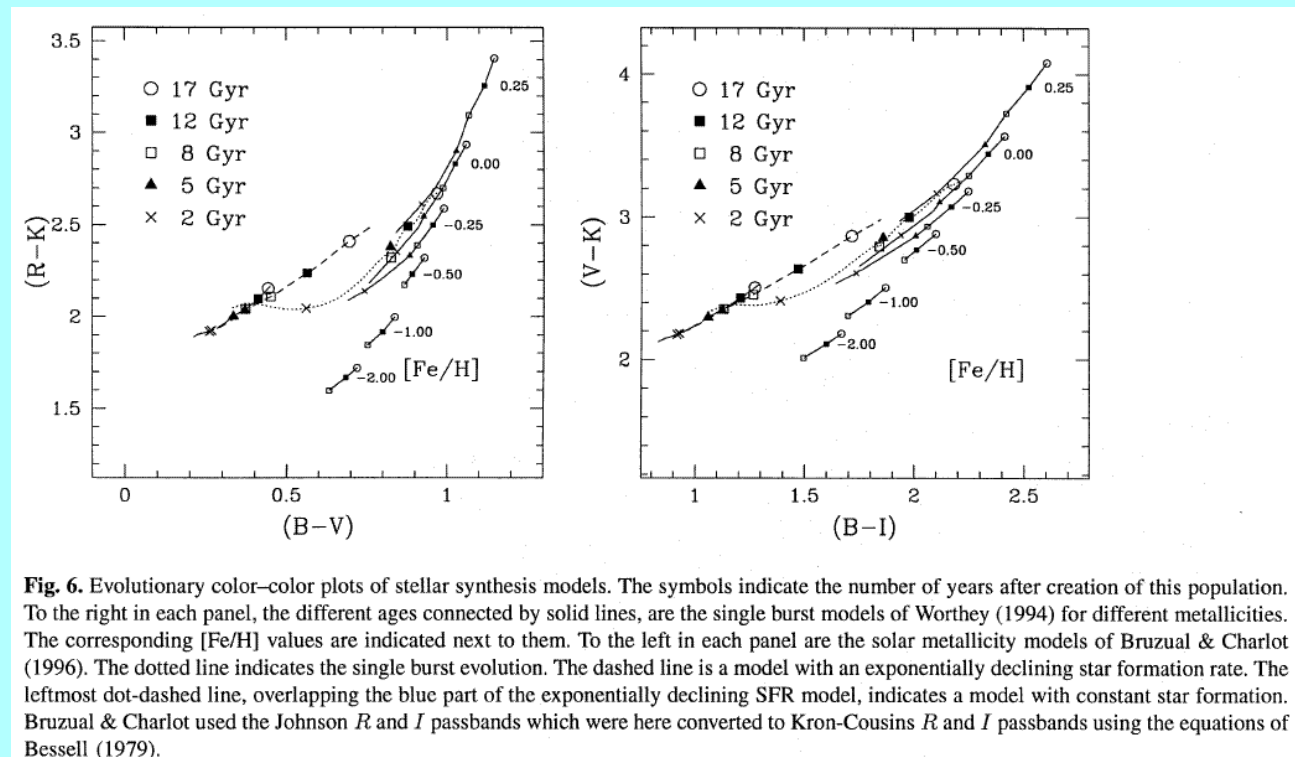


FIG. 3.—Two-color diagram from Shapley-Ames spiral galaxies, along with the model galaxy disk colors described in the text. The three curves correspond to the different mass functions adopted, the Miller and Scalo function (*lowest curve*), the extended Miller-Scalo (i.e., “Salpeter”) function, and the shallow m^{-2} IMF (*top curve*).

- Galaxies have same upper mass limit in IMF ($\approx 50M_{\odot}$).



De Jong* derives models to study the color gradients in disks and among different disks.



- Dust reddening plays a minor role.
- Outer parts have lower average ages and are more metal poor than inner parts of disks.
- Late type galaxies ($T \geq 6$) have lower metallicities and younger average ages.

*A.&A. 313, 377 (1996)

SFR and Schmidt law.

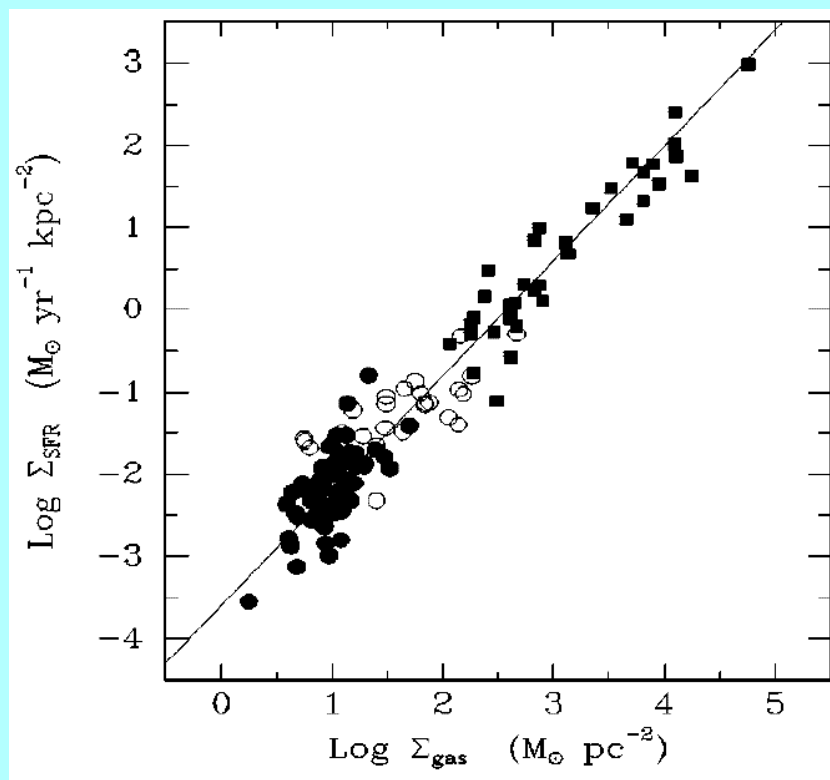
Schmidt* proposed that the star formation rate relates to the gas density as

$$\text{SFR} \propto \rho^2$$

Often this was immediately translated in (observable) surface properties.

The latest result[†] gives

$$\Sigma_{\text{SFR}} = (2.5 \pm 0.7) \times 10^{-4} \left(\frac{\Sigma_{\text{gas}}}{1 M_{\odot} \text{pc}^{-2}} \right)^{1.4 \pm 0.15} M_{\odot} \text{year}^{-1} \text{kpc}^{-2}$$



*Ap.J. 129, 243 (1959)

[†]Kennicutt, Ann.Rev.Astron.Astrophys. 36, 189 (1998)

Population synthesis.

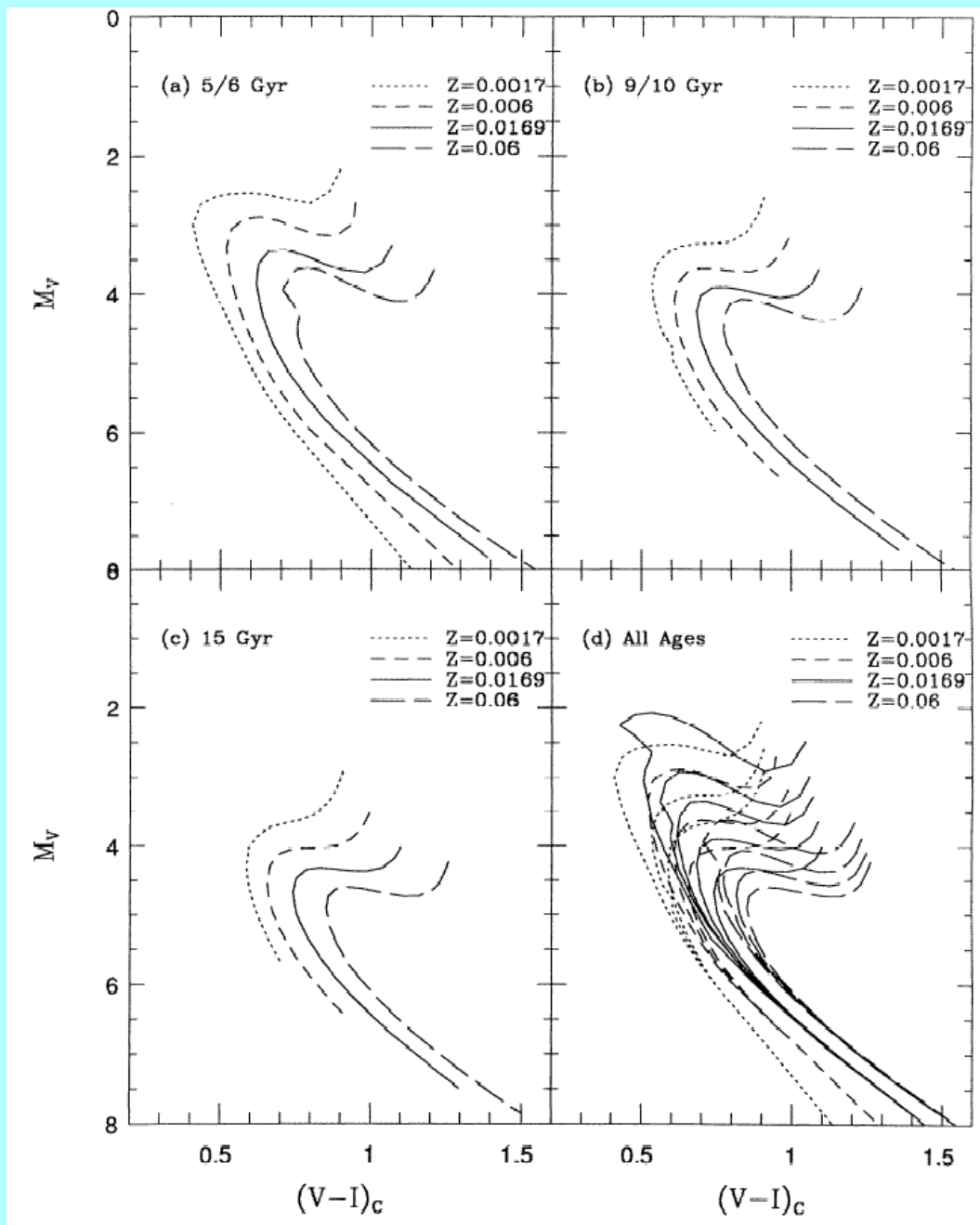
Attempt to fit the intermediate resolution spectra with those of observed stars. Best method now is by fitting integrated spectra of generations of particular age and metallicity*.

The steps are the following.:

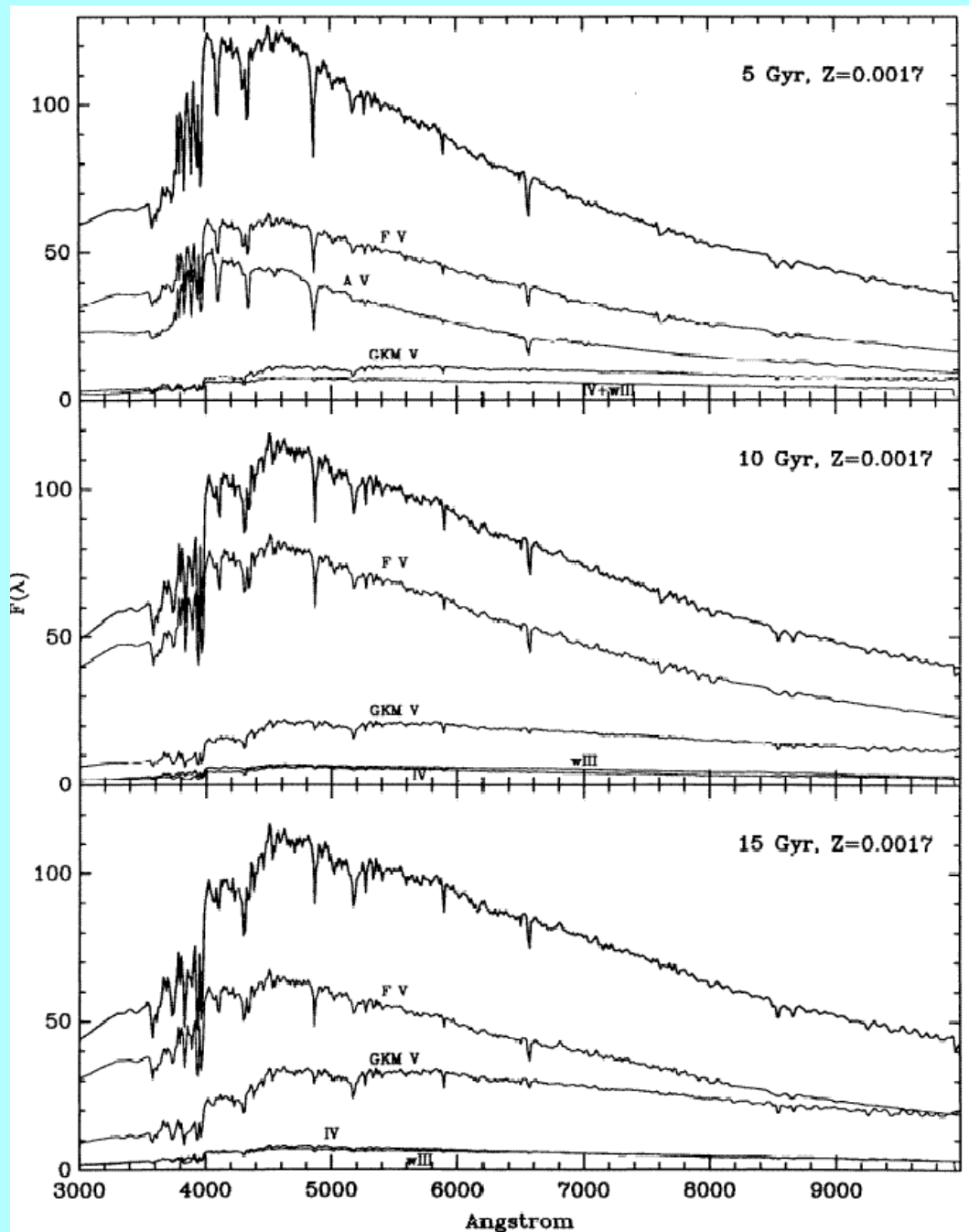
- Measure spectra of stars of various ages and metallicity.
- Synthesize integrated spectra of generations from a set of isochrones.
- Fit using least-squares techniques to galaxy spectra.

*For example Pickles Ap.J. 296, 340 (1985); Ap.J.Suppl. 59, 33 (1985) or Pickles & van der Kruit A.&A.Suppl, 84, 421 (1990) and 91, 1 (1991)

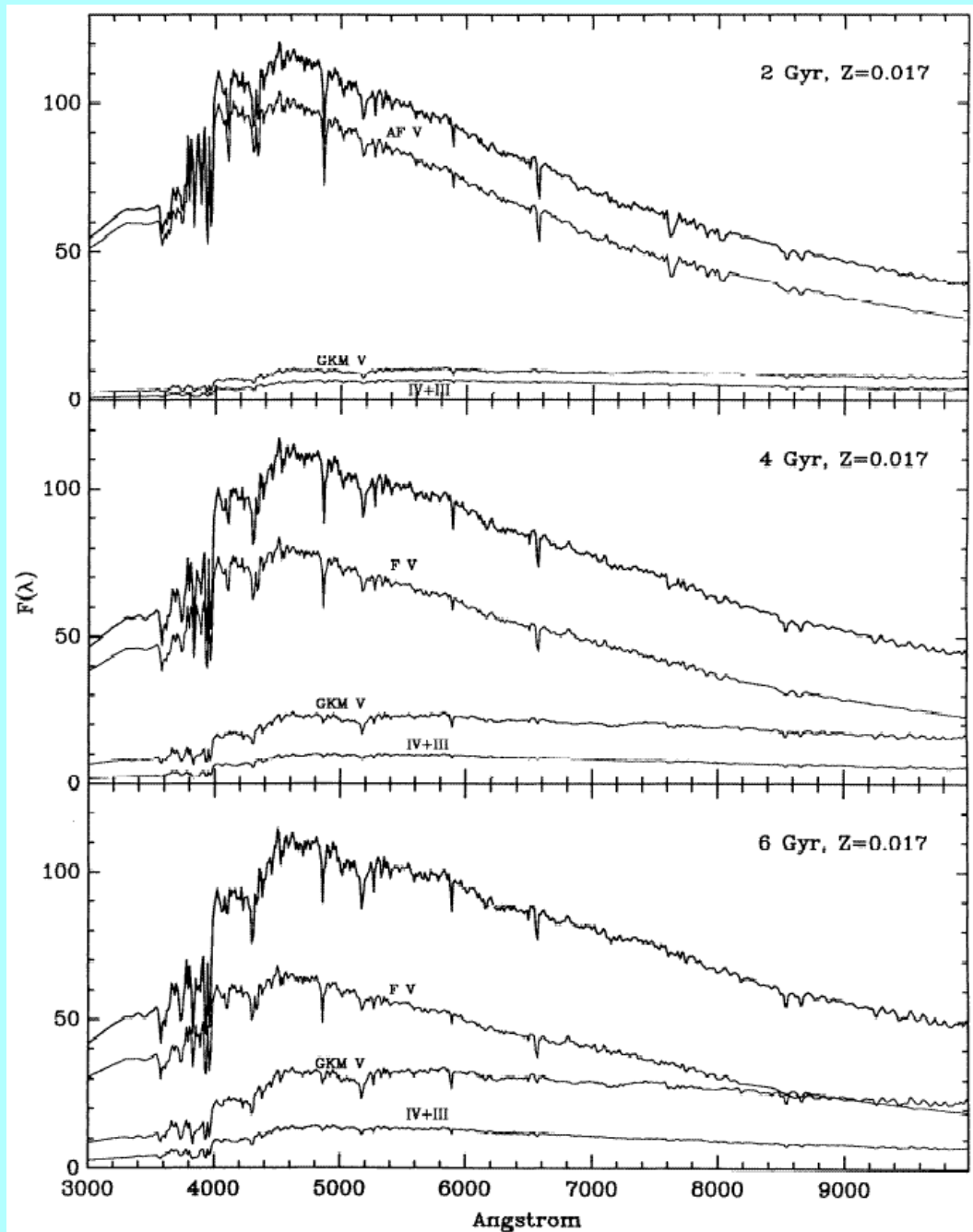
Here is a set of isochrones used by Pickles & van der Kruit.



These are synthesized spectra of a metal poor cluster at three ages.



These are synthesized spectra of a metal rich cluster at three ages.



This is an example of a spectrum of an elliptical galaxy fitted by a set of stellar spectra.

